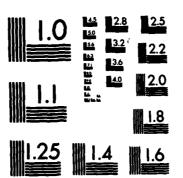
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# NAVAL POSTGRADUATE SCHOOL Monterey, California



## **THESIS**



AERODYNAMICALLY EFFICIENT GRADIENT REFRACTIVE INDEX MISSILE SEEKER LENS

by

Herbert M. Carr III

October 1982

Thesis Advisor:

A. E. Fuhs

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This thesis explores the use of a pointed seeker lens designed using a spherically symmetric gradient refractive index (GRIN). The design helps to solve the current design conflict between optical quality and aerodynamic drag inherent in hemispherical seeker lenses. Equations for lens design and the evaluation of off-axis lens performance have been developed for both a homogeneous version and a GRIN version of the pointed

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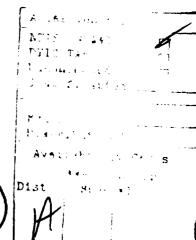
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Aerodynamically Efficient Gradient Refractive Index Missile Seeker Lens

by

Herbert M. Carr III
Captain, United States Army
B.S., University of Texas, 1971

Submitted in partial fulfillment of the requirements for the degree of

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from the
NAVAL POSTGRADUATE SCHOOL
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This thesis explores the use of a pointed seeker lens designed using a spherically symmetric gradient refractive index (GRIN). The design helps to solve the current design conflict between optical quality and aerodynamic drag inherent in hemispherical seeker lenses. Equations for lens design and the evaluation of off-axis lens performance have been developed for both a homogeneous version and a GRIN version of the pointed seeker lens. The homogeneous lens is used as a comparison and a check for the GRIN lens. A FORTRAN program (GISL) has been written and employed to evaluate and compare both the homogeneous lens and many different configurations of possible GRIN lens designs. Results indicate that the GRIN lens has highly superior off-axis imaging performance as compared to the homogeneous lens. The best results were obtained for the GRIN lens with a fifty percent, positive, spherically symmetric gradient index with center of symmetry interior to the lens. Only very slightly inferior performance was observed with a five percent version of the same lens; such a lens possibly can be manufactured today. GRIN lens performance also indicates that for objects offaxis by more than 17.2 degrees a large scale, multiple element sensor array may be required; with such a sensor array, objects off-axis by more than 37.2 degrees may require mirror elements to compensate for image movement.

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#### LIST OF SYMBOLS

Symbol in Equations	FORTRAN Symbol	Equation Number Where First Introduced	Definition	Units
a	A	113	Parameter of gradient index function	Nondimensional
a		41	Defined by Equation 44	
A		24	Defined by Equation 25	
A		137	Defined by Equation 138	
A <sub>1</sub>		161	Defined by Equation 164	
A <sub>2</sub>		160	Defined by Equation 161	
A <sub>3</sub>		169	Defined by Equation 170	
A <sub>5</sub>		179	Defined by Equation 184	
A <sub>6</sub>		178	Defined by Equation 181	
AB	AB	30	Line segment from A to B, from nose to apex of cone	Nondimens: anal
b	В	113	Parameter of gradient index function	Nondimensional
b		41	Defined by Equation 45	
В		24,137	Defined by Equation 26	
<b>B</b> 1		161	Defined by Equation 165	
B <sub>2</sub>		160	Defined by Equation 162	
B <sub>5</sub>		179	Defined by Equation 185	
<sup>B</sup> 6		178	Defined by Equation 182	
BF	BF	2	Line segement from B to F, focal length	Nondimensional
C		41	Defined by Equation 46	

С		24	Defined by Equation 27	
С		137	Defined by Equation 139	
$c_1$		161	Defined by Equation 166	
c <sup>2</sup>		160	Defined by Equation 163	
c <sub>6</sub>		178	Defined by Equation 183	
đ		41	Defined by Equation 47	
đ	ERROR	142	Error parameter in meridian plane ray intercept with outside surface	Nondimensional
D	D2	72	Geometrical length of skew ray in homogeneous lens	Nondimensional
D'	D3	95	Geometrical length of skew ray from cone to image plane	Nondimensional
DYDNN (J)	DYDXN (J)	20	Slope of normal to out- side surface in the meridian plane at the Jth point	Nondimensional
DYDXT (J)	DYDXT(J)	21	Slope of outside sur- face tangent in the meridian plane at the Jth point	Nondimensional
DYDKN <sub>PIM</sub>		51	Slope of normal to out- side surface in the meridian plane. Interpolated value	Nondimensional
e	E	108	Spherical CRIN scalar invariant	Nondimensional
e		41	Defined by Equation 48	
f(x,y,z)			General function of $x$ , $y$ , and $z$	
f <sub>x</sub>		82	Partial derivative of function f with respect to x	

fy		82	Partial derivative of function f with respect to y	
f <sub>z</sub>		82	Partial derivative of function f with respect to z	
i		186	Summation index	
î		53	Unit vector in the x direction	
I	I	3	Number of angular incre- ments in lens design algorithm	
IA			Fraction of absorbed radiant energy during transmission through the lens	Nondimensional
<sup>1</sup> 1	n	17	Angle of incidence with respect to local normal at outside surface	Radians
I'i	IlP	16	Angle of refraction with respect to the local normal at the outside surface	Radians
<sup>1</sup> 2	12	12	Angle of incidence with respect to the local normal at the inside surface	Radians
I'2	I2P	10	Angle of refraction with respect to the local normal at the inside surface	Radians
IR		102	Fraction of reflected radiant energy at the surface interface	Nondimensional
I <sub>T</sub>	NINCTY	102	Fraction of transmitted radiant energy at the interface	Nondimensional
ı	MINCTY (G)	190	Intensity element of intensity summation	Nondimensional

<sup>I</sup> av	XMAVE	190	Average value of ray intensity	Nondimensional
ĵ		53	Unit vector in the x-direction	
J	J	4	Ray Number index	
k	LK ,	63	Direction cosine, x- direction, of outside surface normal	
k'	LKP	83	Direction cosine, z- direction, of inside surface normal	
k		55	Unit vector in y- direction	
K	CIK	58	x-direction cosine of ray exterior to lens surface	
K'	CKP.	63	x-direction cosine of ray inside the lens	
K"	CXPP	92	x-direction cosine of the ray after refraction by the lens	
1	LL	64	y-direction cosine of outside surface normal	
1'	LIP	84	y-direction cosine of inside surface normal	
L	CI	58	y-direction cosine of the ray external to the lens	
L'	CLP	64	y-direction cosine of the ray internal to the lens	
L"	CLPP	93	y-direction cosine of the ray after diffraction by the lens	

10		122	Generalized z-direction cosine of the initial ray direction in GRIN	
m	LM	65	z-direction cosine of outside surface normal	
m		38	Generalized slope of a line segment	
m*	LMP	85	z-direction cosine of inside surface normal	
M	CM	58	z-direction cosine of ray external to the lens	
M'	CMP	65	z-direction cosine of ray internal to the lens	
M"	CMPP	94	z-direction cosine of ray after refraction by the lens	
n		108	Generalized local value of the gradient refractive index	Nondimensional
n <sub>0</sub>		109	Generalized value of the gradient index at the initial intercept point	%ndimm==onal
n <sub>2</sub>	N2	125	Local lens interior value of the GRIN	Nondimensional
N		186	Generalized nth value of summation index	
N <sub>1</sub>	Nl	1	Homogeneous index of refraction of medium 1, exterior to lens	Nondimensional
N <sub>2</sub>	N2	1	Homogeneous index of refraction of medium 2, interior to lens	Nondimensional
N <sub>3</sub>	N3	12	Homogeneous index of refraction of medium 3, behind the lens	Nondimensional
N'		55	Normal vector at the point of intersection at the outside surface	

$\overline{N}_{m}$		53	Normal vector in the meridian plane correponding to $\overline{N}^{\bullet}$	
$\hat{\mathbf{N}}_{\mathbf{i}}$		80	Generalized surface normal unit vector	
N <sub>ti</sub>		104	Ratio of indices of refraction at an interface	Nondimensional
Ñ <sub>po</sub>		148	Unit normal vector to plane of skew ray	
N <sub>pox</sub>	NPOX	149	Defined by Equation 150	
N <sub>poy</sub>	NPOY	149	Defined by equation 151	
N poz	NPOZ	149	Defined by Equation 152	
08	osymb	126	Line segment from the center of symmetry of the GRIN function, O <sub>s</sub> , to B, the origin	Nondimensional
o <sub>s</sub>		112	Location of the GRIN center of symmetry on the lens axis	Nondimensional
P <sub>0</sub>		120	Generalized initial x- direction cosine of the GRIN ray	
$P_1$		76	Defined by Equation 77	
P <sub>2</sub>		76	Defined by Equation 78	
P <sub>3</sub>		76	Defined by Equation 79	
PARL	PARL	40	Defined by Equation 41	
PAR2	PAR2	40	Defined by Equation 42	
PAR3	PAR3	40	Defined by Equation 43	
₫ <sub>0</sub>		121	Generalized y-direction cosine of initial GRIN ray direction	
Q			Perpendicular distance between ray and line parallel to ray through the origin at B in Figure 3	Nondimensional

QP.		37	Line segment defined by Figure 5	Nondimensional
QP'		37	Line segment defined by Figure 5	Nondimensional
r	RAD	108	Radial coordinate in GRIN from O <sub>S</sub>	Nondimensional
r <sub>0</sub>	RO	108	Initial value of r at ray intercept point	Nondimensional
r <sub>0</sub>		146	Unit vector in direction from $O_S$ to intercept point on outside surface	
rox	R0X	151	x-direction cosine of $\hat{\mathbf{r}}_0$	
r <sub>oy</sub>	ROY	150	y-direction cosine of r <sub>0</sub>	
r <sub>oz</sub>	ROZ	150	z-direction cosine of $\hat{\mathbf{r}}_0$	
r <sub>ih</sub>		156	Unit vector from O <sub>s</sub> to point of homogeneous intercept on inside surface	
r <sub>g</sub>	RADH	160	Geometrical radius from O <sub>s</sub> to cone during iteration for the inside surface intercept	Nondimensional
r <sub>PI</sub>	RAD	174	Radius to inside surface intercept as found by interation	Nondimensional
r <sub>11</sub>		103	Reflection coefficient for parallel E vector	Nondimensional
r		104	Reflection coefficient for perpendicular E vector	Nondimensional
R	R	2	Maximum inside radius of cone measured from the lens axis	Meters
R	RAD	33	Radius of circle in the y-z plane	Nondimensional
Ŕ		57	Unit vector in the ray direction	

R'		62	Vector in the direction of refracted ray	
Rz	RZERO	113	Maximum possible radius in GRIN	Nondimensional
s	s	51	Line segment in the meridian plane; see Figure 5	Nondimensional
ST	ST	51	Inclusive line segment in the meridian plane; see Figure 5	Nondimensional
T	T	14	Edge thickness of the lens	Nondimensional
ŭ	U	17	Direction of the ray with respect to the lens axis in the meridian plane	Radians
יט•	UP	13	Direction of the ray inside the lens with respect to the lens axis in the meridian plane	Radians
U"	UDP	4	Angle between ray and lens axis at the focal point in the meridian plane	Radians
v		115	Integration variable	Nondimensional
x		22	General x-coordinate along lens axis	Nondimensional
x'		30	General x-coordinate in the grid plane (tilted)	Nondimensional
$ \mathbf{x} $		2	Absolute value of any quantity x	
<b>*</b> 0	ЖO	40	x-coordinate of outside surface skew ray intercept	Nondimensional
×i	хI	69	x-coordinate of inside surface intercept	Nondimensional

×im	XIM	95	x-coordinate in the image plane (Spot Diagram)	Nondimensional
ж <sub>ІН</sub>	хін	130	First x-coordinate of imaginary homogeneous intercept in GRIN design	Nondimensional
x <sub>p</sub>	ХP	135	x-coordinate of the first intermediate point on the GRIN ray during iteration in the design algorithm	Nondimensional
кiн	XLHP	141	Second x-coordinate of imaginary homogeneous intercept in GRIN design	Nondimensional
 х,,,	ХР	143	Successive values of x during iteration	Nondimensional
*PIM		52	x-coordinate of ray intercept in meridian plane	Nondimensional
х <sub>1</sub> (J)	X1 (J)	14	x-coordinate of the Jth ray on the outside surface in the meridian plane	Nondimensional
х <sub>2</sub> (J)	X2(J)	5	x-coordinate of the Jth ray on the inside sur- face in the meridian plane	Nondimensional
У		22	General y-coordinate (vertical)	Nondimensional
у'		31	y-coordinate in the tilted grid plane	Nondimensional
y <sub>0</sub>	YO	49	y-coordinate of outside surface skew ray intercept	Nondimensional
Yi	YI	170	y-coordinate of inside surface skew ray intercept	Nondimensional

y <sub>im</sub>	MIK	96	y-coordinate of the skew ray in the image plane (Spot Diagram)	Nondimensional
у <sub>ІН</sub>	АТН	131	First y-coordinate of imaginary HIN inter-cept in GRIN design (iteration)	Nondimensional
Y <mark>lH</mark>	Y1HP	137	Second y-coordinate of imaginary HIN intercept in GRIN design (iteration)	Nondimensional
уp	YP	136	x-coordinate of the first intermediate point on the GRIN ray during iteration in the design algorithm	Nondimensional
у <sub>р</sub>	YP	144	Successive values of yp during iteration	Nondimensional
У <sub>С</sub>	YŒNTR	186	y-coordinate of the image centroid in the Spot Diagram	Nondimensional
Yi	YIM(G)	186	Individual y-coordinates of image plane ray intercept	Nondimensional
у <sub>1</sub> (J)	Yl(J)	15	y-coordinate of the Jth ray on the outside surface in the meridian plane	Nondimensional
y <sub>2</sub> (J)	<b>J2(J)</b>	5	y-coordinate of the ray on the inside surface in the meridian plane	Nondimensional
z		32	General z-coordinate (horizontal)	Nondimensional
z '		32	z-coordinate in the grid plane (tilted)	Nondimensional
<b>z</b> <sub>0</sub>	20	50	z-coordinate of the skew ray on the out- side surface	Nondimensional
z <sub>i</sub>	21	71	z-coordinate of the skew ray on the inside surface (cone)	Nondimensional

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<sup>2</sup> im	ZIM	97	z-coordinate of the skew ray in the image plane (Spot Diagram)	Nondimensional
œ.	ALPHA	2	Cone half-angle	Radians
αp	ALFAP	30	Grid plane tilt angle with respect to the lens axis	Radians/Degrees
β	BETA	2	Total angle between lens axis and the point of maximum radius of the cone; measured at the focal point	Radians
Υ	CAMMA	29	Nose half-angle (opaque region)	Radians
δ		120	Defined by Equation 124	
$\Delta U^{\omega}$	DLUDP	3	Angle between successive rays at the focal point	Radians
Δ <b>x</b>		52	Incremental change in x; see Figure 5	
Δ <b>y</b>		52	Incremental change in y; see Figure 5	
ε	EPSILON	109	Sign function (±1)	
ζ	ZETA	134	Angle between instantan- eous GRIN ray direction and the lens axis in the meridian plane	Radians
η		120	Defined by Equation 123	
θ		108	Generalized angular co- ordinate of GRIN ray in the plane of the ray	Radians
θ1		1	Generalized angle of incidence with respect to the local normal	Radians
θ2		1	Generalized angle of refraction with respect to the local normal	Radians
θ <sub>0</sub>	THETA0	108	Reference (at surface intercept) in GRIN	Radians

$\theta^{\mathbf{H}}$	THETAH	125	Iteration trial value of $\theta$	Radians
$\boldsymbol{\theta_{T}}$		125	Total angular GRIN coordinate	Radians
θp	THETAP	173	Trial values of $\theta$ during iteration	Radians
θ <b>.</b>	MENX	173	Revised trial value of $\theta$ during iteration	Radians
μ		54	Angle between meridian plane and the point of intersection on the outside surface	Radians
π	PI	13	Proportionality factor between the circumference and the diameter of a circle; 3.14159	Radians
°r	RMSRAD	189	Spot size; RMS radius of image	Nondimensional
σ <sub>Y</sub>	SIGMAY	188	Standard deviation of y-coordinates of rays in the Spot Diagram	Nondimensional
σ <b>z</b>	SIGMAZ	187	z-standard deviation of rays in the Spot Diagram	Nondimensional
ф	PHI	57	Angle between skew ray and the outside surface with respect to the surface normal	Radians
φ'	PHIP	61	Refraction angle be- tween skew ray and the outside surface normal	Radians
Φį	PHII	90	Angle of skew ray inter- cept with respect to the inside surface normal	Radians
Φ <u>:</u>	PHIIP	91	Refraction angle of skew ray with respect to the inside surface normal	Radians

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Ψ	PSI	118	Angle between skew ray direction and radial direction from O <sub>g</sub>	Radians
Ψο	PSI0	109	Initial value of angle $\psi$ at point of intercept at outside surface	Radians

## **ACKNOWLEDGEMENTS**

The author would like to express his sincere appreciation to Distinguished Professor Allen E. Fuhs for his invaluable guidance, friendship, and time away from sabbatical. Without him this thesis would not have been possible.

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# I. INTRODUCTION

Historically, the design of tactical missiles employing passive or semiactive infrared (IR) seekers has involved a difficult compromise between aerodynamic requirements and optical or seeker requirements. Whereas aerodynamically the missile nose region should be sharp in order to reduce drag, optically it should be hemispherically shaped for image quality and as large as possible to increase aperture and therefore acquisition or tracking range. Some IR homing missiles designed for very short range anti-armor missions have totally ignored nose drag in order to optimize seeker performance while other designs for longer range missiles requiring high cruise velocities and greater aerodynamic efficiency have used the blunted ogive as a compromise. have not been any IR designs which have ignored optics in favor of aerodynamics; nor has there been employed a pointed seeker lens with the desired optical qualities.

In order to increase the performance of optically guided missiles beyond the current state of development, the conflict between aerodynamic requirements and optical restrictions must be resolved. Significant improvement in missile performance by increasing thrust is not likely due to the highly advanced state of propulsion today. One way to resolve the aerodynamics-optics problem is to design a pointed lens which has, if not

imaging quality, enough optical performance to allow the reduction of tracking data.

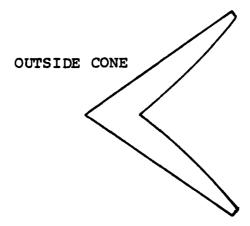
The lens must be in the general shape of a cone or ogive with half angle compatible with the design free stream Mach number in order to maintain an attached shock wave. Such a spike-shaped lens might also be used in the diffuser portion of an optically guided ramjet with nose inlet to conserve stagnation pressure. A diffuser-lens may benefit by a semi-isentropic spike shaped lens.

Poor optical performance due to the pointed shape of sharp lenses has precluded their use. Gradient Refractive Index (GRIN) materials, however, permit the lens designer the freedom to spatially vary the lens index of refraction to compensate for a traditionally poor optical shape. Although the use of GRIN has seen widespread use in fiber optics technology, it has been used infrequently in lens applications until recently. Lens designers are discovering that multipleelement photographic objectives may be redesigned using a two-element gradient lens [1]. It should be noted, however, that such GRIN lens have not yet been successfully fabricated even though large index changes in glass have been accomplished by the diffusion of doped electropolarizable ions [2]. Extensive research is being conducted in the creation of ever larger and more precise gradients. A total change in refractive index of approximately five percent is currently attainable.

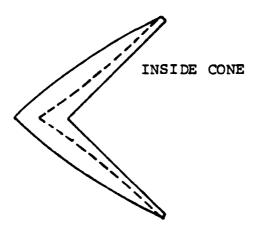
At the Naval Postgraduate School, Frazier [3], Terrell [4], and Amichai [5] have studied GRIN as applied to the sharp lens problem. The brief introduction by Frazier was followed by Terrell who designed a sharp lens having a conical outside surface and a variable inner surface using a homogeneous index (HIN) and then also briefly touched on the GRIN application. Amichai extended Terrell's lens to the GRIN case and included preliminary results from this thesis in a computer routine intended as a framework for lens optimization by following researchers.

This thesis investigates the design and off-axis performance (skew rays) of a GRIN seeker lens having a variable outer surface and a fixed right circular cone as an inside surface. The variable outside surface is determined by the character of the spherical gradient employed and varies from a pseudo ogive to an approximate isentropic spike as shown in Figure 1. First, the theory of the HIN lens is developed, followed by the design of a homogeneous lens intended both as a comparison and a check for the GRIN lens theory and the design which follows. Lens performance parameters are discussed and results presented for both the HIN and GRIN lenses.

Principles and and a series



(a) Terrell's Lens With Conical Outer Surface



(b) Lens Shapes In This Thesis With Conical Inner Surface.

Figure 1. Lens Shapes

# II. THE HOMOGENEOUS LENS

#### A. THEORY

Snell's law is the cornerstone of contemporary lens design in homogeneous optical materials. In the HIN case

Snell's law is used in the familiar form

$$N_1 \sin \theta_1 = N_2 \sin \theta_2 \tag{1}$$

where  $N_i$  is the index of refraction of the material corresponding to surface intercept angle  $\theta_i$ , with respect to the surface normal, at the interface between surfaces. Lens geometry and the relative values of  $N_i$  determine the resultant optical behavior.

### B. ASSUMPTIONS AND SIGN CONVENTION

In order to simplify the design and analysis of the seeker lens problem, certain assumptions have been introduced. Although energy loss upon transmission through the lens at each surface is calculated, it is assumed that the light is monochromatic radiation, time dependent electric and magnetic fields. Light impinging upon the lens is assumed to have a planar wave front as if propagating from an object at infinity. The presence of a shock wave attached to the lens is ignored as are any other regions of expansion or compression in the flow field about the lens [6]. Furthermore, the index of refraction of the free stream is assumed

to be equal to that of the interior space behind the lens even though computer routines were written with the flexibility to process unequal values.

The sign convention used is a right handed system with spatial coordinates positive to the right, up, and out of the page as seen by the reader. All angles are assumed positive counterclockwise from point of reference unless otherwise noted. All linear dimensions are implicitly non-dimensionalized with respect to the maximum radius of the right circular cone forming the inside surface of the lens.

### C. HIN LENS DESIGN

The lens design procedure consists of calculating both the loci of points forming the outside surface and the slope of the surface at each point in the meridian plane. It is convenient to approach the problem by placing a point source of light at the design focal point F and calculating successive refracted ray paths U' using Snell's law at points  $D_i$  on the inside surface in Figure 2. Points  $E_i$  are formed by the intercept of the refracted ray and the slope of the outside surface as extended from the previous point  $E_{i-1}$ . The tangent to the surface which will refract the ray in the desired direction, U, may again be found from Snell's law. The accuracy of the calculated surface increases as the spacing between points decreases or as angle  $\Delta U''$  becomes very small. The final ray is parallel to the lens axis and point  $D_n$  corresponds to point B, which is the origin. The area

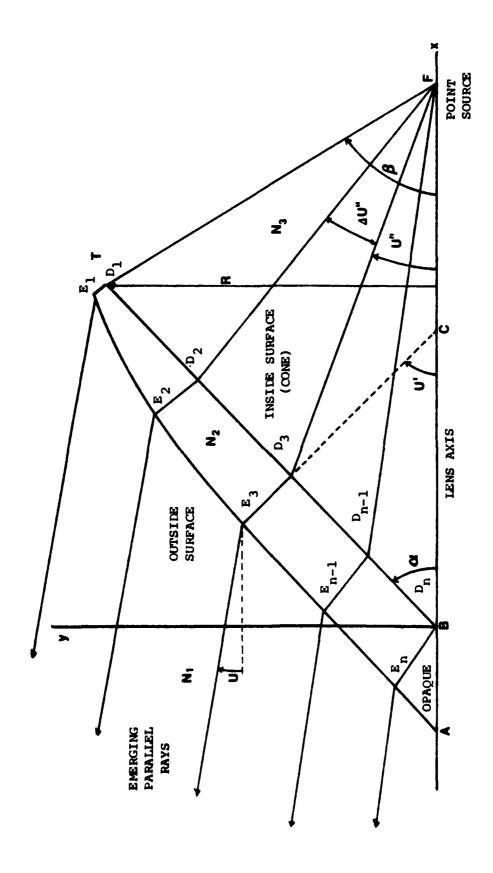


Figure 2. Conceptual Drawing and Coordinate System of Outer Surface Solution

formed by the triangle  $\mathtt{ABE}_n$  is opaque to prevent rays from scattering through the opposite side of the lens.

A particular lens is begun by specifying the focal length BF, the cone half-angle  $\alpha$ , the indices of refraction  $N_1$ ,  $N_2$ , and  $N_3$ , the cone radius at the edge R, the ray direction angle U, and the total number of rays to be traced. The thickness, T, of the lens at the edge  $(E_1D_1)$  must also be specified in order to define which of the family of possible outer surfaces will be calculated. In general, angle U will be taken to be zero in order to investigate objects on-axis at infinity. R will always be set equal to one and to nondimensionalize all linear dimensions every length is implicitly expressed as a ratio with respect to R. The angle  $\beta$  at the focal point is measured from the lens axis to  $D_1$  and is expressed as

$$\tan |\beta| = \frac{R}{(BF - R \cdot \cot a)}$$
 (2)

Thus, if I+l is the total number of rays to be traced, then

$$\Delta U'' = \beta/I \tag{3}$$

and

$$U'' = \beta - J \cdot \Delta U'' \tag{4}$$

where J is the ray number. To begin the lens design the

coordiantes  $x_2(J)$ ,  $y_2(J)$  of point  $D_J$  must be found. Here the subscript 2 refers to the inside surface where  $x_1(J)$ ,  $y_1(J)$  are the coordinates of  $E_J$  on the outside surface.

Following Kingslake [7], Q in Figure 3 may be expressed as

$$Q = BF \sin U'' = x_2(J) \sin U'' + y_2(J) \cos U''$$
 (5)

But,

$$y_2(J) \cdot = x_2(J) \tan \alpha$$
 (6)

so that

BF sin U" = 
$$x_2(J)$$
 [sin U" + tan  $\alpha$  cos U"] (7)

Now,  $x_2(J)$  and  $y_2(J)$  are

$$x_2(J) = \frac{BF \sin U''}{[\sin U'' + \tan \alpha \cos U'']}$$
 (8)

$$y_2(J) = \frac{BF \text{ in } U'' \tan \alpha}{[\sin U'' + \tan \alpha \cos U'']}$$
 (9)

Looking now to find the incident angle  $I_2^1$  which the ray forms with the inside surface at  $D_{,T}$  it can be seen that

$$\tan (\alpha + I_2') = \frac{BF - x_2(J)}{y_2(J)}$$
 (10)

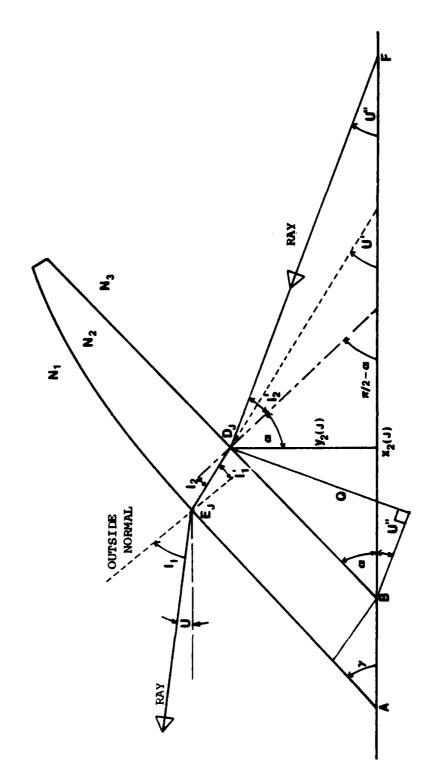


Figure 3. Ray Trace Geometry

Substituting into Equation (1), Snell's law, the refraction angle  ${\bf I}_2$  can now be written as

$$I_2 = \sin^{-1}[(\frac{N_3}{N_2}) \sin I_2']$$
 (12)

The path of the ray inside the lens may now be expressed in the form of the angle U' where

$$U^* = (\frac{\pi}{2} - \alpha) - I_2$$
 (13)

To this stage all rays are treated the same. Now, however, a differentiation must be made between the outermost ray which defines the edge of the lens and all other subsequent rays.

For the outermost ray the thickness, T, of the lens at the edge can be specified which yields  $\mathbf{x}_1(1)$  and  $\mathbf{y}_1(1)$  immediately:

$$x_1(1) = x_2(1) - T \cos U'$$
 (14)

$$y_1(1) = y_2(1) + T \cos U'$$
 (15)

At point  $\mathbf{E}_1$  we may write Equation (1) again to obtain

$$I_1' = \sin^{-1}\left[\left(\frac{N_1}{N_2}\right) \sin I_1\right]$$
 (16)

where angles  $I_1$  and  $I_1'$  are the incident and refracted angles at the outside surface and are not yet known. Angles  $I_1$  and  $I_1'$  may be found by first noting that

$$u + i_1 = u' + i_1'$$
 (17)

then substituting Equation (16) into (17) and rearranging terms to obtain

$$I_1 = U' - U + \sin^{-1}\left[\left(\frac{N_1}{N_2}\right) \sin I_1\right]$$
 (18)

It can be shown after some algebraic manipulation that Equation (18) may be solved for the angle  $I_1$  in the form

$$I_{1} = \sin^{-1} \left\{ \frac{\sin^{2}(U'-U)}{\left[\cos(U'-U) - N_{1}/N_{2}\right]^{2} + \sin^{2}(U'-U)} \right\}^{1/2}$$
(19)

Angle  $I_1^*$  may be found, if desired, by substitution of Equation (19) into Equation (16). Of more importance, however, is the determination of the slope of the outer surface and the surface normal. Now that angle  $I_1$  is known these slopes may be written as

$$DYDXN(J) = \frac{dy}{dx} \Big|_{normal} = - tan(U + I_1)$$
 (20)

and

DYDXT(J) = 
$$\frac{dy}{dx} \Big|_{\text{surface}} = \cot u(U + L_1)$$
 (21)

Equations (20) and (21) may be evaluated by using Equation (19) for  $I_7$ .

Now that all the parameters are known for the first ray, the remainder of the points  $E_J$  and the respective slopes may be found. Each successive ray is traced as before by Equations (4) through (13). Equations (14) and (15), however, may not now be used since the lens thickness along the ray can not be specified. Instead, the intersection of the ray and the surface slope from the previous ray is used to define the new point  $E_{J+1}$ . The intersection is found by first writing the equations of lines representing the ray and the surface tangent. For the ray:

$$y = -x \tan U' + y_2(J) + x_2(J) \tan(U')$$
 (22)

For the surface tangent:

$$y = x \cot (I_1 + U) + y_1(J-1) - x_1(J-1) \cot (I_1 + U)$$
(23)

Equations (22) and (23) are solved simultaneously to yield the coordinates of  $\mathbf{E}_{\mathsf{T}}$  which are

$$y_1(J) = \left\{\frac{A + B}{C}\right\} \tag{24}$$

where:

$$A = \cot(I_1 + U) \cot(U') [y_2(J) + x_2(J) \tan(U')]$$
 (25)

$$B = y_1(J-1) - x_1(J-1) \cot(I_1 + U)$$
 (26)

$$C = 1 + \cot(I_1 + U) \cot(U')$$
 (27)

and

$$x_1(J) = \cot(U')[-y_1(J)+y_2(J)+x_2(J)\tan(U')]$$
 (28)

Now Equations (19), (20), and (21) are used to calculate  $I_1$  and the slopes at  $E_J$ . Therefore, the remainder of the lens surface may be generated. The opaque region at the surface is formed by extending the slope of the surface at  $E_K$  to intersect the lens axis. Here, K refers to the last ray. The nose half-angle,  $\gamma$ , thus formed is given by

$$\tan \gamma = \frac{dy}{dx} | \text{surface, } E_K$$
 (29)

### D. SKEW RAYS

A skew ray is one that begins from an off-axis object point and enters the lens either in front of or behind the meridian plane (z = 0). For every skew ray, there is a corresponding mirror image skew ray on the opposite side of the meridian plane so that two skew rays are traced at the

expense of only one calculation. These two skew rays intersect at the same diapoint.

Large numbers of skew rays are traced through a lens in order to study lens performance at different obliquities. The procedure is to superimpose a grid over the lens aperture and to trace rays through the intersections of the grid, through the lens and onto the image plane. An image plane spot diagram and an energy density plot are then constructed for study.

To accomplish this, the aperture grid has been attached to the nose (opaque region) of the seeker lens at station A; see Figure 3. The plane of the aperture grid is tilted relative to the lens by a variable angle. A transformation between the grid coordinates and the lens coordinates has been derived to connect skew rays from grid to intercept with the outside surface of the lens. Referring to Figure 4, it can easily be seen that

$$x' = [x + AB]\cos \alpha_p - y \sin \alpha_p$$
 (30)

$$y' = [x + AB] \sin \alpha_p + y \cos \alpha_p$$
 (31)

$$z' = z \tag{32}$$

where  $\alpha_p$  is the tilt of the grid plane, AB is the length of the opaque nose portion on the x-axis. Both z and z' are positive out of the page.

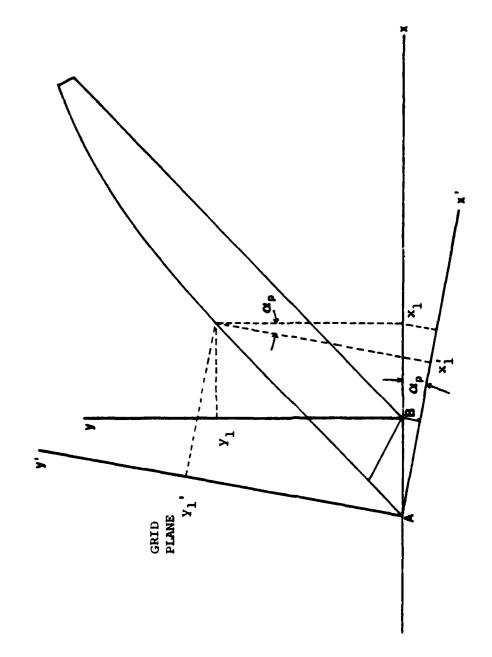


Figure 4. Coordinate Transformation (z and z' are out of the page)

The three dimensional lens outer surface is generated by rotating the array of outer surface coordinates in the meridian plane through  $2\pi$  about the lens axis. Each pair of coordinates  $x_1(J), y_1(J)$  thus describe a circle in the y-z plane of the lens. This circle transforms, however, to an ellipse in the grid plane given by the equation

$$y' = [x + AB] \sin \alpha_p \pm \sqrt{R^2 - z'^2} \cos \alpha_p \qquad (33)$$

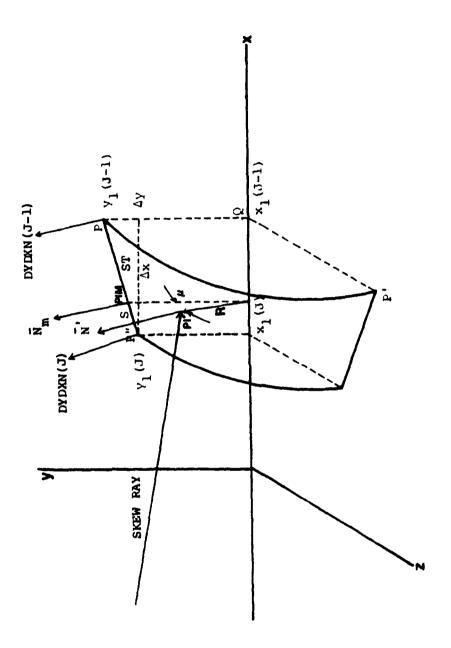
where R is the radius of the circle in the x-y plane. R may be expressed by the familiar equation of a circle

$$R^2 = y^2 + z^2 \tag{34}$$

in the y-z plane of the lens, or by solving Equation (35) for R in the grid plane

$$R^2 = z'^2 + \left(\frac{y'}{\cos \alpha_p} - (x + AB) \tan \alpha_p\right)^2$$
 (35)

With the aid of the foregoing groundwork, the x,y,z coordinates of the ray intercept with the outer surface may be found. To see how this is accomplished, first refer to Figure 5. The skew ray will pass outside the circle formed by rotating some point  $x_1(J)$ ,  $y_1(J)$  and inside the next circle formed by rotating  $x_1(J-1)$ ,  $y_1(J-1)$ . In so doing the skew ray will



Geometry of Skew Ray Intercept With Outside Surface. X, Y Plane is the Meridian Plane.  $\tilde{N}_m$  is the Surface Normal in the Meridian Plane;  $\tilde{N}$ ' the Surface Normal at the Point of Intercept, PI. Figure 5.

intercept an imaginary cone formed by the two circles. Now, the y',z' of the skew ray in the grid plane and a trial value  $\mathbf{x}_1$  (J) from the array of surface points are substituted into Equation (35) and a value for R is calculated. If R is greater than  $\mathbf{y}_1$  (J) but less than  $\mathbf{y}_1$  (J-1) than the appropriate circles have been found. If not, another trial value of  $\mathbf{x}_1$  (J) is picked. In practice the  $\mathbf{x}_1$  (J) corresponding to

$$y_1(J) = z' \qquad (36)$$

is chosen as the first trial value and subsequent trial values are  $x_1$  (J-1),  $x_1$  (J-2) and so forth.

The equation of the cone passing through the two circles may be derived by again referring to Figure 5 and noting that

$$QP = QP' = \sqrt{y^2 + z^2} = R$$
 (37a)

$$QP = y = m[x - x_1(J)] + y_1(J)$$
 (37b)

where the slope m is given by

 $<sup>^{\</sup>rm l}{\rm Recall}$  that a straight line between points  ${\rm E_J}$  was used to form the outer surface in the meridian plane. These approximations are valid only if the spacing between points or circles is very small.

$$m = \frac{y_1(J-1) - y_1(J)}{x_1(J-1) - x_1(J)}$$
 (38)

and therefore

$$R^{2} = \left\{ \left[ \frac{y_{1}(J-1) - y_{1}(J)}{x_{1}(J-1) - x_{1}(J)} \right] \left[ x - x(j) \right] + y_{1}(J) \right\}^{2}$$
 (39)

By equating Equation (39) to Equation (35) and solving for x, the expression for the x-coordinate,  $x_0$ , of the outside surface ray intercept may be found. After some algebra, this reduces to the complicated relationship

$$x_0 = \frac{PAR1}{PAR2} \pm \sqrt{\frac{PAR1}{PAR2} - \frac{PAR3}{PAR2}}$$
 (40)

where:

$$PARl = ab + ce - b^2AB - c^2d$$
 (41)

$$PAR2 = b^2 - c^2$$
 (42)

PAR3 = 
$$a^2 + b^2 (AB)^2 - 2abAB + z' - (e-cd)^2$$
 (43)

and where

$$a = y'/\cos \alpha_p \tag{44}$$

$$b = \tan \alpha_{p} \tag{45}$$

$$c = \left[\frac{y_1(J-1) - y_1(J)}{x_1(J-1) - x_1(J)}\right] \tag{46}$$

$$d = x_1(J) \tag{47}$$

$$e = y_1(J) \tag{48}$$

It has been found that the plus sign in Equation (40) gives the correct values until PAR1/PAR2 becomes greater than  $\mathbf{x}_1(1)$  at which time the negative sign must be used. It now follows from Equations (31), (32) and (40) that the  $\mathbf{y}_0$  and  $\mathbf{z}_0$  coordinates of the intercept are

$$y_0 = \frac{y'}{\cos \alpha_p} - (x_0 + AB) \tan \alpha_p$$
 (49)

and

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$$\mathbf{z}_0 = \mathbf{z}' \tag{50}$$

The next step in tracing the skew ray is to ascertain the direction cosines of the ray inside the lens after refraction at the outside surface. This is accomplished in a series of steps beginning with the determination of the angle  $\phi$  which the skew ray makes with the normal to the surface at the point of intercept. Angle  $\phi$  may be found by taking the scalar product of the direction cosines of the

skew ray and the surface normal. Thus to find  $\phi$ , the surface normal must first be obtained.

As might be surmised by the reader, the direction cosines of the surface normal can be found by taking the gradient of Equation (39) with Equation (37) substituted for R. Several orders of magnitude may be gained in accuracy, however, if the normal  $\overline{N}_{m}$  is found by interpolation in the meridian plane and then rotated to the point of intercept by angle  $\mu$  as in Figure 5. This interpolation procedure involves multiplying the ratio S/ST with the difference between DYDXN(J) and DYDXN(J-1) and adding the product to DYDXN(J).

$$DYDXN_{PIM} = \frac{S}{ST}(DYDXN(J-1) - DYDXN(J)) + DYDXN(J)$$
 (51)

where subscript PIM refers to point of intersection, meridian. Values of the slopes of the normals in the meridian plane, DYDXN(J), are given by Equation (20) and S and ST are the linear separation of points P", PIM and P", P respectively. S/ST is given by

$$\frac{S}{ST} = \frac{\left[ (R - y_1(J))^2 + (x_{pim} - x_1(J))^2 \right]^{1/2}}{(\Delta x^2 + \Delta y^2)^{1/2}}$$
(52)

Therefore, the normal vector,  $\overline{\mathbf{N}}_{\mathbf{m}},$  in the meridian plane is

$$\overline{N}_{m} = (\frac{1}{DYDXN_{PIM}})\hat{i} + \hat{j}$$
 (53)

and since

$$\tan \mu = \frac{z'}{Y_0} = \frac{z_0}{Y_0}$$
 (54)

the normal vector at the point of intersection on the outside surface becomes

$$\overline{N}' = (\frac{1}{\overline{DYDXN_{PIM}}})\hat{i} + \cos(\tan^{-1}\frac{z_0}{y_0})\hat{j} + \sin(\tan^{-1}\frac{z_0}{y_0})\hat{k}$$
 (55)

and

$$|\overline{N}'| = [DYDXN_{PIM}^{-2} + 1]^{1/2}$$
(56)

Now the scalar product of the skew ray vector,  $\hat{R}$ , and surface normal,  $\overline{N}$ , may be performed to find  $\phi$ . Since the acute angle between these vectors is required, the dot product must be written

$$\hat{\mathbf{R}} \cdot \overline{\mathbf{N}}' = |\overline{\mathbf{N}}'| \cos(\pi - \phi) \tag{57}$$

Substituting Equation (55) for  $\overline{N}^{\,\prime}$  and noting that

$$\hat{R} = K\hat{i} + L\hat{j} + M\hat{k} = \cos \alpha_{p}\hat{i} - \sin \alpha_{p}\hat{j}$$
 (58)

the left side of Equation (56) may be expanded to

$$\hat{R} \cdot \overline{N}' = \frac{\cos \alpha_p}{DYDXN_{PIM}} - \sin \alpha_p \cos(\tan^{-1} \frac{z_0}{y_0})$$
 (59)

Solving Equation (56) for the angle  $\phi$  and introducing Equation (58), the expression for angle  $\phi$  becomes

$$\phi = \pi - \cos^{-1} \left\{ \frac{\frac{\cos \alpha_{p}}{DYDXN_{PIM}} - \sin \alpha_{p} \cos(\tan^{-1} \frac{z_{0}}{y_{0}})}{\left[DYDXN_{PIM}^{-2} + \cos^{2}(\tan^{-1} \frac{z_{0}}{y_{0}}) + \sin^{-2}(\tan^{-1} \frac{z_{0}}{y_{0}})\right]^{1/2}} \right\}$$
(60)

Now that the acute angle,  $\phi$ , between the ray and the outside surface normal is known,  $\phi$ , the angle between the refracted ray inside the lens and the surface normal may be found using Snell's law.

$$\phi' = \sin^{-1}(\frac{N_1}{N_2}\sin\phi) \tag{61}$$

The next step is to ascertain the direction cosines K', L', M' of the refracted ray inside the lens. Following Kingslake [7], Figure 6 shows the optical vector relationship between  $\overline{R}$ , the skew ray,  $\overline{R}$ ', the refracted skew ray,  $\overline{N}$ ', the surface normal, and the indices of refraction of the two media. Algebraically this relationship is written as

$$N_2 \overline{R}' = N_1 \overline{R} + (N_2 \cos \phi' - N_1 \cos \phi) \overline{N}'$$
 (62)

By resolving Equation (61) into component form, the direction cosines of  $\overline{R}$ ' may be found. After rearranging, K', L', and M' become

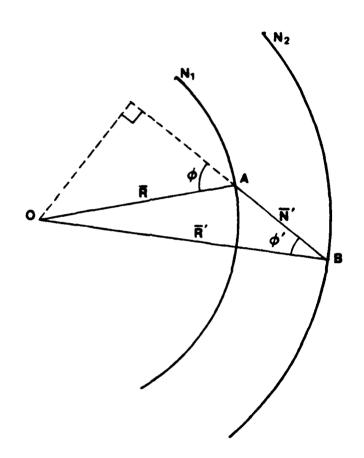


Figure 6. Kingslake's Skew Ray Diffraction Diagram

$$K' = \frac{N_1}{N_2} K + (\cos \phi' - \frac{N_1}{N_2} \cos \phi) k$$
 (63)

$$L' = \frac{N_1}{N_2} L + (\cos \phi' - \frac{N_1}{N_2} \cos \phi) \ell$$
 (64)

$$M' = \frac{N_1}{N_2} M + (\cos \phi' - \frac{N_1}{N_2} \cos \phi) m$$
 (65)

where K, L, and M are given by Equation (57) and k,  $\ell$ , m are

$$k = \frac{DYDXN_{PIM}^{-1}}{|\overline{N}'|}$$
 (66)

$$\ell = \frac{\cos(\tan^{-1}\frac{z_0}{y_0})}{|\overline{N}'|}$$
(67)

$$m = \frac{\sin(\tan^{-1}\frac{z_0}{y_0})}{|\overline{N}'|}$$
(68)

Again,  $|\overline{N}'|$  in Equations (66) through (68) is expressed by Equation (56).

At this stage in the process of tracing the skew ray, the coordinates of the external ray intercept  $\mathbf{x}_0$ ,  $\mathbf{y}_0$ ,  $\mathbf{z}_0$  are known as are the direction cosines K', L', M' of the skew ray inside the lens. The next step is, of course, to find the ray intercept with the inside surface (cone) and the direction of the ray subsequent to refraction. The intermediate steps are similar to, if not identical with,

the foregoing. One major difference, however, is that the inside conical surface may be expressed analytically and the intercept coordinates  $x_i$ ,  $y_i$ ,  $z_i$  may be found exactly without approximation. Proceeding, it can be seen that if  $x_i$ ,  $y_i$ ,  $z_i$  were already known, Equations (63), (64), and (65) could be rewritten as

$$K' = \frac{x_i - x_0}{D} \tag{69}$$

$$L' = \frac{y_i - y_0}{D} \tag{70}$$

$$M' = \frac{z_i - z_0}{D} \tag{71}$$

where D is the distance between surface intercept points.
Upon rearranging:

$$x_i = DK' + x_0 \tag{72}$$

$$y_i = DL' + y_0 \tag{73}$$

$$z_i = DM' + z_0 \tag{74}$$

Furthermore, there exists a relationship between  $x_i$ ,  $y_i$ , and  $z_i$  which is given by the expression for the inside conical surface:

$$y_i^2 + z_i^2 - x_i^2 \tan^2 \alpha = 0$$
 (75)

where  $\alpha$  is the half-angle of the cone. Substituting Equations (72), (73), and (74) into Equation (75) and solving for D, it is seen that

$$D = \frac{-P_1}{P_2} \pm \frac{\sqrt{P_1^2 - P_2 P_3}}{P_2}$$
 (76)

where

$$P_1 = L'y_0 + M'z_0 - K'x_0 tan^2 \alpha$$
 (77)

$$P_2 = L^{2} + M^{2} - K^{2} \tan^2 \alpha$$
 (78)

and

$$P_3 = y_0^2 + z_0^2 - x_0^2 \tan^2 \alpha \tag{79}$$

Correct values for D are obtained by using the minus sign in Equation (76). Now that D is known, values for K', L', M',  $x_0$ ,  $y_0$ , and  $z_0$  are substituted into Equations (72), (73), and (74) to yield the coordinates of the inside surface intercept point.

In the case of the inside surface, the gradient may be used to obtain the surface normal at the point of intercept since this surface has been expressed analytically. If Equation (75) is denoted by f(x,y,z) then

$$\hat{N}_{i} = \frac{|\nabla f|}{|\nabla f|} |x_{i}, y_{i}, z_{i}$$
 (80)

where  $\hat{N}_{i}$  is the unit vector in the direction of the surface normal at the point of intercept. Here,

$$\vec{\nabla} f = \frac{\partial f}{\partial x} \hat{i} + \frac{\partial f}{\partial y} \hat{j} + \frac{\partial f}{\partial z} \hat{k}$$
 (81)

or alternately,

$$\overline{\nabla} \mathbf{f} = \mathbf{f}_{\mathbf{x}} \hat{\mathbf{i}} + \mathbf{f}_{\mathbf{y}} \hat{\mathbf{j}} + \mathbf{f}_{\mathbf{z}} \hat{\mathbf{k}}$$
 (82)

Thus, the direction cosines of the normal k',  $\ell'$ , m' may be expressed as

$$k' = \frac{f_x}{|\overline{\nabla}f|}$$
 (83)

$$\ell' = \frac{f_{y}}{|\overline{\nabla}f|}$$
 (84)

$$\mathbf{m'} = \frac{\mathbf{f_z}}{|\overline{\nabla}\mathbf{f}|} \tag{85}$$

where:

$$f_{x} = -2x_{i} \tan^{2}\alpha \qquad (86)$$

$$f_{y} = 2y_{i} \tag{87}$$

$$f_z = 2z_i \tag{88}$$

and

$$|\overline{\nabla}f| = 2[x_i^2 \tan^4 \alpha + y_i^2 + z_i^2]^{1/2}$$
 (89)

Now that the two unit vectors  $\hat{R}'$  and  $\hat{N}'$  are known the incident acute angle of the skew ray with the inside surface,  $\phi_i$ , may be ascertained by again taking the scalar product as in Equation (57). Thus,

$$\phi_i = \pi - \cos^{-1} [K'k' + L'l' + M'm']$$
 (90)

and using Equation (1) again

$$\phi_{i}^{*} = \sin^{-1}\left[\frac{N_{2}}{N_{3}}\sin\phi_{i}\right] \tag{91}$$

where  $\phi_{1}^{*}$  is the acute angle between the skew ray and the surface normal after refraction at the inside surface. Direction cosines K", L", M" of the skew ray after refraction are found analogously to K', L', and M' in Equations (63), (64) and (65). Here, however, N<sub>2</sub>, N<sub>3</sub>, K', L', M',  $\phi_{1}$ ,  $\phi_{1}^{*}$ , and k',  $\lambda$ ', m' are substituted for N<sub>1</sub>, N<sub>2</sub>, K, L, M,  $\phi$ ,  $\phi$ ', and k,  $\ell$ , m respectively. Thus

$$K'' = \frac{N_2}{N_3} K' + (\cos \phi_i' - \frac{N_2}{N_3} \cos \phi_i) k'$$
 (92)

$$L'' = \frac{N_2}{N_3} L' + (\cos \phi_1' - \frac{N_2}{N_3} \cos \phi_1) \ell'$$
 (93)

and

$$M'' = \frac{N_2}{N_3} M' + (\cos \phi_i' - \frac{N_2}{N_3} \cos \phi_i) m'$$
 (94)

Finally, the skew ray intercept with the image plane may be found. The image plane is treated as another surface along the path of the skew ray, and the intercept coordinates  $x_{im}$ ,  $y_{im}$ ,  $z_{im}$  are easy to find. Equations (72), (73) and (74) may be used again in the form

$$x_{im} = D'K'' + x_{i}$$
 (95)

$$y_{im} = D'L'' + y_i \tag{96}$$

$$z_{im} = D'M'' + z_{i}$$
 (97)

where D', in this instance, is the linear separation between coordinates  $x_i$ ,  $y_i$ ,  $z_i$  at the inside surface and  $x_{im}$ ,  $y_{im}$ ,  $z_{im}$  of the image plane. Furthermore,

$$x_{im} = BF (98)$$

is the equation of the image plane. It follows that

$$D' = \frac{(BF - x_i)}{K''}$$
 (99)

after substitution of Equation (98) into Equation (95). Therefore,

$$y_{im} = \left[\frac{BF - x_{i}}{K''}\right]L'' + y_{i}$$
 (100)

and

$$z_{im} = \left[\frac{BF - x_{i}}{K''}\right]M'' + z_{i}$$
 (101)

Hence, the skew ray has been traced onto the image plane at the focal point of the lens. The  $y_{im}$  coordinate of the corresponding mirror image skew ray is the same as Equation (100). The  $z_{im}$  coordinate, however, is the negative of Equation (101) since the mirror image ray is behind (when viewed along the z-axis) the meridian plane. After tracing a complete set of skew rays through the lens a spot diagram may be plotted. Clearly, the number of rays to be traced depends entirely upon the incremental size of the aperture grid chosen. According to Kingslake [7], at least 100 rays must be traced to give a fair approximation of the actual image. In addition to the spot diagram, an energy density plot may now be constructed by counting the number of rays within progressively larger radii from the image centroid and then plotting the number of rays as a function of radius. Here, each ray is assumed to contain a unit, nondimensional, an amount of radiant energy for convenience.

## E. RADIANT ENERGY LOSS

In reality, each ray loses intensity upon transmission at each interface. Of the total amount of energy contained in each ray, a fraction  $\mathbf{I}_{\mathbf{T}}$  will be transmitted, a fraction  $\mathbf{I}_{\mathbf{R}}$  will be reflected, and a fraction  $\mathbf{I}_{\mathbf{A}}$  will be absorbed by the medium into which the ray is propagating. Since it has been assumed that absorption is negligible, it must be true that

$$I_{T} + I_{R} = 1 \tag{102}$$

Furthermore, the relative amounts of transmitted and reflected electromagnetic energy may be calculated by the well-known Fresnel Equations which state the dependency of  $\mathbf{I}_{\mathbf{T}}$  and  $\mathbf{I}_{\mathbf{R}}$  upon the angle of incidence and the indices of refraction at the interface. That  $\mathbf{I}_{\mathbf{T}}$  and  $\mathbf{I}_{\mathbf{R}}$  are further dependent upon the orientation of the electric vector with respect to the geometry of ray incidence is fundamental to the boundary conditions which govern the form of the Fresnel relations as derived in Hecht-Zajac [8]. Since this thesis examines lens response to monochromatic radiation, the indices of refraction are not considered as a function of wavelength; further, the electric vector orientation is assumed to be rapidly and randomly changing with time. By time averaging field components, it may be seen that the reflectance is

$$I_R = \frac{1}{2}(r_{\perp}^2 + r_{\parallel}^2)$$
 (103)

where

$$r_{\perp} = \frac{\cos \phi - (N_{ti}^2 - \sin^2 \phi)^{1/2}}{\cos \phi + (N_{ti}^2 - \sin^2 \phi)^{1/2}}$$
(104)

and

$$r_{\parallel} = \frac{N_{\text{ti}}^{2} \cos \phi - (N_{\text{ti}}^{2} - \sin^{2} \phi_{i})^{1/2}}{N_{\text{ti}}^{2} \cos \phi + (N_{\text{ti}}^{2} - \sin^{2} \phi)^{1/2}}$$
(105)

Here, N<sub>ti</sub> is the ratio of the index of refraction of the transmission side of the interface to the index of refraction of the incident side. From Equation (102) it now follows that the transmittance through the interface is

$$I_{t} = 1 - I_{R} \tag{106}$$

The total transmittance through the lens is simply the product of  $I_{t}$  at the outside surface with that of the inside surface where  $\phi_{i}$  is substituted for  $\phi$  in Equations (104) and (105).

Total internal reflection of the ray may occur at the inside surface if the incident angle becomes too large. Following Reference (7), this occurs when angle  $\phi_{\hat{1}}$  is equal to or greater than  $\pi/2$ . Thus, Snell's law becomes

$$\sin \phi_i = \frac{N_3}{N_2} \tag{107}$$

and any ray with  $\phi_i$  equal to or greater than this will be

totally internally reflected. In this thesis such rays are labeled "failed rays" since they fail to intersect the image plane. If  $N_2$  is 1.5 and  $N_3$  is 1.0, the incident angle for total internal reflection is 41.81°, or greater.

## F. OPTICAL PATH LENGTH (OPL)

The optical path length of a skew ray is an analytical tool with which the researcher may ascertain the phase of a ray at the end of the path. By so doing, the image diffraction pattern may be constructed which shows the addition or subtraction of amplitude depending upon relative phase. Since each ray must begin with the same phase, monochromatic radiation is used for diffraction experiments. OPL is included here only as a matter of interest. The calculation of optical path length is simply the sum of the geometrical path segments of a ray multiplied by the corresponding index of refraction of the medium for that segment.

## III. THE GRIN LENS

### A. THEORY

This thesis assumes a spherically symmetric, inhomogeneous, isentropic medium in which the refractive index varies from point to point but is independent of direction at each point. The refractive index is a function of the coordinates of the points of the region being considered. The problem of describing the resulting curved paths of rays in such GRIN materials has been solved long ago in the form of a single second order vectorial differential equation. Marchand [9] has shown that the solution to the differential equation in the case of spherical gradients can be written in polar coordinates in the plane of the ray as

$$\theta = \theta_0 + e \int_{r_0}^{r} \frac{dr}{r[n^2r^2 - e^2]^{1/2}}$$
 (108)

Here  $r_0$  and  $\theta_0$  are values of r and  $\theta$  at a convenient reference point on the ray and |e| is a scalar constant along the ray given by

$$e = \varepsilon n_0 r_0 \sin \psi_0 \tag{109}$$

Referring to Figure 7, r is measured from the center of symmetry of the index function; angles  $\theta$  and  $\theta_0$  are measured

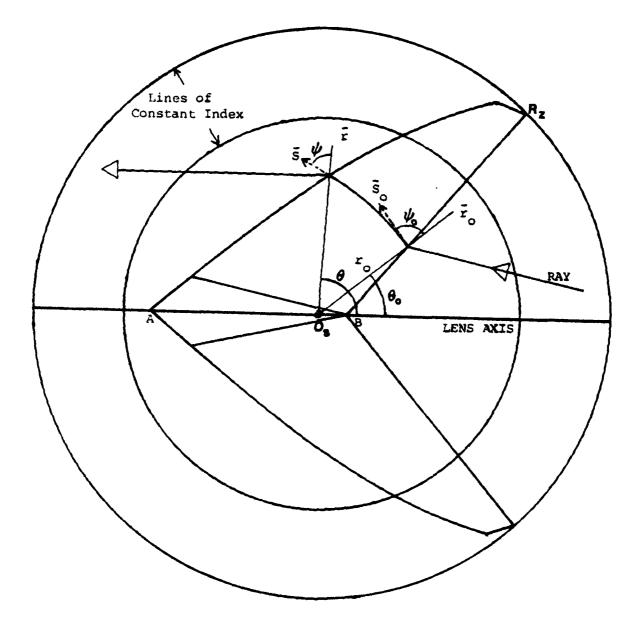


Figure 7. Grin Lens Geometry

counterclockwise from the lens axis; angles  $\psi$  and  $\psi_0$  are measured counterclockwise from the radius vector to the instantaneous ray direction vector.  $\epsilon$  is the sign function expressed as

$$\varepsilon = \operatorname{sgn} |\psi_0| = \pm 1 \tag{110}$$

where

$$|\psi_0| \stackrel{>}{<} \pi/2 \tag{111}$$

In Equation (109),  $n_0$  refers to the index function n evaluated at  $r_0$ . The form of n has been chosen as

$$n = n(O_{S}, r)$$
 (112)

where  $O_S$  denotes the position of the center of symmetry on the lens axis (shown in Figure 7) and r is the radial coordinate from  $O_S$  to the point in question. This form allows the study of the effects on lens performance as  $O_S$  is changed.

More specifically, in order to allow the analytical integration of Equation (108), the expression used for  $n(O_s,r)$  is a generalized version of that employed by Luneburg as described by Marchand [9]. Here

$$n = \left[a + b\left(\frac{r}{R_z}\right)^2\right]^{1/2} \tag{113}$$

whereas Luneburg used specific values for a and b. Generalizing the index function enables the strength, or percent change of the gradient to be varied as well as the algebraic sign of the gradient. If the parameter b is negative, a decreasing parabolic gradient results; conversely a positive b yields an increasing parabolic gradient. By visualizing a plane wave front passing through a spherical gradient, it may be seen that the negative gradient results in light rays bending forward at the center, as shown in Figure 7, whereas a positive gradient has the opposite effect. This principle was used by Wood, as related by Marchand [9], in constructing simple lenses having plane faces and a radial index. The Wood lens acted as a converging or diverging lens depending upon the sign of the gradient used.

Additionally, it must be noted that if the parameter b is equal to zero, Equation (113) reduces to

$$n = \sqrt{a} \tag{114}$$

which is a HIN lens having constant index of refraction.

This fact has facilitated the correlation of GRIN and HIN computer trace algorithms.

Equation (108) may be integrated by a change of variable using the relation

$$v = \left(\frac{r_0}{r}\right)^2 \tag{115}$$

which leads to the solution

$$\theta = \theta_0 - \frac{\varepsilon}{2} \left\{ \sin^{-1} \left[ \frac{2e^2/r^2 - a}{\sqrt{a^2 + 4be^2/R_Z^2}} \right] - \sin^{-1} \left[ \frac{2e^2/r_0^2 - a}{\sqrt{a^2 + 4be^2/R_Z^2}} \right] \right\}$$
(116)

Equation (116) gives  $\theta$  as a function of r. This equation can be easily solved for r as a function of  $\theta$  in the form

$$r = \frac{\sqrt{2} |e|}{\left\{a + \sqrt{a^2 + \frac{4be^2}{R_Z^2}} \sin\left[-2\varepsilon(\theta - \theta_0) + \sin^{-1}\left[\frac{2e^2/r_0^2 - a}{\sqrt{a^2 + 4be^2/R_Z^2}}\right]\right\}\right\}^{1/2}}$$
(117)

The instantaneous direction of the ray at any point  $r,\theta$  may be ascertained by using the invariance of e. Hence,

$$e = \varepsilon nr \sin \psi = \varepsilon n_0 r_0 \sin \psi_0 \qquad (118)$$

and

$$\psi = \sin^{-1} \left[ \frac{n_0 r_0 \sin \psi_0}{nr} \right]$$
 (119)

Furthermore, the orientation of the plane of the ray may be easily deduced since every ray in a spherical medium is a plane curve lying in a plane through the center of symmetry.

Using this fact, Marchand [9] has shown that a suitable conversion from coordinates r and  $\theta$  in the plane of the ray to global Cartesian coordinates may be written in the form

$$x = r \left(\delta \frac{x_0}{r_0} + \eta p_0\right) \tag{120}$$

$$y = r(\delta \frac{y_0}{r_0} + \eta q_0)$$
 (121)

$$z = r(\delta \frac{z_0}{r_0} + \eta \ell_0)$$
 (122)

Here  $p_0$ ,  $q_0$ ,  $\ell_0$  are the initial direction cosines of the ray at  $r_0$ ,  $\theta_0$ ;  $x_0$ ,  $y_0$ , and  $z_0$  are the Cartesian coordinates corresponding to  $r_0$ ,  $\theta_0$ . The parameters  $\delta$  and  $\eta$  are given by Marchand as

$$\eta = \sin \theta / \sin \psi_0 \tag{123}$$

$$\delta = \cos \theta - \eta \cos \psi_0 \tag{124}$$

It should be noted that Equation (117) may become singular for certain rays where  $\psi$  or  $\psi_0$  become very close to zero or  $\pi$ . This singularity may be more easily seen in Equation (123) where  $\eta$  becomes indeterminate as both  $\theta$  and  $\psi_0$  approach zero and/or  $\pi$ . In practice these conditions occur when  $O_S$  is located either far out in object space, coincident with B, or on the image side of the lens; see Figure 8. Positioning

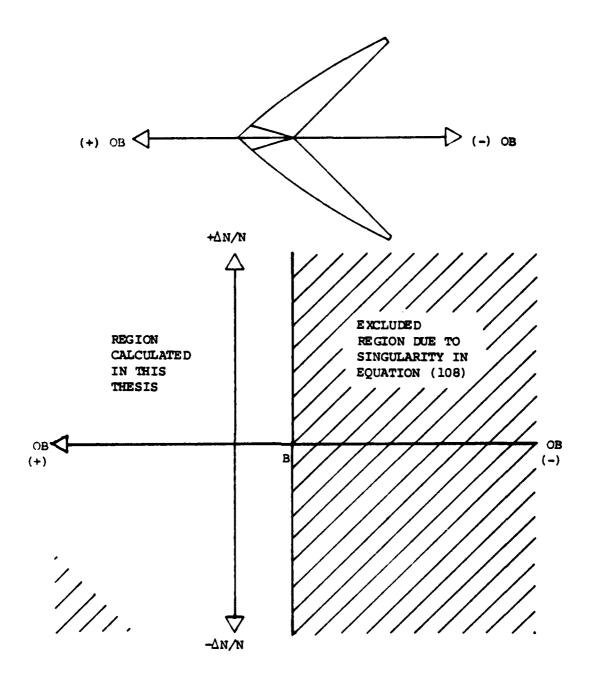


Figure 8. Excluded Regions for Center of Symmetry Due to Singularities

the center of symmetry at these locations has therefore been avoided.

### B. ASSUMPTIONS

The aforementioned assumptions for the HIN case also apply here. Additionally, it must be assumed, as depicted in Figure 7, that the GRIN lens could or soon may be fabricated from a sphere of dielectric material with the required spherically symmetric parabolic gradient.

#### C. GRIN LENS DESIGN PARAMETERS

In the HIN lens, the available design parameters are basic. These include: F, R, T,  $\alpha$ , U, and  $N_2$ . Parameters available for varying the design of the GRIN lens, however, include those of the HIN case but expand the index of refraction variable  $N_2$  into  $O_s$ , a, +b, and -b. These additional lens design parameters greatly expand the lens designer's power to bend radiant energy to his will.

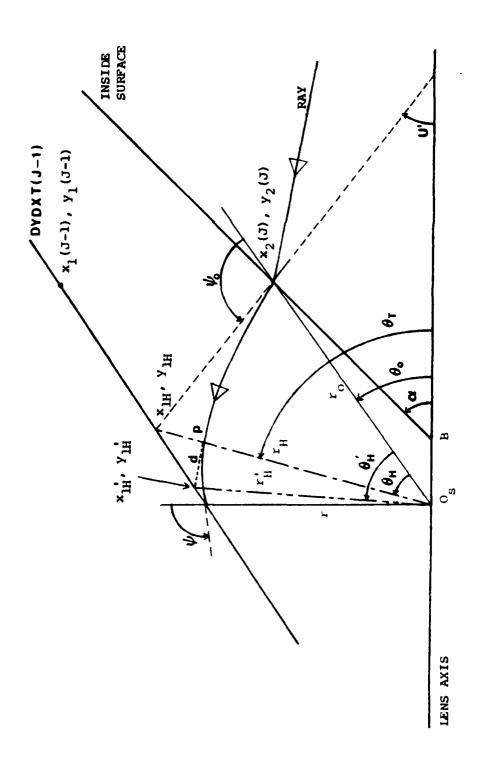
#### D. GRIN LENS DESIGN

The GRIN lens design procedure, although paralleling that of the HIN case, is somewhat more complicated in that the rays are now curved and the index of refraction varies. Accordingly, the same design process is used but with more intermediate calculations required.

The additional calculations arise since the intercept of the GRIN ray with the surface tangent cannot be solved in closed form, and an iterative solution must be used. The reader will note that Equation (117) will yield r if  $(\theta-\theta_0)$  is known. Hence the iterative procedure is to "guess"  $(\theta-\theta_0)$  based on the HIN coordinates  $\mathbf{x}_{1H}$ ,  $\mathbf{y}_{1H}$  which are calculated as in the homogeneous lens. Thus

$$\theta_{H} = \theta_{T} - \theta_{0} \tag{125}$$

in Figure 9 is used in Equation (117) to find point p. tangent to the ray path at point p is now extended to intercept the surface tangent again using the homogeneous intercept relations to find  $x_{1H}^{\prime}, y_{1H}^{\prime}$ . The prime superscripts indicate successive iteration values. The distance d is employed as a measure of the error of point p. If d is not within an acceptable margin then  $\theta_{\,H}^{\,\, t}$  is calculated based on  $\mathbf{x}_{1H}^{\prime},\mathbf{y}_{1H}^{\prime}$  and the procedure repeated to find p' in Figure 9. If d' is not within allowable error then the iteration continues until it is acceptable. In practice, this iteration procedure has proved to be extremely rapid, rarely requiring more than three iterative steps before converging. Slight modifications, however, must be introduced to handle a positive gradient. Furthermore, if Os is located outside the lens proper, a decreasing angular increment must be subtracted from  $\theta_H$  to ensure that each radial vector  $\mathbf{r}_H$ ,  $\mathbf{r}_H^*$ , ... intersects the ray during iteration. These modifications to the intercept iteration procedure are recorded in the program listing for program GISL (for Gradient Index Seeker Lens). Refer to Appendices A, B, and C for a full description of GISL.



D

Figure 9. Grin Ray Intercept Geometry

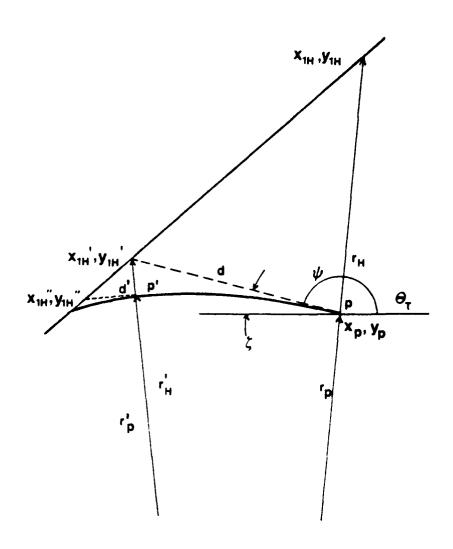


Figure 10. Expanded View, Intercept Geometry

With the overall iterative intercept procedure now being clear, the mathematical details of the GRIN lens design follow. Referring again to Figure 9, the GRIN ray is traced as in the HIN design procedure to obtain the coordinates  $x_2(J), y_2(J)$  on the inside surface in the meridian plane. Now, before Snell's law can be used, the index of refraction must be found at these coordinates. Using Equation (113)

$$n_2 = [a + b(\frac{r_0}{R_z})^2]^{1/2}$$
 (125)

where

$$r_0 = \sqrt{(x_2(J) + OB)^2 + y_2(J)^2}$$
 (126)

and where

$$R_Z = \sqrt{(x_2(1) + OB)^2 + y_2(1)^2}$$
 (127)

Furthermore, by inspection it can be easily seen that

$$\psi_0 = \pi - U' - \theta_0 \tag{128}$$

U' is identical to the HIN case and

Here Equation (125) must be substituted for  $N_2$  in Lation (12).

$$\theta_0 = \tan^{-1} \left[ \frac{y_2(J)}{y_2(J) + OB} \right]$$
 (129)

In Equations (126) through (129) OB is the line segment from  $O_S$  to B and is defined as a positive quantity to the left of B for algebraic ease of manipulation. By substitution into Equation (109) the scalar constant, e, of the ray can now be found. Furthermore, since angle U' is now known coordinates  $\mathbf{x}_{1H}$ ,  $\mathbf{y}_{1H}$  may be calculated using Equations (14) and (15) or Equations (24) through (28).

For the first ray defining the edge of the lens, Equations (14) and (15) are employed, and, since no iteration is required, angle  $(\theta_T^{}-\theta_0^{})$  is used in Equation (117) to find r immediately. Here

$$\cos \theta_{T} = \frac{x_{1H} + OB}{r_{H}}$$
 (130)

and

$$r_{\rm H} = \sqrt{(x_{\rm 1H} + OB)^2 + y_{\rm 1H}^2}$$
 (131)

Thus

$$x_1(1) = r \cos \theta_T - OB$$
 (132)

and

$$y_1(1) = r \sin \theta_T \tag{133}$$

are the coordinates of the first point on the outside surface. Since e is known, Equation (119) is employed to find the angle  $\psi$ . Angle  $\zeta$  in Figure 10 is employed to translate  $\psi$  into the ray direction with respect to the lens axis by the relation

$$\zeta = \pi - (\psi + \theta_{\mathbf{m}}) \tag{134}$$

Therefore, by substituting  $\zeta$  for U' in Equations (19), (20), and (21), DYDXN(1) and DYDXT(1) may be found. Here again, N<sub>2</sub> in Equation (19) must be replaced by n<sub>2</sub> as given by Equation (113) evaluated at  $x_1(1), y_1(1)$ .

Now that the first point on the outside surface is known, along with the surface tangent, the remainder of the K number of rays may be processed to yield the balance of the outside surface. Each subsequent intercept between ray and tangent must be iterated. Thus, unlike the first ray, once  $\mathbf{x}_{1H}$ ,  $\mathbf{y}_{1H}$ ,  $\mathbf{r}_{H}$ ,  $\mathbf{r}_{p}$ ,  $\psi$ , and  $\zeta$  are known,  $\mathbf{x}_{1H}$  and  $\mathbf{y}_{1H}$  are ascertained by the substitution of the coordinates of point p for  $\mathbf{x}_{2}(\mathbf{J})$ ,  $\mathbf{y}_{2}(\mathbf{J})$  in Equations (24) through (28) with U' replaced by  $\zeta$ . The coordinates of point p are

$$x_{p} = r_{p} \cos \theta_{T} - OB \tag{135}$$

$$y_{p} = r_{p} \sin \theta_{T}$$
 (136)

Therefore

$$y_{1H}' = \left\{ \frac{A + B}{C} \right\} \tag{137}$$

where

$$A = \cot(I_1 + U) \cot \zeta [y_p + x_p \tan \zeta]$$
 (138)

$$B = y_1(J-1) - x_1(J-1) \cot[I_1 + U]$$
 (26)

$$C = 1 + \cot(I_1 + U) \cot \zeta$$
 (139)

and

$$I_{1} = \sin^{-1} \left\{ \frac{\sin^{2}(\zeta - U)}{\left[\cos(\zeta - U) - N_{1}/n_{2}\right]^{2} + \sin^{2}(\zeta - U)} \right\}$$
(140)

Note that  $n_2$  in Equation (140) is found from Equation (113) evaluated at  $r_p$ .

Additionally,

$$x_{lh}' = \cot \zeta [-y_{lh}' + y_p + x_p \tan \zeta]$$
 (141)

Now the error, d, may be evaluated as

$$d = \sqrt{(x_p - x_{1H}')^2 + (y_{1H}' - y_p)^2}$$
 (142)

Here d is compared to  $1 \times 10^{-5}$ . If d is larger than this value then the entire procedure is repeated by substituting  $x_{lH}', y_{lH}'$  for  $x_{lH}, y_{lH}$  and so forth. Once the error criteria are satisfied

$$x_1(J) = x_p \tag{143}$$

$$y_1(J) = y_p \tag{144}$$

and the next ray is processed. Correlation between GRIN (with b set equal to zero) and HIN design procedures run with identical parameters has shown agreement to the fifth and sixth decimal places.

### E. SKEW RAYS IN GRIN

GRIN skew rays are handled analogously to the homogeneous case with the same coordinate transformation from grid plane to global coordinates being required. It is only after the initial directions cosines of the GRIN skew ray K', L', M' are found that the differences between GRIN and HIN appear. The only exception to this being the use of Equation (113) in Snell's law for refraction at the interface. Since the GRIN skew rays display curvature in a plane through O<sub>S</sub>, K', L', and M' are constantly changing until intercept with the inside conical surface. Therefore, not only must the plane of the skew ray be analytically described, but the final values of K', L', M' must be found. Due to the nature of

GRIN rays, the procedure for finding the ray intercept with the conical inside surface is different from both the HIN case and the iteration procedure employed in the meridian plane due to the multiplicity of the geometry encountered. The Newton-Raphson iteration routine has been found to be ideal for this purpose.

To begin, the magnitude and direction of the initial radius vector  $\overline{r}_0$  from  $o_s$  for  $x_0$ ,  $y_0$ ,  $z_0$  must be ascertained. The magnitude is given by

$$r_0 = [(x_0 + OB)^2 + y_0^2 + z_0^2]^{1/2}$$
 (145)

Therefore the unit vector in the direction of the intercept is

$$\hat{r}_0 = \frac{x_0 + OB}{r_0} \hat{i} + \frac{y_0}{r_0} \hat{j} + \frac{z_0}{r_0} \hat{k}$$
 (146)

The plane of the ray may be fully described by the vector normal to the plane. Two vectors,  $\hat{\mathbf{r}}_0$  and  $\hat{\mathbf{R}}$  lie in the plane of the ray.  $\hat{\mathbf{R}}$  is the unit vector in the initial direction of the ray after refraction and described using direction cosines as

$$\hat{R} = K'\hat{i} + L'\hat{j} + M'\hat{k}$$
 (147)

Thus, the plane of the ray may be described by the cross product

$$\hat{N}_{p_0} = \hat{r}_0 \times \hat{R}$$
 (148)

or

$$\hat{N}_{p_0} = N_{p_0x} \hat{i} + N_{p_0y} \hat{j} + N_{p_0z} \hat{k}$$
 (149)

where

$$N_{p_{0x}} = r_{0y}M' - r_{0z}L'$$
 (150)

$$N_{p_{0Y}} = r_{0z}K' - r_{0x}M'$$
 (151)

and

$$N_{p_{0z}} = r_{0x}^{L'} - r_{0y}^{K'}$$
 (152)

In Equations (150), (151), and (152)  $r_{0x}$ ,  $r_{0y}$ , and  $r_{0z}$  refer to the x, y, and z components of  $\hat{r}_0$  in Equation (146). Furthermore, the angle  $\psi_0$  between  $\hat{r}_0$  and  $\hat{R}$  may be found from the dot product as

$$\psi_0 = \cos^{-1} [r_{0x}K' + r_{0y}L' + r_{0z}M']$$
 (153)

Now the scalar invariant e may be found. Substituting known values into Equation (109):

$$e = \varepsilon [a + b(\frac{r_0}{R_z})^2]^{1/2} r_0 \sin \psi_0$$
 (154)

where  $R_{\rm Z}$  is unchanged from that found during previous calculations for the lens shape by Equation (127).

With the foregoing groundwork established, the intercept of the ray with the inside surface may be calculated. The Newton-Raphson iteration scheme requires the calculation of the radius vector from  $O_{\mathbf{q}}$  to the cone by geometrical methods and the radius to the ray by GRIN theory. The difference between the two radii is then divided by the difference between the derivatives of the two functions. The resulting quantity is subtracted from the trial angle,  $\theta_{\rm p}$ , in the plane of the ray, to give a new trial angle  $\theta_{D}^{\prime}$ . The process is continued until the difference between radii is less than  $1 \times 10^{-5}$ . The first trial angle is measured to a reference HIN intercept as if the material were homogeneous since the actual GRIN ray curves only slightly. The coordinates of this HIN intercept point are designated  $x_{iH}$ ,  $y_{iH}$ ,  $z_{iH}$  and are derived using the HIN equations as before. To obtain the first trial angle, the scalar product between  $r_0$  and r<sub>iH</sub> is used. Here

$$r_{iH} = [(x_{iH} + OB)^2 + y_{iH}^2 + z_{iH}^2]^{1/2}$$
 (155)

and

$$\hat{r}_{iH} = \frac{x_{iH} + OB}{r_{iH}} \hat{i} + \frac{y_{iH}}{r_{iH}} \hat{j} + \frac{z_{iH}}{r_{iH}} \hat{k}$$
 (156)

Thus,

$$\theta_{p} = \cos^{-1}[r_{0x}r_{iHx} + r_{0y}r_{iHy} + r_{0z}r_{iHz}]$$
 (157)

where  $r_{iHx}$ ,  $r_{iHy}$ , and  $r_{iHz}$  are the x, y, and z components of  $r_{iH}$ , respectively.

In Equation (117)  $\theta_p$  is substituted for  $(\theta-\theta_0)$  to yield r as required by the iteration procedure. The geometrical radius,  $r_g$ , is not so easily acquired. First, note that the equation of the plane of the ray inside the lens is given by

$$N_{p0x}(x-x_0) + N_{p0y}(y-y_0) + N_{p0z}(z-z_0) = 0$$
 (158)

Secondly, the equation of the conical surface is given by Equation (75). The combination of the plane of the ray and the cone yield the loci of possible intercept points on the inside surface. In Cartesian coordinates, the sum of Equations (75) and (158) is

$$x(N_{p0x}-x \tan^{2}) + y[N_{p0y}+y] + z[N_{p0z}+z] - N_{p0x}x_{0}$$
$$- N_{p0y}y_{0} - N_{p0z}z_{0} = 0$$
(159)

Equation (159) must be transformed into coordinates r and  $\theta_{p}$ 

in the plane of the ray. The transformation is made possible using Equations (120) through (124). Upon substitution, and after solving for  $r_q$ , Equation (159) becomes

$$r_g = \frac{-B_2}{2A_2} \pm \sqrt{\frac{B_2^2}{4A_2^2} + \frac{C_2}{A_2}}$$
 (160)

where

$$A_2 = B_1^2 + C_1^2 - A_1^2 \tan^2 \alpha$$
 (161)

$$B_2 = B_1 N_{poy} + C_1 N_{poz} + A_1 N_{pox} + 2A_1 OB tan^2 \alpha$$
 (162)

$$C_2 = OBN_{pox} + OB^2 tan^2 \alpha + N_{pox} x_0 + N_{poy} y_0$$

$$+ N_{poz}^{z}$$
 (163)

and

$$A_1 = \frac{\delta}{r_0} (x_0 + OB) + \eta K'$$
 (164)

$$B_1 = \frac{\delta y_0}{r_0} + \eta L' \tag{165}$$

$$C_1 = \frac{\delta z_0}{r_0} + \eta M' \tag{166}$$

Here,  $\delta$  and  $\eta$  are found from

$$\eta = \sin \theta_{p} / \sin \psi_{0}$$
 (167)

and

$$\delta = \cos \theta_{p} - \cos \psi_{0} \tag{168}$$

which follow from Equations (123) and (124). The plus sign in Equation (160) yields the correct values. Now that r of the ray and  $r_g$  of the surface are known, the derivatives of r and  $r_g$  with respect to  $\theta$  at  $\theta_p$  must be found. It can be shown that for the ray

$$\frac{d\mathbf{r}}{d\theta}\Big|_{\theta_{\mathbf{p}}} = \frac{\varepsilon \mathbf{r}^{3}}{2e^{2}} \cos\left[-2\varepsilon \theta_{\mathbf{p}} + \sin^{-1}(\mathbf{A}_{3})\right]$$
 (169)

where

$$A_3 = \frac{2e^2/r_0^2 - a}{\sqrt{a^2 + 4be^2/R_Z^2}}$$
 (170)

The derivative of  $r_g$  is somewhat more complicated. With persistence, however, it can be shown that

$$\frac{dr_g}{d\theta} = -\frac{1}{2A_2^2} [A_2 \frac{dB_2}{d\theta} - B_2 \frac{dA_2}{d\theta}] + \frac{1}{2} [\frac{B_2^2}{4A_2^2} + \frac{C_2}{A_2}]^{-1/2} [C_3]$$
(171)

where  $A_2$ ,  $B_2$ , and  $C_2$  are given by Equations (161) through (168) and

$$C_3 = \left[\frac{1}{2A_2^4}(A_2^2 B_2 \frac{dB_2}{d\theta} - B_2^2 A_2 \frac{dA_2}{d\theta}) - \frac{C_2}{A_2^2} \frac{dA_2}{d\theta}\right]$$
(172)

Therefore, the r and  $\theta_{p}$  to the intercept are found by iteration of revised trial values

$$\theta_{p}^{\prime} = \theta_{p} - \frac{(r - r_{g})}{(\frac{dr}{d\theta} - \frac{dr_{g}}{d\theta})}$$
 (173)

In practice the quotient of differences in Equation (173) is reduced by a factor of 1.3 to slow convergence and provide stability. The number of iterations, however, rarely exceeds five.

Values for r and  $\theta_p$  are now +ransformed into Cartesian coordinates  $\mathbf{x_i}$ ,  $\mathbf{y_i}$ ,  $\mathbf{z_i}$  by substitution into Equations (120) through (124). Next, the values of the direction cosines K', L', and M' at the intercept are needed. The angle  $\psi$  between the radius vector and the tangent to the ray at intercept may first be deduced from the scalar invariant, e. Hence

$$\psi = \sin^{-1}\left(\frac{e}{n_2 r_{PI}}\right) \tag{174}$$

Where e is known,  $n_2$  is evaluated at  $r_{\rm PI}$ ;  $r_{\rm PI}$  is the radius to the intercept as found by the iteration above. Three constraints on the direction cosines may be written. These are:

- 1) The scalar product of the radius vector,  $\mathbf{r}_{\text{pl}}$ , and the instantaneous ray direction vector,  $\hat{\mathbf{R}}$ , at intercept.
- 2) The scalar product of the normal to the plane of the ray  $\hat{N}_{\text{D0}}$  and  $\hat{R}.$

3) The sum of the squares of the direction cosines must sum to unity.

Mathematically, the above constraints are written as

$$\frac{(x_i + OB)}{r_{PI}} K' + \frac{y_i}{r_{PI}} L' + \frac{z_i}{r_{PI}} M' = \cos \psi \qquad (175)$$

$$N_{p0x}K' + N_{p0y}L' + N_{p0z}M' = 0 (176)$$

$$K'^2 + L'^2 + M'^2 = 1$$
 (177)

To find K', L', and M', Equations (175), (176) and (177) are solved simultaneously. It can be shown that the solution leads to

$$M' = \frac{-B_6}{2A_6} \pm \sqrt{\frac{B_6^2}{4A_6^2} - \frac{C_6}{A_6}}$$
 (178)

$$L' = A_5 - B_5 M' \tag{179}$$

$$K' = (1 - L'^2 - M'^2)^{1/2}$$
 (180)

where

$$A_{6} = \left[1 + \left(\frac{N_{p0z}}{N_{p0x}}\right)^{2}\right]B_{5}^{2} - \frac{2N_{p0y}N_{p0z}}{N_{p0x}^{2}}B_{5}$$

$$+ \left(\frac{N_{p0z}}{N_{p0x}}\right)^{2} + 1 \tag{181}$$

$$B_6 = \frac{2 N_{p0y} N_{p0z}}{N_{p0x}} A_5 - \left[ \left[ 1 + \left( \frac{N_{p0z}}{N_{p0x}} \right)^2 \right] A_5 B_5$$
 (182)

$$C_6 = \left[1 + \left(\frac{N_{p0z}}{N_{p0x}}\right)^2\right] A_5^2 - 1 \tag{183}$$

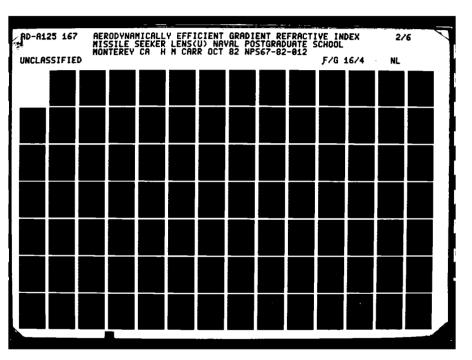
and

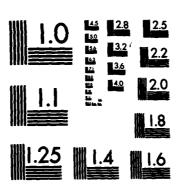
$$A_5 = \frac{N_{p0x} \cos \psi}{N_{p0x} B_4 - N_{p0y} A_4}$$
 (184)

$$B_5 = \frac{N_{p0} x^{C_4} - N_{p0} z^{A_4}}{N_{p0} x^{B_4} - N_{p0} y^{A_4}}$$
 (185)

Furthermore, in Equations (184) and (185),  $A_4$ ,  $B_4$ , and  $C_4$  are the coefficients of K', L' and M' in Equation (175).

Now that the direction cosines of the ray are known at the point of intersection with the inside surface, the index of refraction,  $\mathbf{n}_2$ , is computed by substituting  $\mathbf{r}_{\mathrm{PI}}$  into Equation (113). The remainder of the skew ray trace to the image plane is identical to the homogeneous procedure.





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$$B_6 = \frac{2 N_{p0y} N_{p0z}}{N_{p0x}} A_5 - \left[ \left( 1 + \left( \frac{N_{p0z}}{N_{p0x}} \right)^2 \right] A_5 B_5$$
 (182)

$$C_6 = \left[1 + \left(\frac{N_{p0z}}{N_{p0x}}\right)^2\right]A_5^2 - 1 \tag{183}$$

and

$$A_5 = \frac{N_{p0x} \cos \psi}{N_{p0x} B_4 - N_{p0y} A_4}$$
 (184)

$$B_5 = \frac{N_{p0}x^{C_4} - N_{p0}z^{A_4}}{N_{p0}x^{B_4} - N_{p0}y^{A_4}}$$
 (185)

Furthermore, in Equations (184) and (185),  $A_4$ ,  $B_4$ , and  $C_4$  are the coefficients of K', L' and M' in Equation (175).

Now that the direction cosines of the ray are known at the point of intersection with the inside surface, the index of refraction,  $n_2$ , is computed by substituting  $r_{\rm PI}$  into Equation (113). The remainder of the skew ray trace to the image plane is identical to the homogeneous procedure.

## IV. LENS PERFORMANCE PARAMETERS

The function of the seeker lens is to focus electromagnetic energy either reflected from or emitted by the target onto a detector. Angular displacement of the target with respect to the missile body axes as well as target angular rate information are both desired outputs from the seeker. Hence, it is not only important just to be able to detect the target by focusing energy into a spot on the detector, it is equally important that this spot be as small as possible to enable the precise position of the spot on the detector to be discerned.

The ability of a lens to focus an object to a small spot does not guarantee the quality of the image. For a FLIR optical system extensive effort is expended to obtain an image with minimum aberration [10]. Seeker optics, however, are generally non-imaging devices where the pressure of the different aberrations does not detract from the function of the seeker as long as a tight image is maintained [11].

Accordingly, the most important parameter by which seeker lens performance is judged is that of spot size at different obliquities. Since the image found at the focal point is not necessarily circular nor equally dense, the standard deviation in the y and z directions with respect to image centroid is used to define spot size.

Therefore, once the coordinate pairs of all the skew rays have been calculated in the image plane, the first step in

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the analysis of lens performance is to calculate the image centroid [12]. If the spot diagram is composed of N rays, the centroid location,  $y_c$  is

$$Y_{C} = \frac{1}{N} \sum_{i=1}^{N} Y_{i}$$
 (186)

Since there is symmetry about the x-y plane, the z coordinate of the centroid will always be zero.

The next step in finding the spot size is to find the standard deviations of the spot diagram in the y and z directions. This is accomplished by summing the squares of the differences of the intercept coordinates with respect to the centroid and then dividing by the number of rays. The standard deviations then, are given by

$$\sigma_{z}^{2} = \frac{1}{N} \sum_{i=1}^{N} z_{i}^{2}$$
 (187)

and

$$\sigma_{y}^{2} = \frac{1}{N} \sum_{i=1}^{N} (Y_{i} - Y_{c})^{2}$$
 (188)

The spot size,  $\sigma_r$ , is now defined by

$$\sigma_{\mathbf{r}} = \sqrt{\sigma_{\mathbf{z}}^2 + \sigma_{\mathbf{y}}^2} \tag{189}$$

Of further interest in appraising lens performance is the energy density of the image as a function of radius from the centroid. In nondimensional form, this is simply the number of rays in the spot diagram within a succession of circles of increasing size overlaid about the centroid. Here, each ray is assumed to carry a unit amount of radiant energy. This type of plot facilitates the comparison of different lens designs by detailing the distribution of energy within each image. Clearly, it is desirable to have as much energy as possible concentrated very close to the image centroid. Between two lenses with equal spot sizes, the preferred lens has more energy concentrated within a smaller radius.

Every ray, however, does not deliver an equal amount of energy to the focal plane. It is prudent, therefore, to include as a performance parameter the average ray intensity. Again, for N rays, the intensity I of each ray is summed and normalized by N to yield

$$I_{av} = \frac{1}{N} \sum_{i=1}^{N} I_{i}$$
 (190)

Finally, each lens design is checked for "failed rays".

The reader will recall that these rays fail to intercept the image plane due to total internal reflection, total external reflection, or failure to intercept the inside surface within the bounds of the lens. Hence, a lens design with fewer "failed rays" or no "failed rays" at all is a preferred lens.

Although there are many other performance criteria by which lenses are compared, the foregoing parameters are more than sufficient to judge the merit of preliminary seeker lens designs. It should be noted, however, that notwithstanding the fact that the image centroid and standard deviations were used as stepping stones to obtain image spot size, they have significant meaning of their own. The standard deviations  $\sigma_{_{\boldsymbol{\mathcal{U}}}}$  and  $\sigma_{_{\boldsymbol{\mathcal{Z}}}}$  inform the lens designer as to the horizontal and vertical spread of the image. Image centroid location,  $\mathbf{y}_{\mathbf{c}}$ , at increasing obliquities is of obvious importance since excessive displacement will cause the image to miss the detector entirely and would dictate the necessity for a second lens element to dampen the movement. Furthermore, since Line of Sight (LOS) measurement accuracy to the target is highly dependent upon the linearity of y as a function of the lens tilt angle, lens designs which exhibit a greater degree of such linear behavior are the preferred designs.

To summarize, it is sufficient to note that although spot size is the most important of the performance parameters, every other parameter has a significant impact on the performance of a particular lens.

# V. RESULTS FOR THE HOMOGENEOUS LENS

The performance of the homogeneous lens is presented primarily as a comparison with which to compare the performance of the GRIN lens. Here, the relationship of spot size to increasing lens obliquity and lens thickness are presented as well as the image centroid movement as a function of obliquity. Additionally, the reader is introduced to the four basic computer plots used to display the results: lens shape, object plane with superimposed skew ray grid, image plane or spot diagram, and energy density. These plots were generated on the VERSATEC Plotter using arrays of data points produced by program GISL on the IBM 3033 mainframe computer.

To begin, Figure 11 shows the homogeneous lens shape. The first of the four basic plots, the lens shape plot, presents the lens side view in the meridian plane. The outer surface (curved) and the inner surface (cone) are constructed by connecting the points  $\mathbf{x}_1(\mathbf{J})$ ,  $\mathbf{y}_1(\mathbf{H})$  and  $\mathbf{x}_2(\mathbf{J})$ ,  $\mathbf{y}_2(\mathbf{J})$ , respectively, by straight lines. At the apex of the lens about the lens axis is a trapezoidal region which represents the opaque nose area. In the legend are listed the lens design parameters and the significant calculated dimensions of the lens. Since all linear dimensions are implicitly normalized with respect to the maximum inside radius, R, the lens may be scaled up or down by multiplying each dimension by a factor of  $\mathbf{R}_{\text{new}}/\mathbf{R}$ . Beginning at the top, the parameters listed in the legend are explained in Table 1.

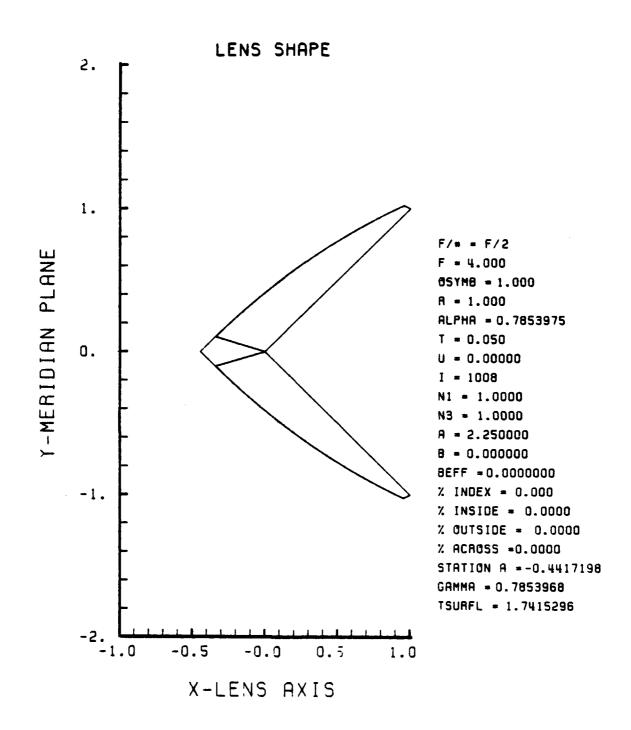


Figure 11. Homogeneous Lens Shape for  $N_2 = 1.5$ 

TABLE 1
Explanation of Lens Shape Plot Legend

PARAMETER	TYPE OF PARAMETER	MEANING
F/#	DESIGN	F number. $F/# = F/2R$
F	DESIGN	Focal length from B.*
OSYMB	DESIGN	OB. Line segment O to B (positive to left).** Immaterial in HIN
R	DESIGN	Maximum radius of cone*
ALPHA	DESIGN	$\alpha$ cone half-angle, * radians
T	Design	Edge thickness*
ט	DESIGN	<pre>Incident ray offset angle (design)*, radians</pre>
I	DESIGN	Number of iterations. I+l = number of rays
N <sub>1</sub>	DESIGN	Free stream index of refraction
N <sub>3</sub>	DESIGN	Index of refraction of interior lens cavity
A	DESIGN	a in $n_2(r) = \sqrt{a + b(r^2/R_Z^2)}$ , gradient refractive index function. $N_2 = \sqrt{a}$ in HIN.
В	DESIGN	b in n <sub>2</sub> (r). Zero in HIN
BEFF	CALCULATED	b effective = $b/R_Z^2$
% Index	DESIGN	Percent change in $n_2(r)$ from $r = 0$ to $r = R_Z$
% Inside	CALCULATED	Percent change in n <sub>2</sub> (r) along inside surface from lens axis to edge
% Outside	CALCULATED	Percent change in n <sub>2</sub> (r) along outside surface from opaque region to edge

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## Table 1 (Continued)

PARAMETER	TYPE OF PARAMETER	MEANING*
% Across	CALCULATED	Percent change in n <sub>2</sub> (r) across lens from lens axis to outside surface at the thickest point
STATION A	CALCULATED	x-coordinate at nose of lens"
GAMMA	CALCULATED	$\gamma$ nose half-angle of opaque nose region, radians
TSURFL	CALCULATED	Total outside surface length from Station A to the edge

<sup>\*</sup>Refer to Figure 2 for clarification

The HIN lens shape has a convex outer surface with maximum thickness on axis of almost ten times the edge thickness. Although the lens has a good aerodynamic shape resembling an ogive, the outer surface is not a circular arc nor can a single analytical function be fitted to the array of points describing the surface. Note that the nose half-angle,  $\gamma$ , is almost identical to the cone half-angle,  $\alpha$ .

All lenses have been designed with a cone half-angle of  $45^{\circ}$  which is approximately the maximum angle for which aerodynamically efficient lens shapes may be designed, considering a free stream Mach Number not to exceed three. Without exception, overall lens performance is more severely degraded as angle  $\alpha$  is reduced.

Table 2 explains the legend of Figure 12, which is the second basic plot. Here the lens is depicted as seen from the skew ray grid plane. The lens tilt angle,  $\alpha_{\rm p}$ , causes the equally spaced (in J) circles descriging the surface of the lens to appear as ellipses. In Figure 12, the grid spacing has been reduced from 0.1, which is normally used, to 0.3 to allow identification of individual rays for correlation with the image plane spot diagram. Although the small number of rays used is not sufficient to given an accurate definition of spot size, the number is sufficient to describe where rays in the object plane are being focused in the image plane by the HIN lens. The skew rays in Figure 12 have been numbered in the order in which they were processed. Actually, only rays 1 through 19 were actually raced; 20 through 33 are

TABLE 2

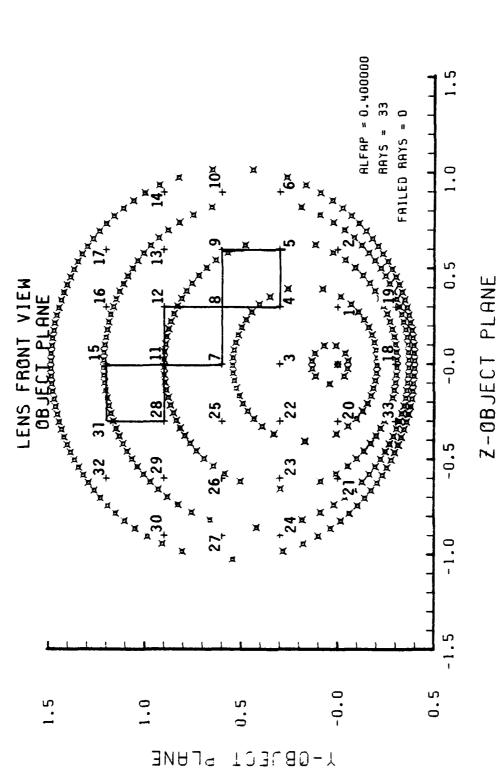
Explanation of OBJECT PLANE Plot Legend

PARAMETER	TYPE OF PARAMETER	MEANING
ALFAP	Analysis	$\alpha_{p}$ . Lens tilt angle, radians
Rays	Analysis	Total number of rays processed
Failed Rays	Analysis	Total number of rays failing to pass through the lens*

Failed Rays are indicated on the plot by a diamond superimposed on the grid location of the ray.

mirror image skew rays. Ray 30 corresponds to ray 14, for example, Here, the lens has been tilted by 0.4 radians, or 22.9 degrees, and 33 rays have been processed of which none have faired to intercept the image plane. The staircase pattern has been added in this case in order to show the resulting distortion present in the image plane (Figure 13).

The Spot Diagram in Figure 13 is an example of the third basic computer plot; see Table 3. Unlike most Spot Diagrams, this example has the individual rays numbered for comparison with Figure 12; also the resulting distorted staircase pattern is sketched. By cross-referencing individual rays between Figures 12 and 13, it is possible to recognize where certain areas of the lens are focusing rays in the image plane. Rays 1, 4, 20, and 22 about the opaque nose region form a coma tail which contributes most of the image spread. Rays 33,



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Figure 12. Example Object Plane

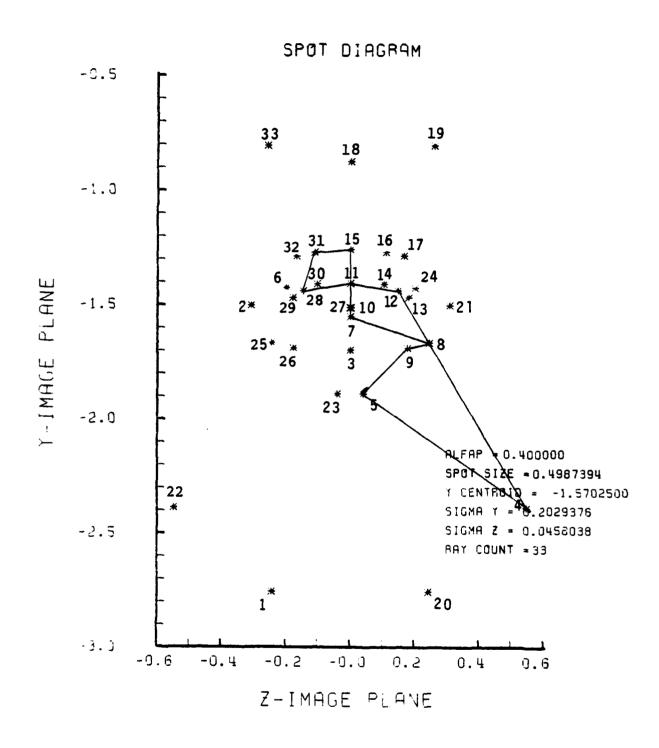


Figure 13. Spot Diagram for HIN Lens Design Shown in Figure 11

TABLE 3

Explanation of Spot Diagram Legend

PARAMETER	TYPE OF PARAMETER	MEANING
ALFAP	Analysis	<pre>a<sub>p</sub>, lens tilt angle, radians</pre>
SPOT SIZE	Performance	$\sigma_{\mathbf{r}}$ , see Equation (189)
Y-CENTROID	Performance	y <sub>C</sub> , y-coordinate of image centroid; see Equation (186)
SIGMA Y	Performance	σ <sub>y</sub> , y-standard deviation; see Equation (188)
SIGMA Z	Performance	$\sigma_z$ , z-standard deviation; see Equation (187)
RAY COUNT	Performance	Number of rays striking the image plane

18, and 19 from the bottom portion of the homogeneous lens are imaged at the top and are widely separated from the core of the image. In general, the regions of the lens which have been found to contribute the bulk of the widely spaced rays are the immediate nose region and the lower portion of the lens.

Furthermore, the distortion present in the image of the staircase pattern clearly shows that regions closest to the nose yield the greatest distortion. Horizontal lines are switched end for end and tilted approximately 45 degrees.

The upper portion of the lens performs the best. Rays 3, 7, 11, and 15 in the meridian plane are focused on the  $y_{\text{IM}}$  axis in a fairly tight region; rays 10, 14, 17, 27, 30,

and 32 about the upper periphery in the object plane are all focused within the image core.

Table 3 explains the legend of the spot diagram. It is seen that the standard deviation in the y-direction is approximately five times that of the z-direction. This elongation of the image is not readily evident in the spot diagram since the ordinate and abscissa have not been plotted with equal increments. This results from the great disparity between Spot Diagrams of the various lens designs studied. It is important, therefore, for the reader to take careful note of the relative sizes of the  $y_{\text{TM}}$  and  $z_{\text{TM}}$  axes.

Of primary importance is the spot size in Figure 13. The value for spot size is adversely affected by the poorest performing regions of the lens. Were it not for these errant rays, the spot size would be considerably smaller. The image intensity pattern is benefitted, however, by the fact that the rays spread the farthest from the centroid contribute significantly less energy per ray than those being focused in the core of the image. Table 4 lists the relative intensities of the primary skew rays plotted in Figure 13. It is seen that the high intercept angles experienced by the rays closest to the bottom of the lens and, to a lesser extent, those near the nose, result in higher reflectivity and lower transmission through the lens.

In the legend of Figure 14, the average of ray intensity is given. The Spot Diagram Energy Density distribution may be seen at a glance. The fraction of energy (number of rays)

TABLE 4

Skew Ray Intensities of HIN Lens

### a) Numberical Order

Ray	Intensity	Ray	Intensity
1	0.746	11	0.919
2	0.789	12	0.918
3	0.921	13	0.910
4	0.892	14	0.887
5	0.877	15	0.918
6	0.811	16	0.916
7	0.920	17	0.910
8	0.917	18	0.640
9	0.904	19	0.576
10	0.870		

# b) In order of Descending Intensity

Intensity	Ray(s)	<u>Intensity</u>	Ray(s)
0.921	3	0.887	14
0.920	7	0.877	5
0.919	11	0.870	10
0.918	12,15	0.811	6
0.917	8	0.789	2
0.916	16	0.746	1
0.910	13,17	0.640	18
0.904	9	0.576	19
0.892	4		

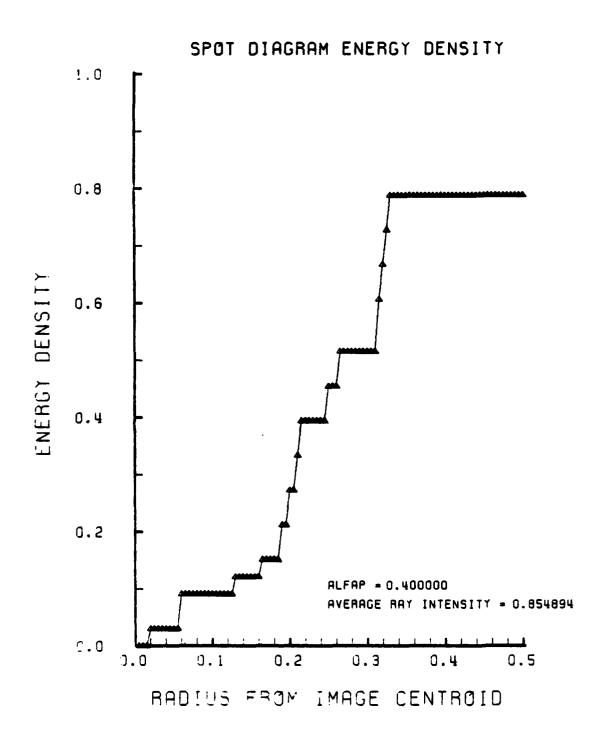


Figure 14. Nondimensional Encircled Energy Plot for HIN Lens Design of Figure 11

shown as a function of radius from the Spot Diagram centroid is normalized with respect to the total number of rays traces (Rays, in Figure 12) regardless of whether all of the rays successfully intercepted the image plane. Thus the fourth basic plot may be used in conjunction with the Spot Diagram to further define the image concentration with respect to the centroid.

The response of the HIN lens to increasing tilt angles is given in Figures p-1through p-23 where the index of refraction of the lens has been set at 1.5. Whereas the spot size at  $\alpha_p = 0.0$  is very small, that at  $\alpha_p = 0.7$  radians is quite large at 67% of the lens radius. Figures 15 and 16 summarize spot size and centroid locations for the lens. Spot size growth, although somewhat irregular at the higher angles, is pronounced. Furthermore, the centroid movement is seen to be an approximately linear function of  $\alpha_p$  until  $\alpha_p = 0.3$  and easily exceeds the radial dimension of the lens at higher tilt angles.

Figure 17 shows that the HIN lens may be slightly improved by increasing the edge thickness, T, by a small amount. The lens may be otherwise tuned to improve performance at certain tilt angles by designing the lens with U slightly greater than zero. These performance improvements are practically insignificant, however, and neither lens tuning by the parameter U nor T produce improvement across the spectrum of tilt angles. Instead, an improvement at one  $\alpha_p$  usually has resulted in a degradation at others.

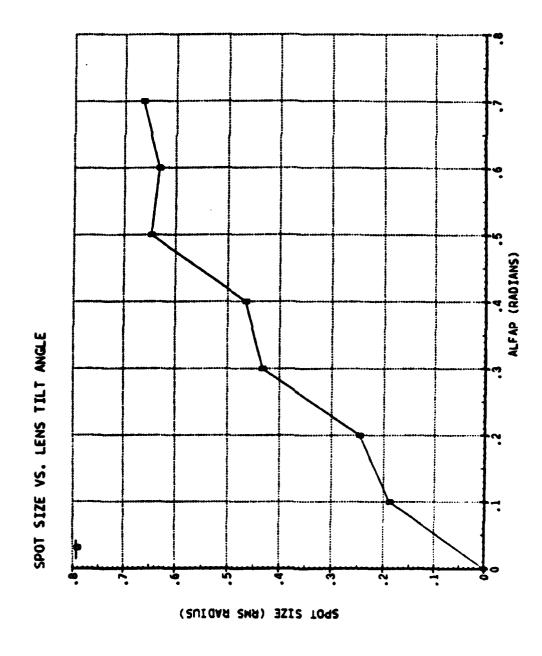


Figure 15. HIN Lens Spot Size versus  $\alpha_p$  for  $N_2 = 1.5$ 

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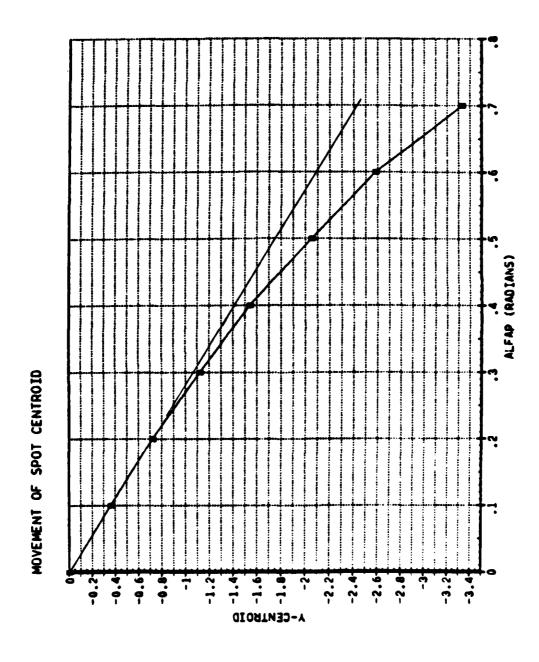


Figure 16. HIN Lens Centroid Movement versus  $\alpha_p$  at  $N_2 = 1.5$ 

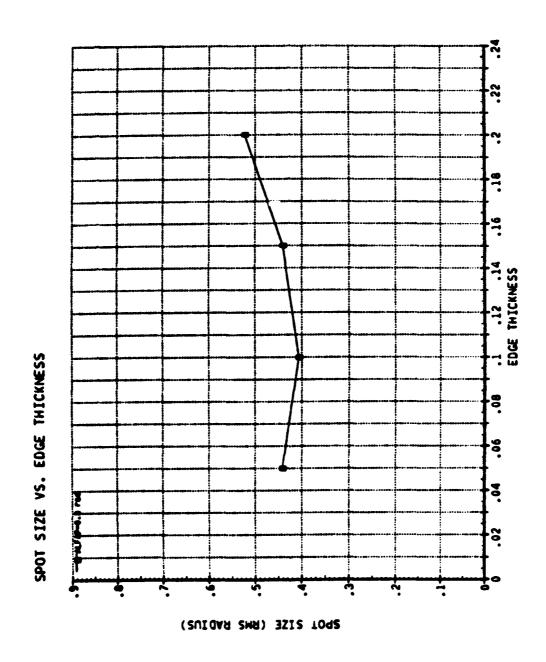


Figure 17. HIN Lens Spot Size versus Edge Thickness at  $\alpha_p = 0.3$  Radians and  $N_2 = 1.5$ 

Finally, the homogeneous lens is shown with index of refraction of three at the intermediate tilt angle of 0.3 radians in Figures D-24 through D-27. The lens shape required to accommodate the higher index of refraction is seen to be thinner and displays less outside surface curvature than the HIN lens with  $N_2 = 1.5$ . Spot size is significantly reduced. Since infrared lenses such as germanium generally have fairly high refractive indices, improved lens performance at these higher values is encouraging.

## VI. GRIN LENS RESULTS

The performance of the GRIN lens is similar in many respects to the homogeneous lens. The relationship between characteristic regions of the GRIN lens, such as the lower lens portion, and where these regions image bundles of rays is identical to the HIN lens as depicted in Figures 11 and 12. The growth of spot size with increasing  $\alpha_p$  and the respective movement of image centroid typical of the homogeneous lens is clearly displayed by the GRIN lens as well.

The measure of the superiority of the GRIN lens, therefore, lies in the successful correction or improvement of the deficiencies seen in the HIN lens. Here, the reduction of spot size is of primary concern.

As a modest example of the ability of the gradient refractive index to reduce spot size, the GRIN lens design shown in Figure 18 is examined. This lens is very similar in shape to the homogeneous lens with  $N_2=1.5$ . Note, however, that unlike the HIN lens, the nose half-angle,  $\gamma$ , is slightly larger than that of the cone, indicating more outside surface curvature. The object plane for this lens, at  $\alpha_p=0.4$ , is shown in Figure 19. Here, the error in focusing rays present in the image plane has been superimposed over the grid plane.  $\gamma_{IM}$  coordinate error contours (with respect to the centroid) are shown in the left half of the plane and  $z_{IM}$  contours in the right. These contours vividly show that the largest

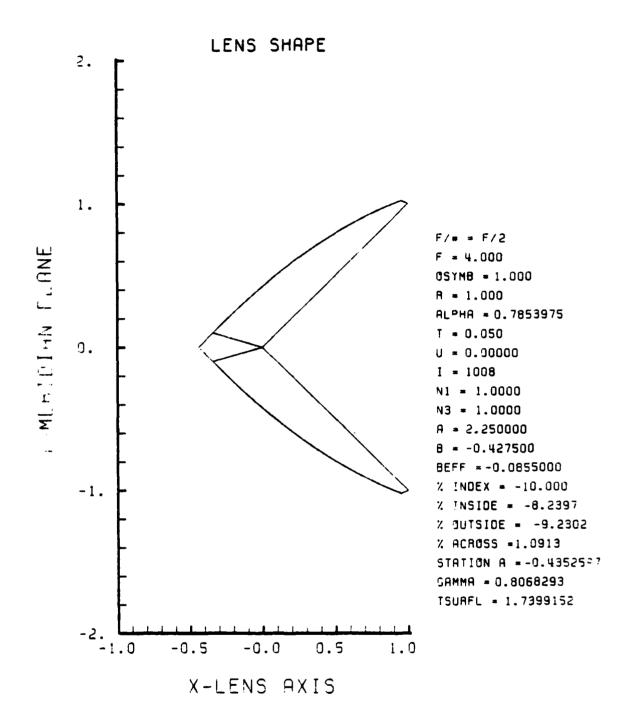
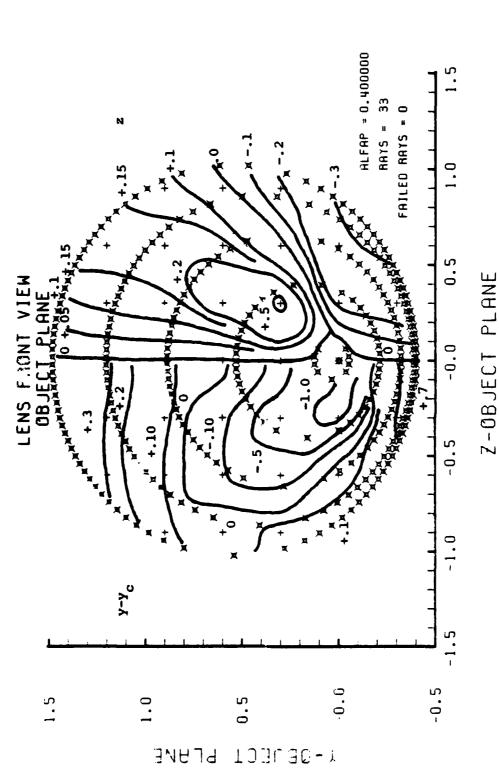


Figure 18. Example GRIN Lens Design with 10% Negative Gradient at OB = 1.0



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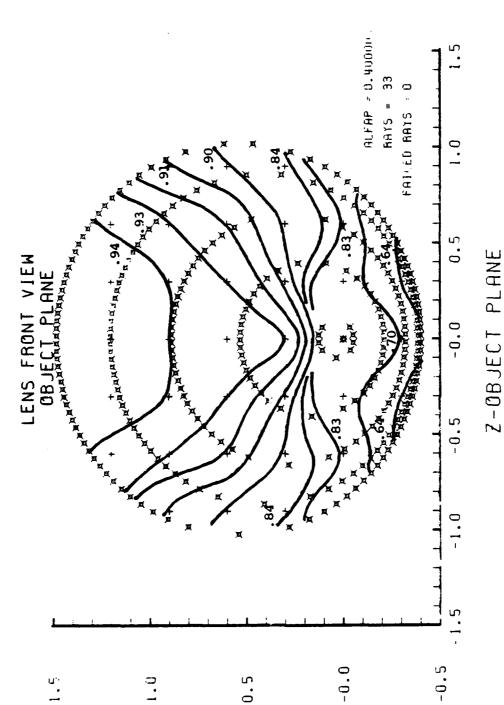
Displace-Example GRIN Lens Object Plane. Contours in the Left Half Depict the Difference Between the  $\,\mathrm{y}\,$  Coordinate in the Image Plane and Image Centroid. The Right Hand Contour Shows the Corresponding ment in the Image Plane. Figure 19.

errors in both y and z stem from the nose and bottom of the lens. The best performance is contributed by the upper central region. The reader will note that negative errors in z occur below the  $y_0 = 0$  plane of the lens and that negative errors in y straddle the  $y_0 = 0$  plane.

The corresponding intensity contours are displayed in Figure 20. Of significance here is that regions which perform relatively well in y and z error also perform well in transmitting energy. Regions of the lens which are characterized by relatively high angles of incidence, therefore, perform the poorest and have the most to gain from better combinations of gradient index and  $O_c$ 

The Spot Diagram, Figure 21, displays an image pattern virtually identical with the HIN lens in Figure 11. The GRIN spot size,  $\sigma_{\mathbf{r}}$ , however, is slightly smaller by approximately 2% and both  $\sigma_{\mathbf{y}}$  and  $\sigma_{\mathbf{z}}$  are correspondingly smaller which indicates superior performance by the GRIN version. Image centroid for the GRIN is displaced further than that for HIN. The average intensity transmitted by the GRIN lens compares favorably at 0.88 as opposed to the average intensity of the HIN at 0.85. Further evidence of superiority of the GRIN lens is seen in the form of a steeper slope to the Encircled Energy Plot (Figure 22) than that of the HIN lens in Figure 13.

Although the example GRIN design displays marginal superiority, some GRIN designs do not. In particular, it



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Intensity Contours of the Example GRIN Lens Shown in Figure 45 Figure 20.

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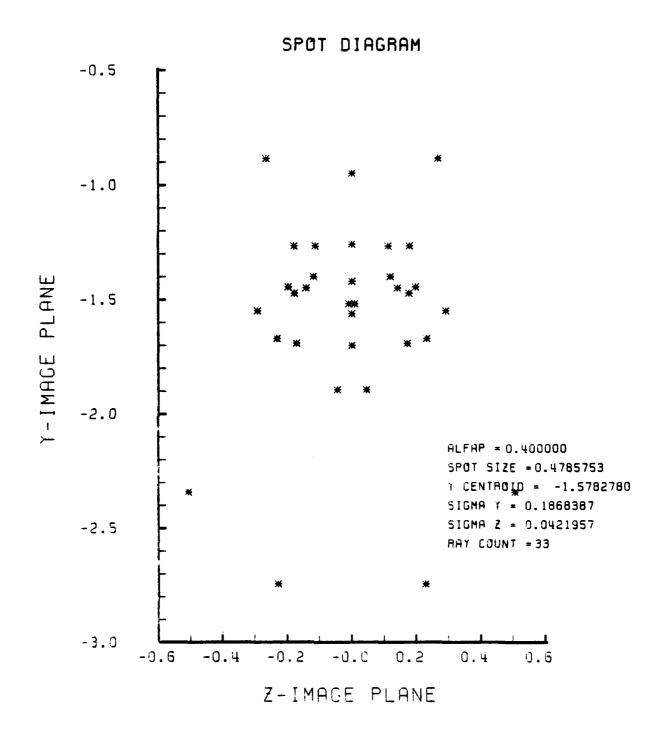


Figure 21. Example GRIN Spot Diagram at  $\alpha_p = 0.4$  Radians for GRIN Lens Design Shown in Figure 18

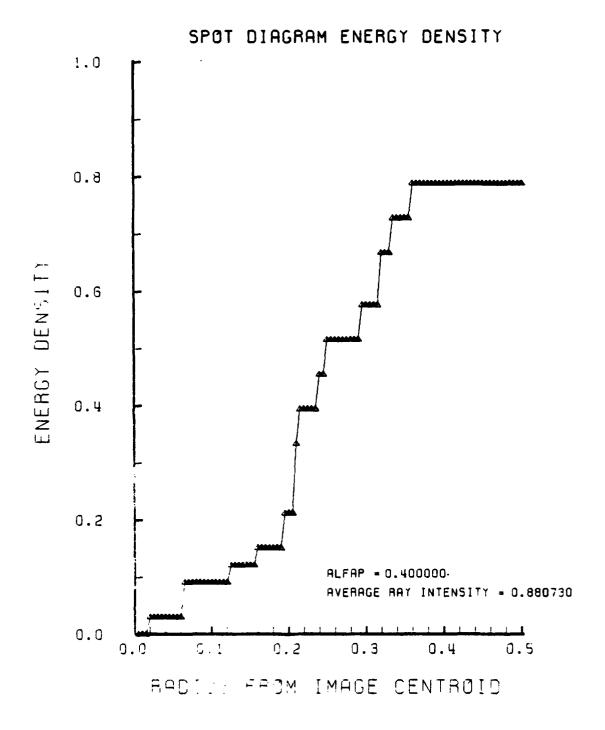


Figure 22. Example GRIN Encircled Energy Plot Corresponding to Figure 21

has been found that as the center of symmetry of the index function,  $O_s$ , is moved further away from the interior of the GRIN lens, performance deteriorates.

Two series of GRIN lens designs have been examined at the intermediate tilt angle of 0.3 radians in order to define the spectrum of lens performance. Both design series explore the GRIN lens for  $O_S$  inside, outside, and far outside the interior of the lens. The first series, Figures E-1 through E-92, are GRIN designs in the "low range" of refractive index with the index parameter "a" set at 2.25. This first set is compared to the HIN lens with  $N_2 = \sqrt{a} = 1.5$ . The second, or "high range" series is for a = 9.0, as compared to the HIN at  $N_2 = 3.0$ , and are displayed in Figures F-1 to F-128. Both the high and low range lenses exploit gradient changes of 5, 10, 25, and 50 percent both positive and negative, where possible. A negative 50% gradient in the low range is, of course, not possible since an index of below 1.0 would result. The center of symmetry was not located at  $x \ge 0$  due to the aforementioned singularity encountered at very small angles.

The lens shapes of the GRIN designs shown vary widely. The outside surface may be either convex or concave and exhibits a thinner profile at higher values of refractive index. Although all of the resultant GRIN lens shapes are superior aerodynamically to the hemispherical seeker lenses currently in use, the convex version of the lens has more

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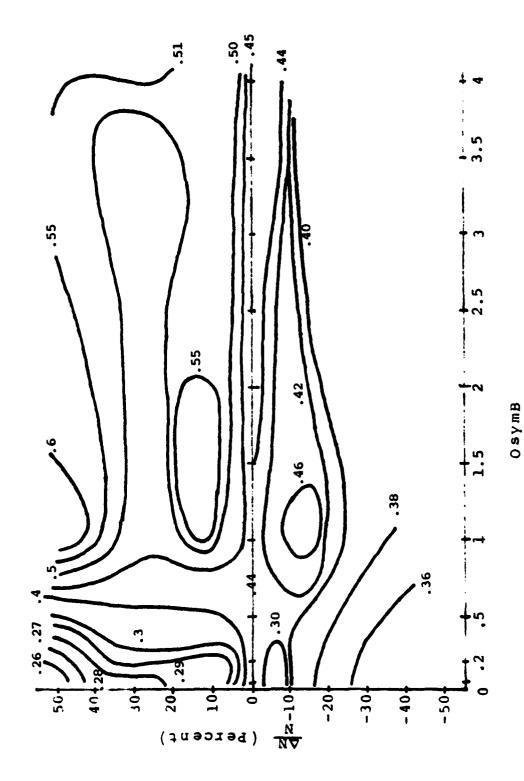
obvious applications for a Sidewinder type tactical missile, whereas the concave lens shape is more applicable to the diffuser of a ramjet with nose inlet.

મિલ્લોએક સ્થાપનો લોક્સ માર્ચિક સોક્સ કોલ્સો કોલ્સોએક જેવી છો જેવે જેવા છે. જે કોલ્સો ફોલ્સોન્સો કોલ્સોન્સો કોલો

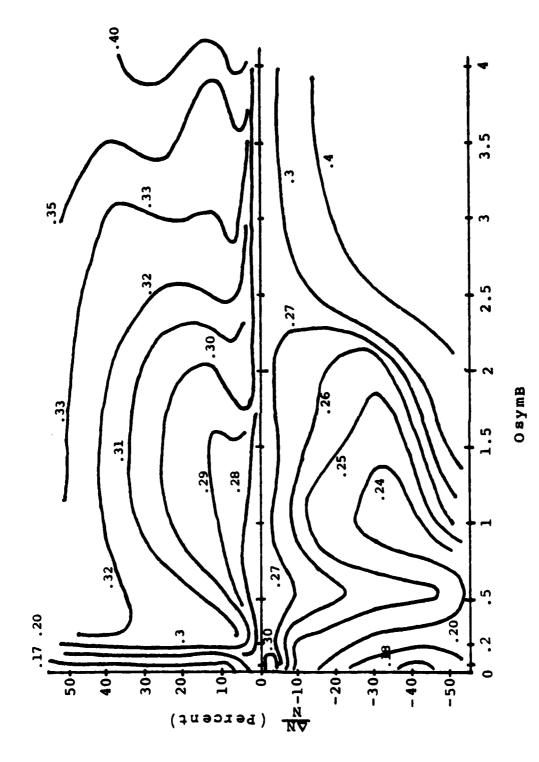
Object plane diagrams show that for lens designs exhibiting smaller spot size, fewer or none of the skew rays failed. This fact underscores the success of GRIN in controlling image deterioration contributed by the nose region of the lenses. Although the lower lens and nose regions still create the largest of the image spread, the extent is reduced.

Figures 23 and 24 are contour diagrams summarizing spot size performance of GRIN lens designs in the low and high range respectively. From these diagrams, it can be seen immediately that the best performance is obtained from lenses having the center of symmetry inside the lens. From this fact, it may be deduced that large changes in refractive index are desired along the surface of the lens rather than across. Except for isolated regions at -5%, increasing gradients produce smaller spot sizes. The positive gradient in the high range, however, has an almost constant spot size from +5% to +50% where an improvement of only 0.6% is seen. In both the low and high ranges the positive gradient at 50% and at OB = 0.05 has proved to be the best performing combination. Again, as in the HIN lens the high range exhibits the smaller spot size.

The best GRIN lens is now further examined over the range of obliquities up to 0.8 radians. Furthermore, since it has



Contour Plot Summary of GRIN Lens Spot Size Performance for  $\alpha_D=0.3~\text{radians},~\alpha=2.25$ Figure 23.



Contour Plot Summarizing GRIN Lens Spot Size Performance at  $\alpha_p$  = 0.3 radians, a = 9.0 Figure 24.

been determined that placing the focal point at x=2.0, improves spot size performance at lower values of  $\alpha_p$ , another series of plots is given for F/l. With the focal point at x=1.5, severe degradation results; for F=3.0 only modest differences with F/2 are evident. Appendices G and H show the performance of the "best" GRIN lens at F/l and F/2.

Here it may be fully recognized that the resulting behavior of a "good" gradient refractive index seeker lens is far superior to that of the HIN lens for  $\alpha_p \leq 0.6$  radians; see Table 5. The reader will note that there are no failed rays even at 0.8 radians (45.8 degrees). Although centroid movement at F/2 is slightly greater than the HIN lens, by changing to F/1, this can be corrected; see Figure 25. The F/1 version, moreover, displays a range of spot sizes significantly smaller than that for the F/2 lens below  $\alpha_p \stackrel{>}{=} 0.54$  radians or the HIN lens below 0.68 radians; see Figure 26. That the lens displays more desirable performance with a shorter focal length is a surprising but highly desired result. Since seeker optics systems are volumn limited, any reduction in the lens focal length is beneficial to the final design and packaging requirements.

It is significant for the reader to note that centroid movement beyond a value of  $\pm 1.0$  exceeds the physical dimensions of the lens radius. If the interior of the lens mount is of similar dimension, it follows that the requirement for  $y_{CENTR}$  to be less than  $\pm 1.0$  restricts the F/l lens to approximately 0.65 radians of tilt whereas the F/2 lens is further restricted to less than 0.3 radians. The superior performance



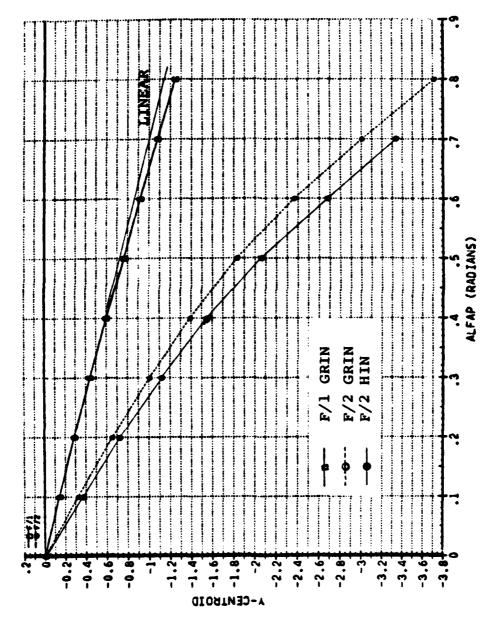


Figure 25. "Bes:" GRIN Lens Centroid Moment for F/N and F/2



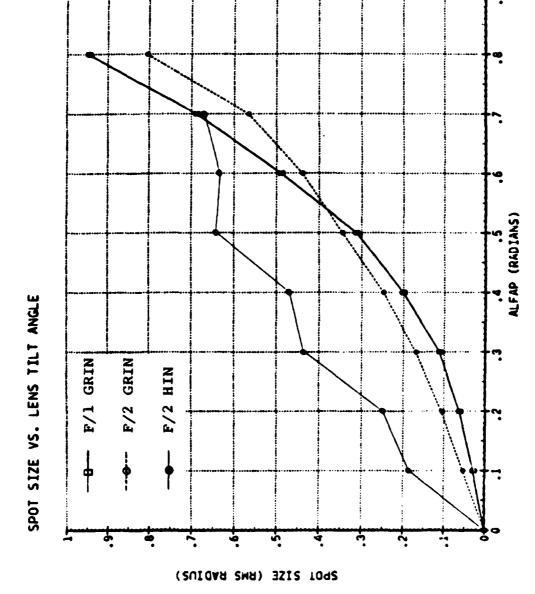


Figure 26. "Best" GRIN Lens Spot Size Performance for F/Numbers of F/l and F/2

TABLE 5

Comparison of HIN and GRIN Lens Designs

Performance Parameter or Feature	HIN	GRIN*
Shape of Outer Surface	Convex	Concave
Fraction of Failed Rays at $\alpha_p = 0.7$ radians	0.143	0.0
Linearity of Centroid Motion; $\alpha_p$ for $\Delta y/y$ deviation of 10%	0.4 radians	0.8 radians
Spot Size for ap		
$\alpha_p = 0.1$ radians	0.190	0.029
$\alpha_{\perp} = 0.2$ radians	0.247	0.063
$\alpha_n = 0.3 \text{ radians}$	0.442	0.112
$\alpha_n = 0.4$ radians	0.466	0.201
$\alpha_n = 0.5$ radians	0.637	0.313
$\alpha_{\perp} = 0.6$ radians	0.632	0.490
α = 0.7 radians	0.667	0.688
$\alpha_{p}^{r} = 0.8 \text{ radians}$	N/A	0.947
Spot radius for 80% of energy @ a = 0.4 radians	Does not attain 80% within 0.50	0.28
Average relative intensity of skew rays 0 $\alpha_p$ = 0.4 radians	0.82	0.44

<sup>\*&</sup>quot;Best" GRIN lens with 50% positive gradient, O = 0.05, in the F/l configuration.

of the F/2 lens above 0.54 radians is, therefore, unusable. Attempts to improve upon this performance by increasing the edge thickness, T, or by slightly adjusting the parameter U, as with the homogeneous lens only produced uniformly degraded performance in every respect.

Although the "best" GRIN lens resulted from a positive 50% gradient change in refractive index, it has been noted that this configuration was only slightly better than the same lens with a positive 5% change. Furthermore, since refractive index gradients of five percent or better have already been produced, it is entirely feasible that if the precise parabolic change could be controlled, this lens could be produced today.

Despite the obvious success of GRIN in controlling spot size growth and image centroid movement, a penalty in the form of reduced ray intensity has been paid. That increasing spot size performance is tempered by a loss of intensity may be seen by comparing "good" GRIN encircled energy plots with that of the HIN lens. Note that the HIN lens with  $N_2 = 3.0$  also loses  $\sim 50\%$  to intensity; see Figure 44. This loss in intensity is partially offset, however, by the GRIN lens with an increased number of rays transmitted and by the reduced spot size.

One drawback to the use of the GRIN lens as a self sufficient, single element lens is still that of the relative sizes of image and detector. Hence, even though GRIN has

significantly reduced spot size below  $\alpha_p \cong 0.6$ , that size is still significant at tilt angles above 0.3 radians. In order to use the lens without a secondary focusing element or mirror arrangement, a large composite sensor array would be required.

#### VII. CONCLUSION

Gradient refractive index materials may be employed to design a pointed seeker lens which exhibits optical performance far superior to that obtainable with conventional homogeneous optical material. A fifty percent, positive, parabolic gradient index with center of symmetry interior to the lens was found to yield the best performance although a five percent version of the same lens was only very slightly inferior; this lens may possibly be fabricated today.

For objects off-axis by more than 0.3 radians (~17.2 degrees) a secondary lens element may be required unless a large scale multiple element sensor array is employed. With such arrays, objects off-axis by more than 0.65 radians (37.2 degrees) may require Cassegrainian or other mirror elements to compensate image movement.

#### VIII. RECOMMENDATIONS FOR FUTURE WORK

This thesis investigated a spherically symmetric GRIN seeker lens with inside conical surface. Future studies should investigate:

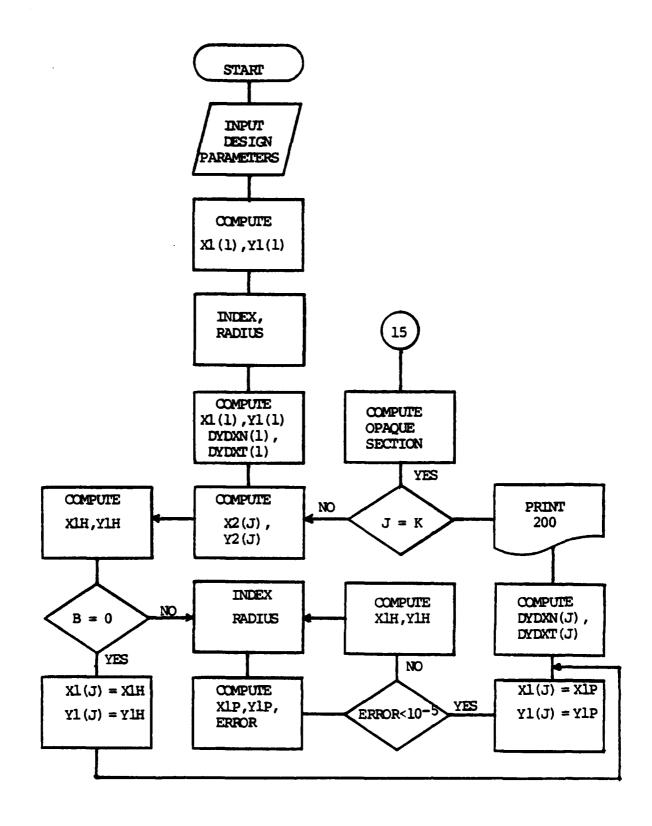
- a) The elimination of the singularity problem by modification of the theoretical equations for GRIN using numerical techniques. See Reference 13.
- b) The optimization of the GRIN lens with both inside and outside curvature.
- c) The effect on lens performance due to an attached shock wave.
- d) The effect on lens performance when the object is no longer in the far field and wave front curvature must be taken into account.
- e) The performance of the lens using wavelengths corresponding to atmospheric windows.
- f) The effect of a radially symmetric gradient index on lens performance.
- g) The feasibility of adding an anamorphic gradientindex lens located on the missile body to increase off-axis tracking/acquisition capability to 90° and beyond. See Figure 1 of Reference 14.

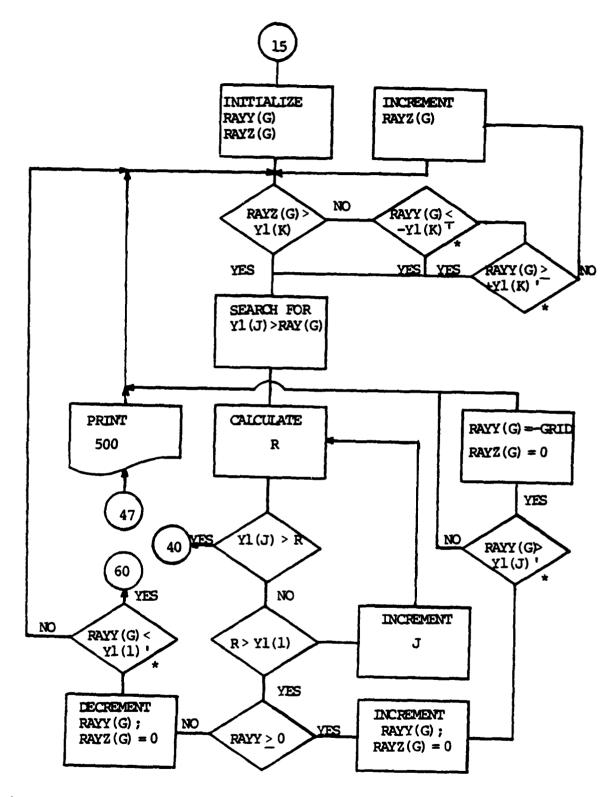
#### APPENDIX A

#### COMPUTER LOGIC FLOW DIAGRAM FOR PROGRAM GISL

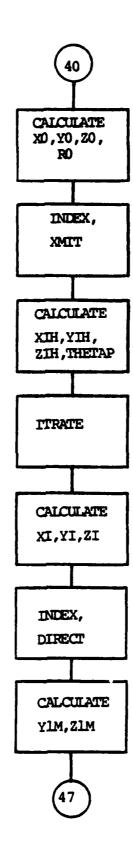
This appendix together with Appendices B and C describe the FORTRAN program GISL. GISL may be used to design either a homogeneous (B = 0), or a GRIN lens to the user's specifications by changing the design parameter where indicated by comments in the input section of the program. Additional software required to run GISL and plot the results are not described since these system-dependent procedures do not apply elsewhere.

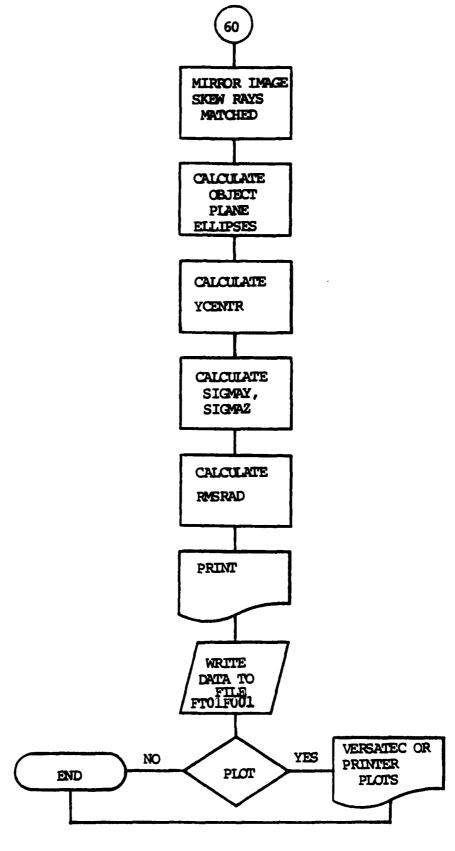
The flow diagrams which follow provide the reader with the information necessary to follow the fundamental logic of the main program and subroutines of GISL. Not shown is Subroutine DIRECT which has been derived from Subroutine DIRECT as described by Amichai [5] with only slight modifications.

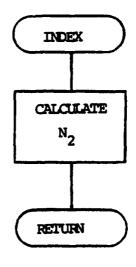


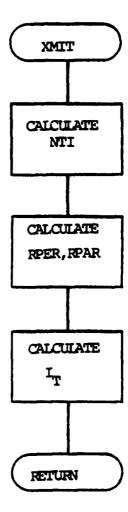


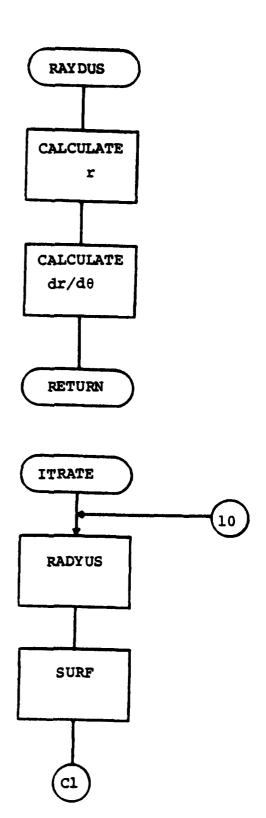
<sup>\*</sup>Here, the prime indicates Y-coordinate as transformed into the grid plane.

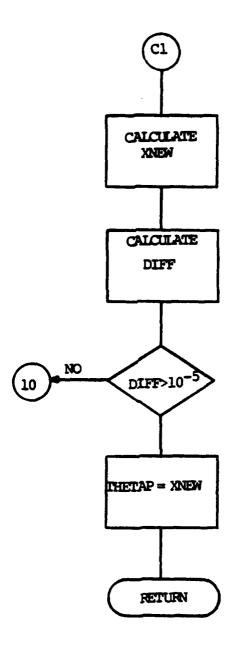


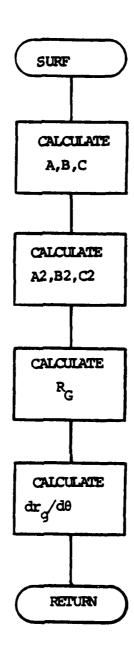












## APPENDIX B

## LISTING OF FORTRAN PROGRAM GISL

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THE FOLLOWING ARE CONSTANTS A & B FOR N2(R).
NOTE THAT IF B=0.0, N2(R) NO LCNGER VARIES AS WHICH CORRESPONDS TO THE HOMOGENEOUS LENS WIT A=9.00
B=+11.25
THE FOLLOWING IS THE DISTANCE FROM 0.0 TO THE SYMMETRY OF THE GRADIENT INDEX ALCNG THE X-AX POSITIVE SENSE IMPLIES THAT OSYM IS TO THE LECCNVERSELY. THE NEGATIVE SENSE IMPLIES THAT OSYM IS TO THE CONSYMB=+0.05
N3=1.0
BETA=ATAN(R/(F-R\*COTAN(ALPHA)))
PI=3.141592653589793
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F=4.0
INSIDE SURFACE CONE HALF-ANGLE:
ALPHA=0.7853982
INCIDENT RAY OFFS ET ANGLE:
U=0.000000
NUMBER OF ITERATIONS (MUST BE AN I = 100
MAXIMUM RADIUS OF INSIDE SURFACE R=1.0
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T=0.05
OR, IF IT IS DESIRED TO HAVE THE TO T=0.05
OR, IF IT IS DESIRED TO HAVE THE TO T=0.05
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VARY = 0
SIZE OF GRID INCREMENT:
GRID=0.3
ELLNUM=4
RUMBER OF ELLIPSES = ENUM ELLNUM MUST BE EVEN AND I/ELLNUM MUST BF AN INTEGER WITH NO REMAINDER ENUM=ELL NUM+1
ELL IS THE Z COORDINATE INCREMENT FOR ELLIPSE PLOTS ELL=GRID/7.5 KNOWNS FOR TRACE RAY:

SHAPE (SEE INPUT ABOVE) GLM LENS ALGER ITHM

AT THIS POINT, IF "THKNES" IS 1, THE REMAINDER OF THE PROGRAM
IF THKNES - EQ. 1) T=0.05 TO T=0.55 BV A 0.05 INCREMENT STEP.
60 TO 3
CONTINUE
T=T+0.05
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BF=F/R
WRITE(6,1900)
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J=1 UDP(J)=BETA X2(J)=(BF\*SIN(UDP(J)))/(SIN(UDP(J))+(TAN(ALPHA))\*COS(UDP(J))) Y2(J)=(X2(J))\*TAN(ALPHA) I 2P(J)=-ALPHA+ATAN(BF-X2(J))/Y2(J)) I 2P(J)=-ALPHA+ATAN(BF-X2(J))/Y2(J)) RO=SQRT((X2(J)+GSYMB)\*\*Z+Y2(J)\*\*Z) RZENO=RO I F(USYMB .LT. 0.0) RZERO=2.0\*PO CALL INDEX(RO, N20) I 2(J)=ARSIN((N3/N20)\*SIN(I2P(J))) UP(J)=PIZ-ALPHA-I2(J) BASE=X2(J)+CSYMB

J), Y1(J), X2(J), Y2(J), UDP(J), I2P(J), I2(J), UP(J), L)\*DLUDP {UDP(J)))/(SIN(UDP(J))+(TAN(ALPHA))\*CNS(UDP(J))) \* TAN(ALPHA) +OSYMB)\*\*2+Y2(J)\*\*2) ROOO=RO SINCE **TEXT** THE FL = 0.0 10 J= 2.K L=J-1 LDP(J) = BE X2(J) = (BF Y2(J) = (RF Y2(

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12 P(J) == ALPHA+ATAN ((BF-X2(J))/Y2(J))

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F(DSYMB .LT. 0.0)N2X1K)

ALL INDEX(RAD,N2X1K)

ALL INDEX(ROOO,N 2X2K)

ACL INDEX(ROOO,N 2X2K)

RCNTE=-(N2X2K-N2X21)/N2X2K)*100.0

RCNTE=-(N2X2K-N2X1K)/N2X2K)*100.0

RCNTO=-(N2X1K-N2X1K)/N2X1K)*100.0

RCNTO=-(N2X1K-N2X1K)/N2X1K)*100.0

RCNTO=-(N2X1K-N2X21)/N2X1K)*100.0

RCNTO=-(N2X1K-N2X21)/N2X1K)*100.0
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TAY=Y1(N)-Y1(P)

TAY=Y1(N)-Y1(P)

TAY=Y1(N)-Y1(P)

SQRT((RADIUS-Y1(P))\*\*2+(XO-X1(P))\*\*2)

SQRT(OEL TAX\*\*2+DELTAY\*\*2)

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TOP SP/ST

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CCNTINUE Y G = 9 5999.

CCNTINUE Y G = 9 5999.

TF ( SQURAY . EQ. 1) WRITE(6,500)G, RAYY (G), RAYZ (G), XD, YD, XI, I F ( SQURAY . EQ. 1) WRITE(6,500)G, RAYY (G), XDIAPT(G), XDIAPT(G) . YDIAPT(G) REFLECTION REFRACTEC INSID .6T A T AND CMPP IP) - (NZ/N3) \*COS (PHII)
SINES OF INSIDE EXTERNAL I
CKP-NUMB\*LKP
CCMP-NUMB\*LLP
CMP-NUMB\*LMP ALTINERNAL REFLECTION: TIRA) GO TO 52 ((N2/N3) \*SIN(PHII)) OF TRANSMITTED INTENSITY IDENCE AND TOTAL INTERNAL 47 CCNTI 457

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WRITE(6,802) G,RAYY(G),RAYZ(G)
                                                                                                                 WRITE(6, 600)G, RAYY(G), RAYZ(G)
                                                                                                                                                                  WRI TE(6, 700)G, RAYY(G), RAYZ(G)
                                                                                                                                                                                                                 WRITE(6,800)G,RAYY(G),RAYZ(G)
                                                                                                                                                                                                                                                                                                                                                                                             ) COUNT=COUNT+1
                                                                                                                                                                                                                                                                                               MIRROR IMAGE RAY MATCHING FOLLOWS
                                              +GRID
O) COUNT=COUNT+1
                                                                                                                                                                                                                                                                                                                           .EQ.1) WRITE(6,401)
                                                                                                                                                                                                                  . EQ.
                                                                                                                                                                                                                                                                  . EQ.
                                                                                                                                                                                                6)=0.0
                                                                                                                                                                                                                                                                                                                  CONTINUE
IFCSQURAY
S=6-1
CCNTINUE
IFCSAYZ(S)
CONTINUE
                                                                                                                                                                                                                                     (L)
                                                                                                                                                                                                                                                                                                                  9
45
                                                                                                                                                                                      52
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                                                                                                                                      51
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CCC

ITNCTY(S) EQ. 1) WRITE(6,50116,S,RAYY(G),RAY2(G),OPL(G), G),NTNCTY(G) ELY(NUMB, COCRD) = (X1(NLMBR) + AB) \*SIN(ALFAP) + SQRT((Y1(NUMBR)) \*\*2-ELZ(NUMB, COORD) | \*\*2) \*COS(ALFAP) ELZ = ELZ(NUMB, COORD) | F(ELLIPS . EQ. 1) WRITE(6,1100) NUMBR, COORD, ELY(NUMB, COORD), ELZ(NUMB, COCRD) | ELZZ+FLL ELZ(NUMB, COORD) = ELZZ+FLL F(ELZ(NUMB, COORD) . GE. Y1(NUMBR)) GO TO 73 .6E. 2) GD TD 64 .EQ. 1) WRITE(6,601)G,S,RAYY(G),RAYZ(G) TINUE FLAG(S) .GE. 3) GO TO 65 F(SQURAY .EQ. 1) WRITE(6,701)G,S,RAYY(G),RAYZ(G) O TO 62 . EQ. 1) WRITE(6,801)G, S, RAYY (G), RAYZ(G) .EQ. 11 WRITE(6,8031G, S, RAYY(G), RAYZ(G) CONTINUE
IF (FLAG(S) .EQ. 4) GO TO 67
IF (SQURAY .EQ. 1) WRITE (6,801) G,S,RAYY (G),RAYZ
GO TO 62
CONTINUE
IF (SQURAY .EQ. 1) WRITE (6,803) G,S,RAYY (G),RAYZ (G) TO 62
CONTINUE
RAYS=G-1
BRAYS=G-1
IF (SQURAY .EQ. 1) WRITE (6,900) RAYS, DRAYS, COUNT .EO. 1) WRITE(6,1000)ENUM GENERATE OBJECT PLANE ELLIPSES IF(ELLIPS .EO. 1) WRITE(6,10 NLMBR=1 NUMB= 1 CCNTINUE CCORD=1 ELZ(NUMB,COORD)=- Y1(NUMBR) CONTINUE 49 70 19 ပပပ

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\*83) F.ALPHA, U. I. R. T. NI, N. 2, N. 3, AL FAP, GRID, STAT NA, GAMMA, RAYS, GGT, IANGLE, THICK, OS YMB, TSURFL, BEFF, PRCNTO, PRCNTE, PRCNTT, B. ., 2F9.7, 14,4F9.7/\*\*, 3F9.7,F9.5,F9.7/\*\*, 614,4F11.7/1X, SPOT DIAGRAM ENERGY DENSITY VS. RADIUS FROM CENTROIC: G=1,RAYS FLAG(G) GT 0) GQ TQ 81 SDRAC(G) .LE . ROC(GG)) SUM4(GG)=1+SUM4(GG) OUTPUT FOR PLOT INTO FILE FTO1F001 184)XI(J)YI(J),X2(J),Y2(J) ÎN(GG)= SUM4(GG)/RAYS E(6,1400) GG, ROC(GG), FRACTN(GG) • EQ. .AND. THKNES GC=GG+1 ROC(GG)=RCC(HH)+0.005 SUM4(GG)=0.0 IF(ROC(GG).GT.0.5) GC TO 80 [66] =0.005 (66] =0.0 [NUE PLOT OUTPUT CONTIN WRI TE 83 52 80 887 81 82 86

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ENGTH
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21.1/14X,
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ITA:'//2X,'NOSE HALF ANGLE =
OPAQUE SURFACE LENGTH = ',
F7.5/2X,'TOTAL LENS SURFACE
                                                                                                                                                                                                                                                                                                                                                                           ZH-
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PHS
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                                                                                                                                                                                                                                                            , EL Z(NUMB, COORD
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SAHICKNO
FRO: NA
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OR PE
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3F10.7,2E13
FRCEPT WITH
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7 / 5 X ; 
7 - 5 / 5 / 5 / 5 / 5 / 5 / 7 + 10 N ZE
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0.7/20x,4F10.7/
0.5E SECTION DAT
A = (F9.5/2x)
FACE LENGTH = (
                                                                                                                                                                                                                                                                                                                                                                        AND:
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ING COMMA
IS PRINTER
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300 FORMAT(10/114,7F10,7
1F7.5,2X, STATION SE
1 TRANSPARENT SURFACE
1 = , F7.5// )
0 FOPMAT(11, SKEW RAY
1 GRID= , F9.7,2X, SEF
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                          7/8,
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4.7F10
4.2F10
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800
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700 FORMAT(1X,14,2F10.77,* INTERCEPT CUTSIDE BOUNDARY OF 2ND SURFACE'/16180826
401 FORMAT(1X,14,2F10.77,**) INTERCEPT CUTSIDE BOUNDARY OF 2ND SURFACE'/16180826
401 FORMAT(11,11,11,12,14,2F10.77,**) INTERCEPT CUTSIDE BOUNDARY OF 2ND SURFACE'/16180826
501 FORMAT(11,11,11,12,14,2F10.77,**) INTERCEPT CUTSIDE OF BOUNDARY OF 2ND SURFISSORY
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, RO, SORT, COS, SIN, ARSIN
              FUNCTION
E INDEX
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                   CENTE
                                                                                                                                                                                                          BROUTINE TRANSMIT CALCULATES THE TRANSMITTED INTENSITY ACH RAY AT BOTH THE OUTSIDE AND INSIDE SURFACES USING RESNEL EQUATIONS.
                 S
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                   FROM
          ALCULATES THE INCEX OF REFRACTION A E CENTER OF SYMMETRY. THE VALUES OF ARE USER INPUTS.
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                CALCULATES THE RADIAL DIMENSION IE RAY AT THE ANGLE SPECIFIED. 5 D(R)/D(THETA) FCR SUBR ITRATE.
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                    SE S
                                                                                                                                                                                                                                                                                                                                                                                                                                                                CSP+SQ)
SP-SQ)/((NTI**2)*CGSP+SQ)
**2+RPAR**2)
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                  LT
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                R, DRDTEG, FI
ON, ARCSYN, F
                                                                                                                                                                                                                                                                            SUBROUTINE XMIT(PLI,NZ,XM1,FACE)
REAL PHI N3 XM1,NTI,NI,NZ,SQRT,COS
INTEGER FACE
CCMMON/XM/N1,N3
                                                                               (R,N)
RO,E,EPSLON,SQRT
RC,EPSLON,PI2,J
ERO)**2)
                                                                                                                                                                                                                                                                                                                                                                               17 I=N3/N2
I (PHI )**2)
I ) GO TO 10
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 S.
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 ANGL B
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                              SLBROUTINE RADYUS (RC, AI
REAL ANGLE, R, E, RA T, A, B,
INTEGER J, FILTER
CGMMON A, B, E, RZERC, EPSI
RAT=SQPT (A**2+4.*B*(E**)
ARC SYN=ARSIN ((2.0*(E/R)
                                                                                                                                                                                                                                                                                                                                                            NTI=N2/N1
IF (FACE NTI ** 2- (SIN (PACE NTI ** 2) * (CCS PECOS (PHI) ** 2) * (CCS PACE NTI *
                INDEX CA
FROM THE
A AND B
                                                                               INDEX (R
P.RZER C
E.RZER C
E.RZER C
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                UBROUTINE RADYUS
OF SYMMETRY TO TH
RADYUS ALSO FINDS
                                                                                  BAB €
                SUBROUTINE
OF RADIUS
CONSTANTS
                                                                              SLBROUTINE
REAL R.N. A
COMMON A.B
N=SQRT(A+B
RETURN
```

[2, J ZERO\*#2] - A) /RAT)

C, EPSLON, PI2 +B+(E++2)/RZE )+(E/RC)++2-A

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S

•O\*EPSLON\*ANGLE+ARCSYNJ ERE F TO ST OR 2 0) ARCSYN=2\*PIZ-ARCSYN 0N\*2\*0\*ANGLE+ARCSYN) 1 °LE\* 0\*0) GO TO 10 E //SQRT(A+RAT\*FRACTN) - C\*989) GO TO 10 R\*\*3)/(2.0\*E\*\*2))\*COS(-; ı 7 =2 LTER= RACTN) .LT. 0.0) FILTER FRACTN=SIN(-EPSCON\*
IF((A+RAT\*FRACTN))
IF((A+RAT\*FRACTN))
IF(FRACTN)

SLBRGUTINE SWITCH CALCULATES THE PGINT ON THE RAY WHI DETERMINE WHETHER PSIP SHOULD BE EVALUATED IN THE IN QUADRANT. SLBRGUTINE SWITCH (RC, THTACR) REAL THTACR, RATIE, EPSLGN, A2B, PI2, RO, SQRT, ARSIN CCMMGN A, B, E, RZEGO, EPSLGN, PI2, RO, SQRT, ARSIN RATESORT (A\*\*2+4.\* B\*(E\*\*2) / RZERG\*\*2) THTACR=0.5\*(PI2-ARSIN((2.0\*(E\*\*2/RO\*\*2)-A)/RAT)) IF(THTACR .LT. 0.0) THTACR=THTACR+2.0\*PI2

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THETA SLBROUTINE SURF CALCULATES BOTH THE RADIUS TO THE LOCUS OF THE INTERCEPT OF THE RAY PLANE AND THE SIDE SURFACE AND THE DERIVATIVE OF THE RADIUS WRT GIVEN THE ANGLE THETA. SURF IS DESIGNED PRIMARILY FOR USE WITH SUBROUTINE ITRATE.

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ER) CLLP,CMMP,N20,ALPHA AN, COTAN, SORT HETA, R, RDOT, SIGN, FILTE C1 OSYMB, YO, ZO, RO, CKKP C1, NPOX, NPOY, NPOZ, PSI CX, NP OY, NP OZ, SIN, COS, SUBROUTINE SURFITHE CCMMON/SUR/PSIO, XC COMMON/DIR/AI, BI, CI REAL N20, OSYMB, NP CI

ALFA=COS(THETA)-SIN(THETA)\*COTAN(PSIO)
BETA=SIN(THETA)/SIN(PSIO)
A=ALFA\*(XO+CSYMB)/RO+BETA\*CKKP
B=ALFA\*YO/RO+BETA\*CLLP
C=ALFA\*ZO/RO+BETA\*CKMP
A2=B\*\*2+C\*\*Z-(A\*\*Z)\*(TAN(ALPHA)\*\*Z)
BZ=B\*\*2+C\*\*Z-(A\*\*Z)\*(TAN(ALPHA)\*\*Z)

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S THE VECTOR DIRECTION COSINES OF INTERCEPT WITH THE INSIDE SURFACE. 0 CALCULATES POINT OF NE DIRECT (RAY AT THE LBROUTIN A SKEW F

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.0*APP
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                          CR. RRX -LT. 0.0) SIGN=+1.0

AND. RRX -LT. 0.0) SIGN=-1.0

FZ*RZ1/RRX1**2+1.0+(NPFY/NPFZ)**2

/(RRX**2)*(RRY-NPFY/NPFZ*RZ)

X)**2-1.0

P)+SIGN*SQRT(BPP**2-4.0*APP*CPP)/(2,CLP)

(RRY-NPFY/NPFZ*RZ)*CLP)/RRX
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                  FY/NPFZ*RRZ)) .GT. EPS) GO TO 47
)/R RX
T((1.0-CKP**2)/(1.0+(NPFY/NPFZ)**2)
if2*(LP
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                    0.0) SIGN=-1.0
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                       P**2-1 0
P**2-APP*CPP1/APP
                                      **2
| RRZ **2
| **2-1.0
| BP P**2-APP*CPP |/APP
| RX*CKP| /RRZ
                                                                                                                                                                                                                                                                                                                                                                                        EPS) GO TO 45
RY-NPFY/NPFZ*RRZ
                                                                                                                                                                                                                                                                                                                                                                                                       EFS! 50 10 45
RRY-NPFY/NPFZ*RRZ
CLP
.O-CLP**2-CMP**2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                       9
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x x + 2 + C L P

x + 2 + 1 • 0 1 + C L P

I GN+5QR T (BP P
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                        .1.0
NPFZ
                                                                                                                                                                                                                                                                        SIR // RX
SQR T(1 . 0-CKP**2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                    EPS160
CLP=0.0

APP=1.00+(FRX/RR2) **2

CPP=COS(PSIR) *RR2 **2

CRP=BPP/APP+SQR | RR2 | RR2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                            45
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CKP=(-NPF Y\*C LP-NPFZ\*CMP)/NP FX

RETURN

LIF (BB = GT = C=0) S 1GN=-1.0

AF=(NPFY/NP FX) \*\*2 +1.0

BP=(NPFY/NP FX) \*\*2 +1.0

BP=(NPFY/NP FX) \*\*2 +1.0

BP=(NPFY/NP FX) \*\*2 +1.0

BP=(NPFY/NP FX) \*\*2 +1.0

CP=2.0\*NPFY\*NPFZ/NP FX\*\*2

AFP=BP=(C\*BP) \*\*2-CP\*AA/BB

BPP=2.0\*AA\*CC\*BP/BB\*\*2-CP\*CC/BB

CFP=BP\*(CC/BB) \*\*2-1.0

CLP=BPP/(2.0\*APP) +S 1GN\*SQRT(BPP\*\*2-4.0\*APP)/(2.0\*APP)

CKP=(-NPFY\*CLP-NPFZ\*CMP)/NPFX

RETURN

END SIGN, THETA, R, THETAP, FILTER)
EPSLON, PIZ, J
I, N POX, NP OY, NP OZ, P SIP
OS YMB, YO, ZO, RO, ČKKP, CLLP, CMMP, NZ O, ALPHA THETA, RADG, ROOTG, FILTER I) RETURN PETURN RADH, ROOTS, SIGN, FILTER I) RETURN PERFORMS ITERATION TO F THE GRIN SKEW-RAY AN NEWTON-RAPHSON ITERAT 0 LOOP=0 CALL RADYUS (RO,THETA,RAD IF (FILTER . EQ. 1) RETURN CALL SURF (THETA,RADH,RD) IF (FILTER . EQ. 3) RETURN FXDOT=ROOTG-ROOTS XNFW=THETA-FXDOTS XNFW=THETA-FXDOTS XNFW=THETA-FXDOTO IF (DIFF . GT . 0.00 001) GO RETURN C CONTINUE SOCOE SUBROUTINE ITRATE (COMMON A B, E, RZER CCOMMON/DIR/AI, BI, CCOMMON/SUR/PSI 0, X O INTEGER LOOP, FILT ER шO INE ITRATE INTERCEPT OUSING THE SUBROUTI OF THE IN SURFACE U 10 S

GISI1080 GIS11090 GIS11100 GIS11110 GIS11120

IF(LOOP .GE. 90) FILTER=3
hRITE(6,641C)
6410 FORMAT(/1x, 'SOLUTICN NOT FOUND')
RETURN
END

## APPENDIX C

## SAMPLE TABULAR OUTPUT FROM PROGRAM GISL

= 0.003220.9795850 0.9430425 0.9430423 0.2992280 0.4861692 0.1072993 0. £780985 0.1776611 0.8777914 0.8304005 C. 1028165 0.9258599 1.0022020 0.9676142 0.9676141 0.2088804 0.4765170 0.1039242 0.8258599 1.0022020 0.1039242 0.9168330 0.9946885 0.9594508 0.9594507 0.3056630 0.4797347 0.1050406 0.9168330 0.4797347 0.1050406 0.9616396 1.0320053 1.0000000 0.9999998 0.3217505 0.4636472 C.0955450 0.6858528 0.1720029 0.8769178 0.8318774 0.9438142 1.0171518 0.9838630 0.9838598 0.3153155 0.4700826 0.1017177 0.943814511 0.9527428 1.0245905 0.9919431 0.9919429 0.3185330 0.4668645 0.1006271 0.6847707 0.1727901 0.8770413 0.8316687 DEL TA LOP = 1.00000 0.9348540 1.0096884 0.9757503 0.9757501 0.3120980 0.4732593 0.6825813 0.1743850 0.8772871 0.8312529 0.9871492 0.9512604 0.9512602 0.3024455 0.4829524 0. 6792321 0.1768308 0.8776608 0.8306212 12P FUCAL LENGTH FROM STATION ZERO = 4.000C0 INDICES OF REFRACTION: N1 = 0.3217505 RADIUS = 1.00000 TABULAR OUTPUT I TERATIONS X 2 GISL EDGE THICKNESS = 0.0500000 BETA 3 ALPHA = 0.7853982 7 11 INCIDENT ANGLE = 1.00000I LENS FARAMETERS 0.8986831 Z 2 æ



0.9719930 0.9347960 0.9347959 0.2960104 0.4893869 0.1084422 0.6769556 0.1784965 0.8779076 0.8302040 0.1095940 11 0.8712009 0.9567289 0.9182188 0.9182186 0.285555 0.4958213 0.1107548 0.8781521 0.8297514 0.8627660 0.8162262 0.8162261 0.2509654 0.5344318 0.1254203 0.6599776 0.1909401 0.8793746 0.8277252 24 0.7482984 0.8547372 0.8075203 0.8075202 0.2477480 0.5376495 0.12€7056 0.7482984 0.8276040 0.1191421 0.8945618 0.8507230 0.8507228 0.2638354 0.5215619 0.12037£1 0.6650197 0.187244£ 0.8790098 0.8283439 0.9643753 0.9265221 0.5265219 0.2927930 0.4926046 0.6758038 0.1793403 0.8780314 0.8299952 14 0.8433996 0.9336274 0.8931330 0.8531329 C.2759229 0.5054744 0.8433996 0.6711047 0.1827840 0.8784947 0.8292129 0.9024346 0.8592664 0.8592662 0.2670529 0.5183452 0.6662558 0.1863384 0.8789114 0.8285100 0.9258732 0.8847126 0.8847124 0.2767054 0.5086521 0.6699068 0.1836621 0.8786031 0.8290300 0.9102768 0.8677794 0.8677792 0.2702704 0.5151275 0.6674823 0.1854381 0.8788039 0.8286913 0.8866592 0.8421465 0.8421463 0.2606179 0.5247796 0.6637738 0.1881607 0.8791218 0.8281551 0.8787258 0.8335392 0.8335391 0.2574005 0.5279973 0.6625181 0.1890816 0.8792174 0.8279540 0.7676280 0.8707618 0.8248989 0.8248987 0.2541829 0.5312141 0.6612529 0.1900077 0.8792992 0.8278561 0.9490560 0.9098864 0.9098862 0.2862579 0.4990390 0.6734728 0.1810462 0.8782556 0.8296165 0.91808 98 0.8762615 0.8762614 0.2734879 0.5119698 0.687879 0.119698 0.9413573 0.9015246 0.9015244 0.2831404 0.5022567 0.6722935 0.1819105 0.8783699 0.8294234 0.7579840 0.8619708 15 0.7963278 0.7867595 0.7772329 9 0.8895560 0.8527024 0.8340593 18 0.8058172 16 0.8246826 17 0.8152681 10 0.8803961 <del>:</del>3

38 0.6082788 0.7387393 0.6819450 0.6819448 0.2027029 0.5826944 0.1457528 0.6082788 0.6396450 0.2056431 0.8796603 0.8272479 0.1332816 0.6491598 0.7726042 0.7185560 0.7185559 0.2155729 0.5698245 0.1401095 0.6452883 0.2016122 0.8798180 0.8269823 3C 0.6893276 0.8058743 0.7545682 0.7545680 0.2284429 0.5569547 0.1346271 0.8270478 0.7976097 0.7456194 0.7456192 0.2252254 0.5601724 0.1359826 0. 6494152 0.1986350 C.8798047 0.8270C46 0.6592678 0.7809767 0.7276143 0.7276142 0.2187904 0.5666068 0.1387237 0.6592678 0.8466741 0.200613C 0.8798138 0.8269892 0.6390076 0.7641952 0.7094607 0.7094606 0.2123554 0.5730422 0.1415054 0.6390076 0.6438924 0.2026151 0.8798357 0.8270029 0.1471885 0.6693305 0.7893112 0.7366352 0.7366350 0.2220075 0.5633891 0.1373481 0.6693305 0.6480498 0.1373481 27 0.7189997 0.8304576 0.7811996 0.7811995 0.2380954 0.5473C16 0.77189997 0.8272754 0.7301762 0.6726955 0.6726953 C.1994854 0.5859120 0.6382093 0.2066600 0.8795860 0.8273730 0.7385736 0.8466769 0.7987810 0.7987808 0.2445304 0.5408671 0.7385736 0.8274904 0.8222976 0.7723581 0.7723579 0.2348775 0.5505193 0.6534515 0.1957040 0.8796960 0.8271877 0.6288099 0.7557480 0.7003270 0.7003269 0.2051379 0.5762599 0.0642865 0.2036191 0.8797619 0.8270767 0.7288064 0.8385835 0.7900075 0.7900074 0.2413129 0.5440835 0.7288064 0.8873380 0.6992603 0.8141024 0.7634806 0.7634804 0.2316604 0.5537370 0.692603 0.8170944 0.6185674 0.7472628 0.6911556 0.6911554 0.2C59204 0.5794767 0.6410708 0.2046333 0.879413 0.8271113 0.7091515 0.6793501

40 0.5875602 0.7215751 0.6634064 0.6634063 0.1562678 0.5891297 0.1486340 0.58755260

0.1500890 0.6955217 0.6352974 0.6352973 0.1866153 0.5987819 0.1530283 0.6323695 0.2107540 0.8791294 0.8281425 0.1605376 0.5965770 0.5289415 0.5289414 0.1512228 0.6341746 0.1698255 0.6155723 0.2219762 0.8759543 0.8335C92 0.1729792 0.6332145 0.5682372 0.5682371 0.1640528 0.6213048 0.1636023 0.6217955 0.2179340 0.8774970 0.8308581 0.6690865 0.6068150 0.6068149 0.1769629 0.6084350 0.1575071 0. 6278507 0.2138370 0.8785779 0.8290725 0.7042491 0.6447080 0.6447079 0.1858329 0.5955642 0.151553 0.7129317 0.6540776 0.6540775 0.1930503 0.5923465 0.6353088 0.2087043 0.8794014 0.8276839 0.6058107 0.5388345 0.5288344 0.1544403 0.6309569 0.60581071392 0.2209763 0.8764024 0.8327499 0.5779552 0.5090140 0.5090139 0.1447878 0.64061C0 0.6124186 0.2239618 0.8749874 0.8351490 0.6779414 0.6163516 0.6163515 0.1801803 0.6052173 0.6293528 0.2128095 0.8787854 0.8287225 0.6241291 0.5584818 0.5584816 0.1608753 0.6245225 0.6241291 0.2189507 0.8771533 0.8314753 0.6867517 0.6258453 0.6258451 0.1833978 0.6019596 0.6308858 0.2117817 0.8789694 0.8284122 0.66 01877 0.5972365 0.5972363 0.1737453 0.6116517 0.6263800 0.2148613 0.8783325 0.8294867 0.6512430 0.5876139 0.5876138 0.1705278 0.6148694 0.6248602 0.2158903 0.8780947 0.8298882 0.6422523 0.5779477 0.5779476 0.1673103 0.6180871 0.6223320 0.2169136 0.8778104 0.8303685 0.6149944 0.5486809 0.5486808 0.1576578 0.6277393 0.6186990 0.2199641 0.8767840 0.8321041 0.5872923 0.5190016 0.5190015 0.1480053 0.6373523 0.6139987 0.2229703 0.8754758 0.8343204 46 0.5242630 45 0.4519505 0.5771319 0.5666530 0.5561299 0.5349353 52 0.4591874 0.4481649 56 0.4147875 44 0.5455577 51 0.4701589 54 0.4370912 50 0.4810801 0.5027711 55 0.4259651 48

0.3466337 0.5207836 0.4480610 0.4480609 0.1254827 0.6599143 0.1825695 0.6028283 0.2296700 0.8711029 0.8417643 0.5685652 0.4989785 0.4989784 0.1415703 0.6438268 0.1745651 0.6108327 0.2249327 0.8744001 0.8361464 0.3809454 0.5496220 0.4787614 0.4787613 0.1351352 0.6502622 0.1777536 0.3809454 0.6076442 0.2268587 0.8731809 0.8382204 60 0.3695607 0.5400661 0.4685786 0.4685785 0.1319178 0.6534799 0.1793550 0.6060428 0.2278067 0.8725201 0.8393459 0.1890293 72 0.2288720 0.4206347 0.3422866 0.3422865 0.0933077 0.6920894 C.1987091 0.5866887 0.2376738 0.8601224 0.8607012 64 0.3234909 0.5012654 0.4273381 0.4273379 0.1190478 0.6663497 0.1857961 0.5996017 0.2315134 C.8697081 0.8441503 0.2883955 0.4715261 0.3958564 0.3958563 0.1093952 0.6760018 0.190646 0.192263 0.4614829 0.3852540 0.3852539 0.1061777 0.6792195 0.65931340 0.2347438 0.8651602 0.8519691 70 0.2524339 0.4411972 0.3638837 0.3638836 0.0997427 0.6856549 0.0907427 0.6856549 0.0907427 0.88561365 0.3581234 0.5304537 0.4583452 0.4583451 0.1287003 0.6566976 0.2287461 0.8718375 C.8405C95 0.4309509 0.3531137 0.3531135 0.0965252 0.6868726 0.5882948 0.2369844 0.8614810 0.8583391 0.3001441 0.4815053 0.4064043 0.4064041 0.1126127 0.6727841 0.5963685 0.2330918 0.8671477 0.8485447 0.3922781 0.55912 06 0.4888944 0.4888943 0.1383528 0.6470445 0.3922781 0.55912 0.2259064 0.8738380 0.8371021 0.3118377 0.49142 32 0.4168978 0.4168977 0.1158302 0.6695674 0.5918837 0.5979854 0.2321913 0.8678473 0.8473423 0.2647396 0.4513741 0.3745971 0.3745970 0.1C256C2 0.6824372 0.25647396 0.5915185 0.2355235 0.8640264 0.8539279 0.4035584 0.2765932 0.2408782

Contract of

87 0.0431264 0.2564759 0.1725258 0.1725258 0.0450451 0.7403520 0.2215379 0.0431264 0.5638599 0.2431734 0.8303474 0.9139132 85 0.0684555 0.2794743 0.1960034 0.1960033 0.0514801 0.7339175 0.2187108 0.55 0.0684555 0.5666870 0.2430628 0.8353659 0.9047449 88 0.0304019 0.2448380 0.1606832 0.1606832 0.0418276 0.7435697 0.2229148 0.5624831 0.2431445 0.8277161 0.9187540 0.1803538 0.2786361 0.2983926 0.2583925 0.0804377 0.7049602 0.2050733 0.1803538 0.5803245 0.2400801 0.8538936 0.8716030 77 0.1681040 0.3679463 0.2872689 0.2872688 0.0772201 0.7081769 0.2066433 0.1681040 0.35787545 0.2405922 0.8521645 0.8746502 78 0.1558043 0.3571758 0.2760841 0.2760840 0.0740027 0.7113946 0.2082030 0.5771948 0.2410615 0.8503491 0.8778595 79 0.1434576 0.2463323 0.2648367 0.2648366 0.0707852 0.7146123 0.2097508 0.1434576 0.5756471 0.2414904 0.8484600 0.8812100 0.1310653 0.3354020 0.2535262 0.2535262 0.0675676 0.7178300 0.2112853 0.1310653 0.5741125 0.2418711 0.8464741 0.8847443 0.1186289 0.3243900 0.2421523 0.2421522 0.0643501 0.7210477 0.2128052 0.1186289 0.8243900 0.2422076 0.8444154 0.8884209 0.0936257 0.3021082 0.2192101 0.2192100 0.0575152 0.7274821 0.2157956 0.0936257 0.3021082 0.2157956 0.2047106 0.3997850 0.3204582 0.3204581 0.0868727 0.6985248 0.2019055 0.2047106 0.3997850 0.2389493 0.8571633 0.8658658 0.1925563 0.3892482 0.3094552 0.3094551 0.0836552 0.7017425 0.2034936 0.1925563 0.5819042 0.2395362 0.8555795 0.8686407 0.2168158 0.4102467 0.3314016 0.3314016 0.090~902 0.6953071 0.2003103 0.5168158 0.5850875 0.2383270 0.8586721 0.8632289 0.1061490 0.3132926 0.2307140 0.2307140 0.0611327 0.7242644 0.1061490 0.313292517 84 0.0810609 0.2908359 0.2076402 0.2076401 0.0546576 0.7306598 0.5681344 0.2429250 0.8377470 0.9004240 86 0.0558102 0.2680214 0.1842988 0.1842987 0.0482626 0.7371352 0.5652616 0.2431462 0.8329000 0.9092394 74 82

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0.2242649 0.2255866 0.2293640 0.2377607 0.2305551 0.2338991 .2093453 0.1247303 0.1247303 0.0321751 0.7532218 0.5585199 0.2427126 0.8193703 0.9342635 0.1729540 0.0881205 0.0881204 0.0225226 0.7628749 0.5548428 0.2417436 0.8103788 0.9512465 0.1356313 0.0508290 0.0508290 0.0128701 0.7725270 0.5514987 0.2402332 0.8008499 0.9695645 0.1481775 0.0633368 0.0633368 0.0160876 0.7693103 0.5525734 0.2407974 0.8040817 0.9633141 0.1229768 0.0382423 0.0382423 0.0096526 0.7757447 0.5504659 0.2396098 0.7975732 0.9759415 0.08436 C5 0.0000002 0.0000002 0.0000001 0.7853969 0.5476371 0.2373897 0.7875214 0.9957625 .2212741 0.1367862 0.1367862 0.0352526 0.7500651 0.5598112 0.2429150 0.8222255 0.9289307 0.1973162 0.1126012 0.1126012 0.0285576 0.7564395 0.5572602 0.2424498 0.8164404 0.5397660 •1851857 0 • 1003983 0 • 1603982 0 • 6257401 0 • 7596572 0 • 5560338 0 • 2421274 0 • 8134441 0 • 9454245 0.1606186 0.0757670 0.0757669 0.C153051 0.7660526 0.5536888 0.2413010 0.8072599 0.9572054 0.2331046 0.1487704 0.1487703 0.0386101 0.7467874 0.5611330 0.2430580 0.8250054 0.5237653 0.1102147 0.0255759 0.0255759 0.0064350 0.7789624 0.0545469 0.2389274 0.7942556 0.9824401 0.0973431 0.0128291 0.0128291 0.0032176 0.7821801 0. 5485333 0.2381868 0.7909024 0.9890516 0 0 0 0.0176410 51-0.0079843 0699940.0-46 95-0.0596255 58-0.0986636 99-0-1117308 100-0.1248229 101-0.1379396 96-0.0726097 97-0.0856237 0.0048451 52-0.0208471 93-0-0337411

II V

NCSE HALF ANGLE = 0.78327 STATION I OPAQUE SURFACE LENGTH = 0.11956 TRANSPARENT SURFACE LENGTH =1.45205 TCTAL LENS SURFACE LENGTH =1.57160

NOSE SECTION DATA:

	0.36591
ENSTINDEX OF REFRACTION DATA THE INDEX GRADIENT IS: N2(R) = SQRT((9.00)+(5.351)*R**2) WITH THE CENTER OF SYMMETRY LOCATED AT 0.050 (PINUS IMPLIES THE OSYM IS ON THE IMAGE SIDE).	IT CHANGE OF THE LENS INDEX GRADIENT: DE SURFACE FROM CENTER TO FDGE: 49.88858 DE SURFACE TO OLISIDE SURFACE AT THICKEST PCINT: IDE SURFACE FROM CENTER TO EDGE: 49.34212

15KE BAY TRACE PARAMETERS:

	12	YDIAPT	-2210370 -0.137F+01	-5424452 -0.111E+01	-8703845 -0.103E+01	0.0 0.100F+06	0.2314239 0.368F+01	-5320314 -0.183E+01	-8564321 -0.127F+01	0.0
ERS ABOVE.	1 \	XDI AP T	0.1287202 0 .126E+01	0.2018253 0 .775E+00	0.4872602 0 .433E+C0	0.1305212 0.100E+06	0.1139907 0 .588E÷01	0.0265410 0.442E+00	0.1123655 0 .162E+00	0.3610038 0
SEE LENS PARAMETERS ABOVE.	1 X	NTNCTY	0.30000 0.1307886-0.1494352 0.2557881-0.1287202 0.2210370 5.04583168 -2.1801453-0.1453156 0.5102614 -0.126E+01 -0.137F+01	0.60000 c0 0.5384317-0.3217840 0.6207671-0.2018253 0.5424452 4.45545169 -1.4455128-0.2271336 0.4481558 -0.775E+00 -0.111E+01	0.50000 CO 0.5583030-0.4993026 0.9974844-0.4872602 0.8703845 4.09892273 -1.1834154-0.2429521 0.3635889 -0.433E+CO -0.103E+O1	4 C.3C00000 0.0 0.0109681 0.2269354 0.1305237 0.1305212 4.85042477 -1.6185007 0.0 0.5535472 0.100E+06	0.3000000 0.1424869 0.1713301 0.2579772 0.1139907 0.2314239 59103 -1.9860687 0.3676555 0.5334762 -0.588E+01 0.368F+01	C.3 CCO000 0.6000000 0.4431463 0.0442133 0.5327001 0.0265410 0.5320314 4.60996151 -1.7169704 0.0316657 0.4739231 0.442E+00 -0.183E+01	7 C.3 CC0000 0.90000 00 8111566-C.1113788 0.8637780-0.1123655 0.8564321 4.27614594 -1.4330120-0.1158373 0.3925338 -0.162E+00 -0.127F+01	0.2600275 0.4473460 0.3610061 0.3610038
OOO SEE LI	γO	N WIZ	-0.1494352 453156 0.53	-0.3217840 271336 C.44	-0.4993026	0.2269354	9.1713301 576555 0.53	0.0442133	-C.1113788	0.4473460
0000€=0	ÛX	WIA	1307886- 1453-0.1	5128-0.2	9583030- 14154-0-24	0.0109681 15007 0.0	0.1424869	9704 0.03	8111566- 10120-0-11	0. 26002 75
C00000 GRID= 0.3000000	RAY Z	7	.30000 c0 (168 -2.180	.6000000 169 -1 .445	. 50000 CO (273 -1 .183	.0 477 -1.618	300000000	600000000000000000000000000000000000000	9000000 594 -1 -433	0.0
ALFAP= 0.4C0	RAYY	OPL				.3C00000 0 4.85042	C.3 C00000 0 0 4 0 9 3 1 5 9 9	36600000 e.	.3 0C0000 0 4.27614	E 0.6000000 0
ALI	RAY		1 0.0	2 0.0	3 0.0	4	2	, j	, C	ů W

0.100E+06 0.0 0.100E+06 16 1.1999998 0.0 0.7783279 0.8796349 0.8306387 0.8306248 0.0 0.100E+06 4.496C8326 -1.3109426 0.0 0.3974670 0.100E+06 0.8597310 -0.150E+01 13 0.9000000 0.3000000 0.5713329 0.6414395 0.6435080 0.5810593 0.2765256 4.53502178 -1.4060097 0.1153301 0.4473754 0.262E+01 -0.283E+01 0.5661314 -0.205E+01 17 1.1999998 0.30000 0.8187981 0.8625243 0.8675712 0.6187382 0.2869341 4.47258854 -1.3300571 0.0765226 0.3879091 0.136E+01 -0.211E+01 15-C.3000000 0.0 0.4567428-0.6129578 0.5509578-0.5509512 C.0 0.100E+06 4.38281155 -0.5380237 0.0 0.4409877 0.100E+06 0.100E+06 20-0.3C00000 0.3000000 0.6123005-0.6787266 0.6885747-0.6313385 0.2748162 4.28067303 -0.9621467-0.2031004 0.4222168 -0.118E+01 -0.822E+00 21-0.3000000 0.6000000 0.9397941-0.8171884 0.9814265-0.7927174 C.5786383 4.05062485 -0.5386383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436383 -0.8436840 -0.8436884 -0.843684 -0.844684 -0.844684 -0.844684 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84468 -0.84488 -0.84488 -0.84488 -0.84488 -0.84888 -0 9 0.6000000 0.300006C 0.3363692 0.4150693 0.4317581 0.3448501 0.2557798 4.65313148 -1.5641460 0.1932485 0.5009179 0.166E+02 -0.711E+01 10 C.6000000 0.6000000 0.5420605 0.3281044 0.6192006 0.2887600 0.5477414 4.5335353596 -1.5893183 0.1321886 0.4537812 0.130E+01 -0.219E+01 15 0.9000000 0.9000000 0.9483795 0.4820271 0.5865439 0.4622759 0.8715199 4.30145073 -1.4332008 0.0648745 0.3594230 0.465E+00 -0.159E+01 C.58C7928 -0.185E+01 18 1.1599998 0.6000000 0.9357986 0.8130573 0.9743750 0.7823852 4.40475750 -1.3604498 0.1072606 0.3611534 0.908E+00 11 0.6000000 0.9000000 0.8309218 0.2059760 C.8806287 0.1906774 4.31594181 -1.4822788 0.0079514 0.3868141 0.252E+00 12 0.9000000 0.0 0.5948853 0.6639546 0.5948910 0.5948853 4.56061363 -1.3724556 0.0 0.4602101 0.100E+06 14 0.9000000 0.6000000 0.7219925 0.5777416 0.7806184 0.5374358 4.452C9980 -1.4449883 0.1326690 0.4111363 0.121E+01 0. 1COE+C6 0.5169492 4.66151619 -1.4761887 0.0

MIRRCR IMAGE SKEW RAYS FULLOW

IMAGE RAY	E RAY	r RAYY	RAYZ	OPL	MIZ WIA		NTNCTY
22	<b>~</b>	1 0.0	-0.3 00000	5.04583168	5.04583168 -2.1801453 C.1453156	156 0.	0.5102614
23	2	0.0	-0° 6 C000CJ	4.45555169	4.45555169 -1.4455128 0.2271336		0.4481558
54	ĸ	0.0	0000006 •0-	4.09892273	4.09892273 -1.1834154 0.2429521		0.3635889
25	2	0.300000	0.30000 00-0.3 000000	4.93159163	-1.9860687-0.3676595		0.5334762
26	9	0.300000	0.3000000-0.6 (00000	4. 60996151	-1.7169704-0.0316697		0.4739231
27	7	0°30000C	0.3000000-0.90000000	4.27614594	-1.4330120 0.1158373		0.3925338
28	6		0.600000-0.3000000	4.65313148	-1.5641460-0.1922485		0.5009179
52	10		0.6000000-0.600000000000000000000000000	4.53353596	-1,5893183-0,1321886		0.4537812
30	11	0.60000 C	0.60000 00-0.900000000000000000000000000000	4.31594181	-1.4822788-0.0079514		0.3868141
31	13		C. 90000 co-0.3 co0000	4.53502178	-1.4060097-0.1153301		0.4473754
32	14		0.900000-0.600000000000	4.45209980	-1.4449883-0.1326690		0.4111363
33	15		0.900000-0.9000000-0	4.30145073	-1.4332008-0.0648745		0.3594230
34	17	1.199995	1.1999958-0.3 000000	4.47258854	-1.3300571-0.0765226		0.3879091
35	18	1.199595	1.1995558-0.6 000000	4.40475750	4.40475750 -1.3604498-0.1072606		0.3611534
36	20.	-0.300000	20-0.300000-0.3000000	4.28067303	4.28067303 -0.9621467 0.2031004	004 C.	C.4222168
37	21.	-0.30000	21-0.3000000-0.6000000	4.05062485	4.05062485 -0.9309335 0.2432328	328 0,	0.3676401
ENC	CF SI	CF SKEW RAY TRACE.	RACE.				

TOTAL NUMBER OF RAYS TRACED = 37

TOTAL NUMBER OF MIRR CR IMAGE RAYS = 16

TOTAL NUMBER OF RAYS STRIKING IMAGE PLANE = 3

## IMAGE PLANE SPOT CIAGRAM ANALYSIS:

THICKNESS =0.0500000 U =0.0	SS =	0.05	000000	ح	0 • 0=		ALFAP :	-0.40000	00	ALFAP =0.4000000 R =1.0000000	_
CENTROI	<b>7</b> : C	CENT	R = 0.0	• 0	YCENTR	11	CENTROID: 2CENTR = 0.0, YCENTR = -1.4490051	51			
STANDARE	) DE	VIAT	I CNS:	SI	GMAY =	o	0911858	SIGMAZ	14	STANDARD DEVIATIONS: SIGMAY = 0.0911858 SIGMAZ = 0.0264883	
PMS SPC	1 51	2E:	R MSR AD	11	PMS SPCT SIZE: RMSRAD = 0.3430366	99					

SPOT DIAGRAM ENERGY DENSITY VS. RADIUS FROM CENTRCID:

GG RADIUS FRACTIGN

POT DIAGRAM CALCULATIONS COMPLETE.

END OF PROGRAM

PRINTER" PLOTTER" "CHARTS PRINTER ENTER ENTER THE 90 OTTER F018 ೧೯ಡ CW ING TRIER CT TER TAIN PLOTS THE FOLLC FOR PRI FOR PLC F PROGRAM CBT **0**F END TOSI

## APPENDIX D



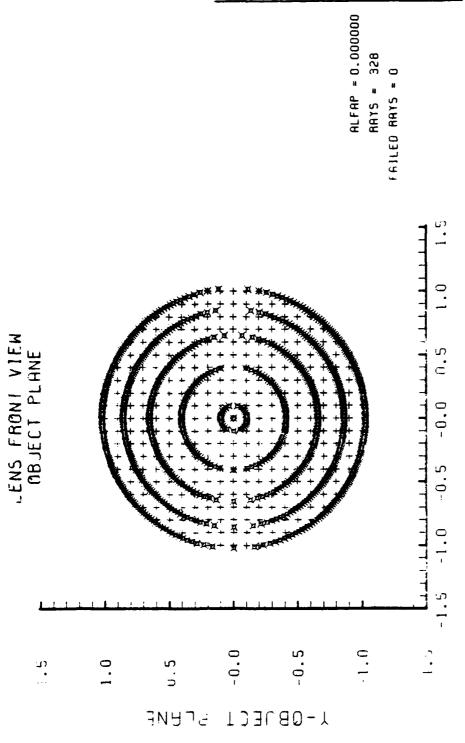


Figure D-1. Object Plane at  $\alpha_{\rm p}$  = 0.0 Radians for Hin Lens Design Shown in Figure 11

Z-OBJECT PLANE

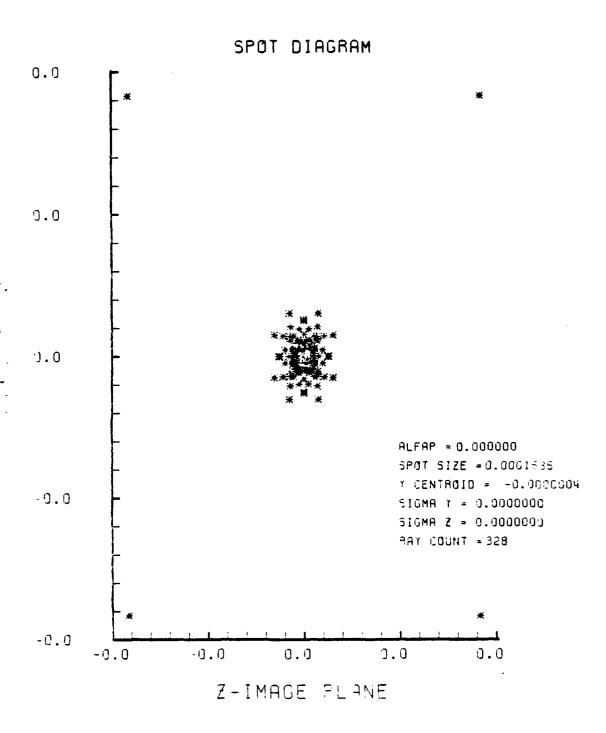


Figure D-2. Spot Diagram Corresponding to Object Plane of Figure D-1

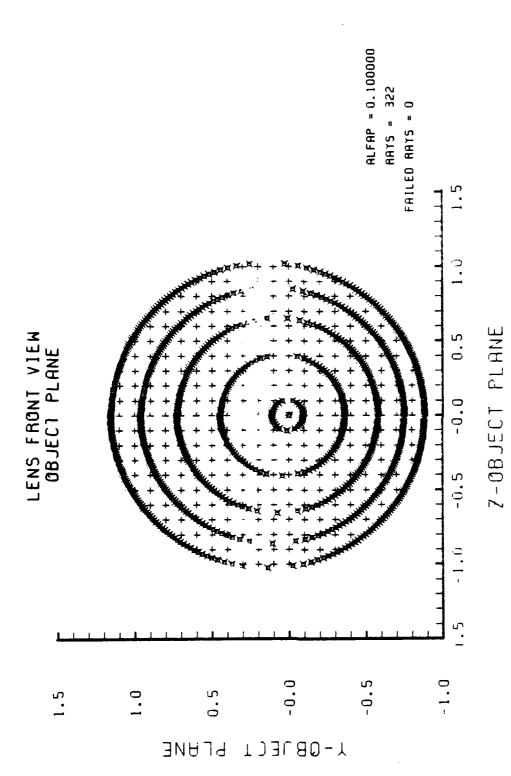


Figure D-3. Object Plane at  $\alpha_{\rm p}=0.1$  Radians for HIN Lens Shown in Figure 11

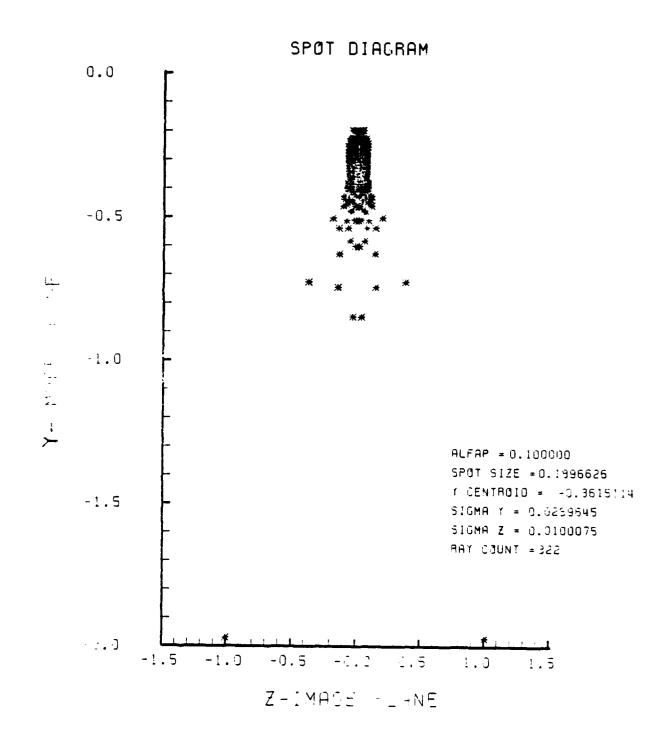


Figure D-4. Spot Diagram Corresponding to Figure D-3

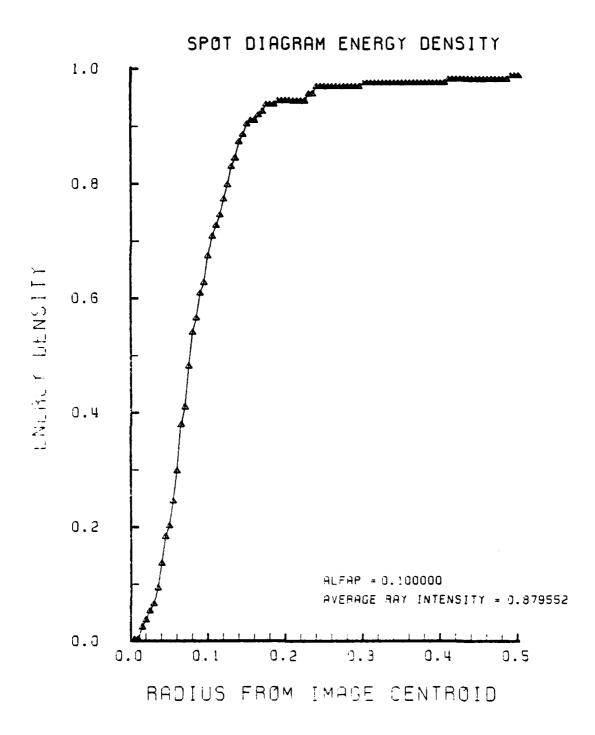


Figure D-5. Encircled Energy Plot for the Spot Diagram of Figure D-4

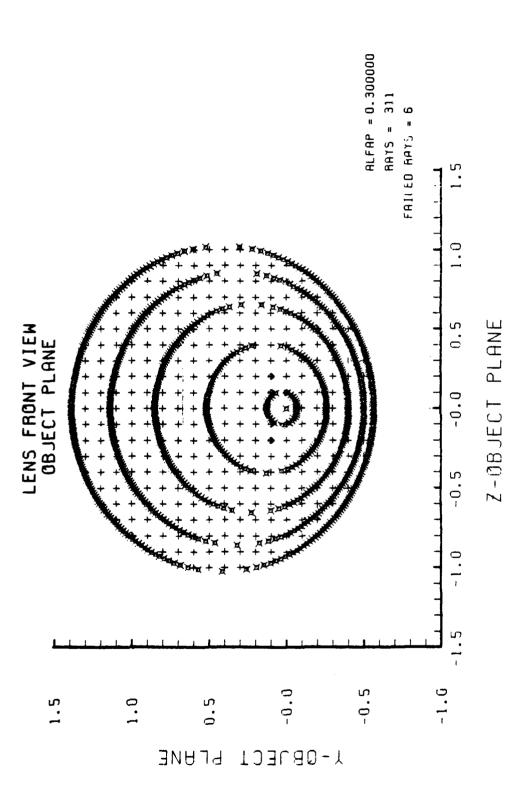


Figure D-6. HIN Object Plane at  $\alpha_p$  = 0.2 Radians,  $N_2$  = 1.5

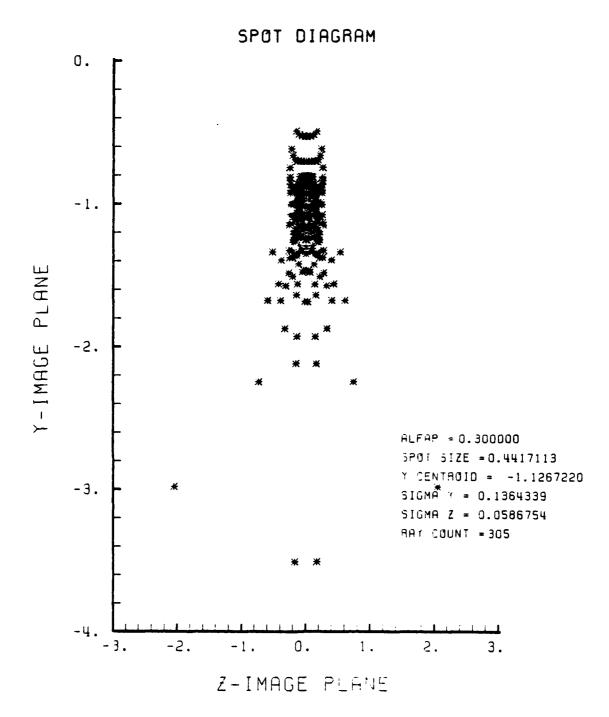


Figure D-7. Spot Diagram of HIN Lens at  $\alpha_p = 0.2$  Radians,  $N_2 = 1.5$ 

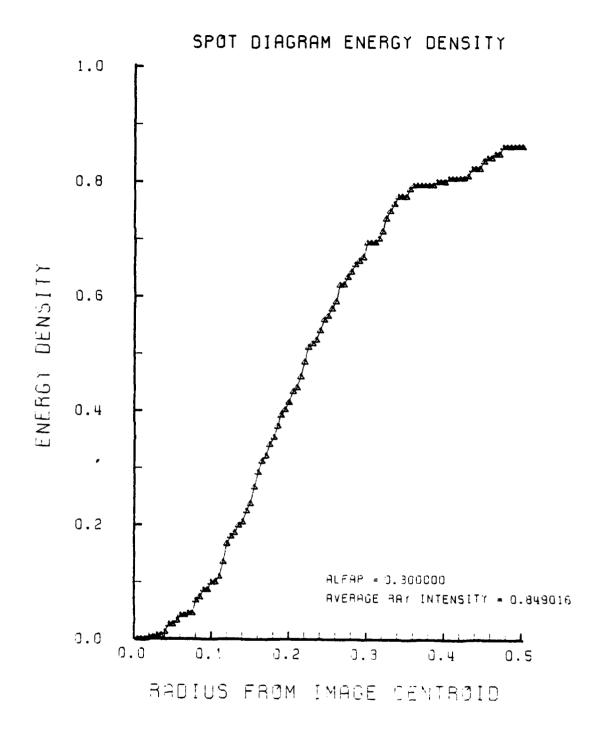


Figure D-8. Encircled Energy Plot of HIN Lens at  $\alpha_p$  = 0.2 Radians,  $N_2$  = 1.5

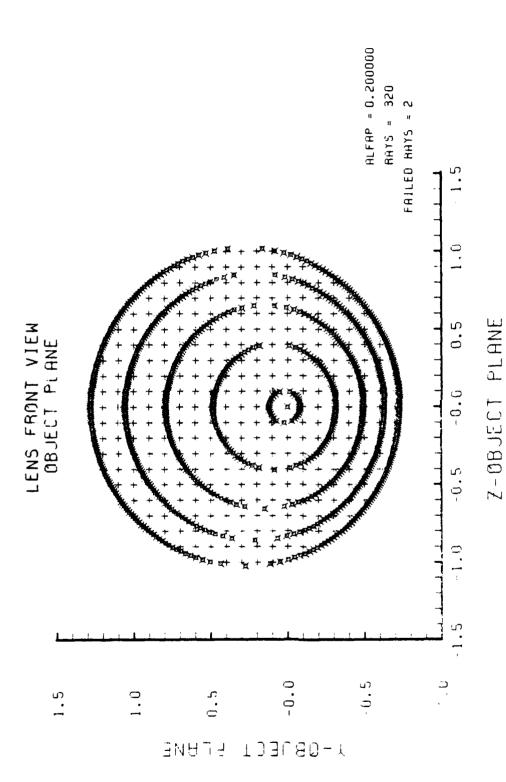
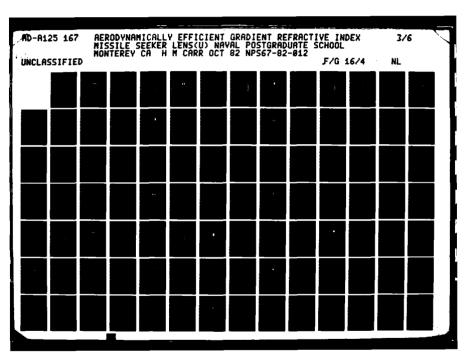
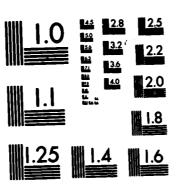


Figure D-9. HIN Object Plane at  $\alpha_p=0.3$  Radians,  $N_2=1.5$ 





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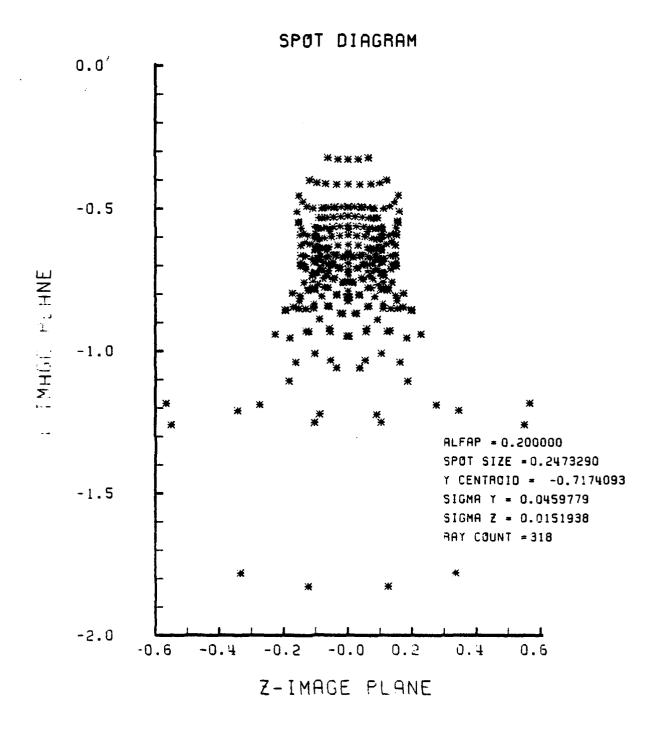


Figure D-10. Spot Diagram of HIN Lens at  $\alpha_p = 0.3$  Radians,  $N_2 = 1.5$ 

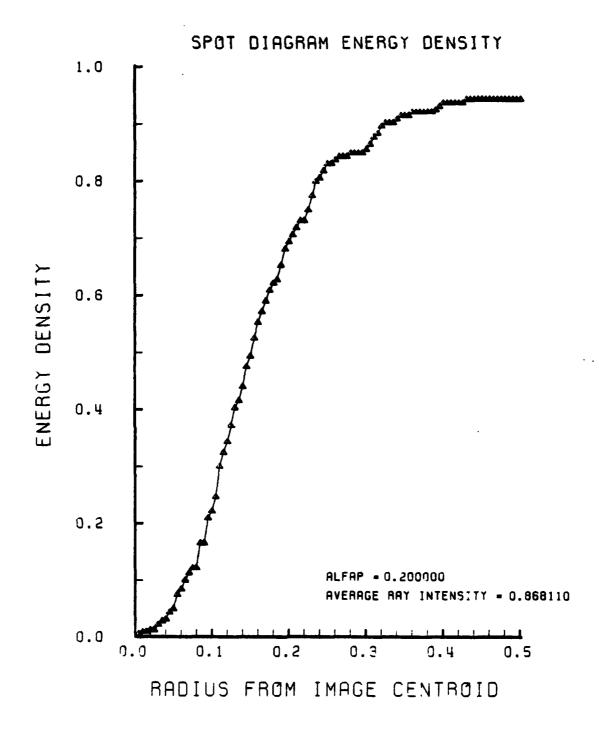


Figure D-11. HIN Lens Encircled Energy at  $\alpha_p = 0.3$  Radians,  $N_2 = 1.5$ 



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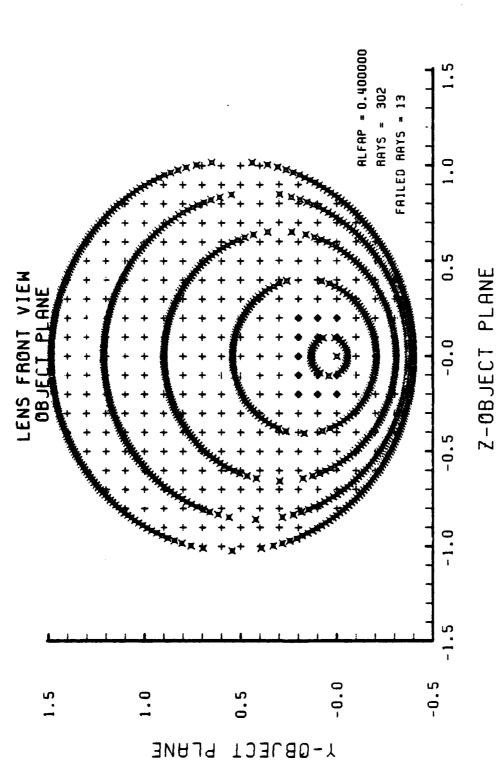
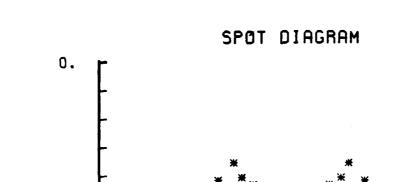


Figure D-12. HIN Object Plane at  $\alpha_p=0.4$  Radians,  $N_2=1.5$ 



-3.

\* \*

ALFAP = 0.400000

SPOT SIZE = 0.4655773

Y CENTROID = -1.5407810

SIGMA Y = 0.1617901

SIGMA Z = 0.0549725

AAY COUNT = 289

-u. \* \* \* \* -1.00-0.75-0.50-0.250.00 0.25 0.30 0.75 1.00

Z-IMAGE PLANE

Figure D-13. Spot Diagram of HIN lens at  $\alpha_p = 0.4$  Radians,  $N_2 = 1.5$ 

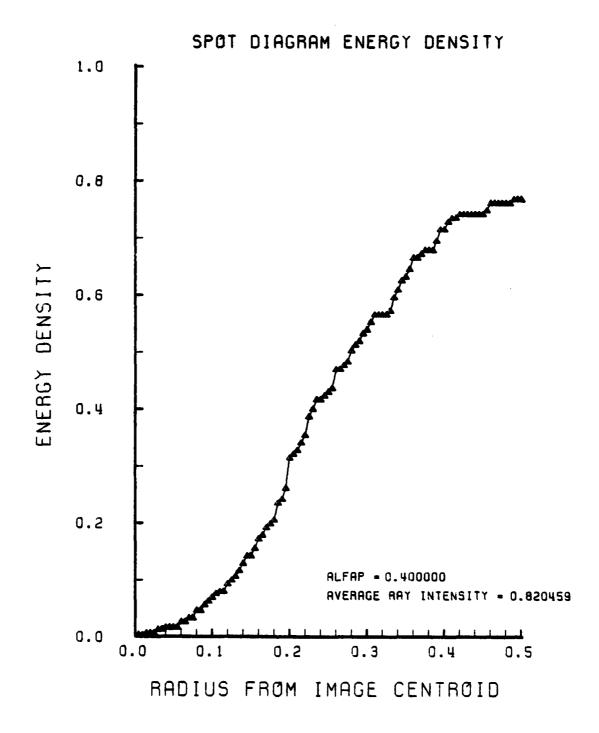
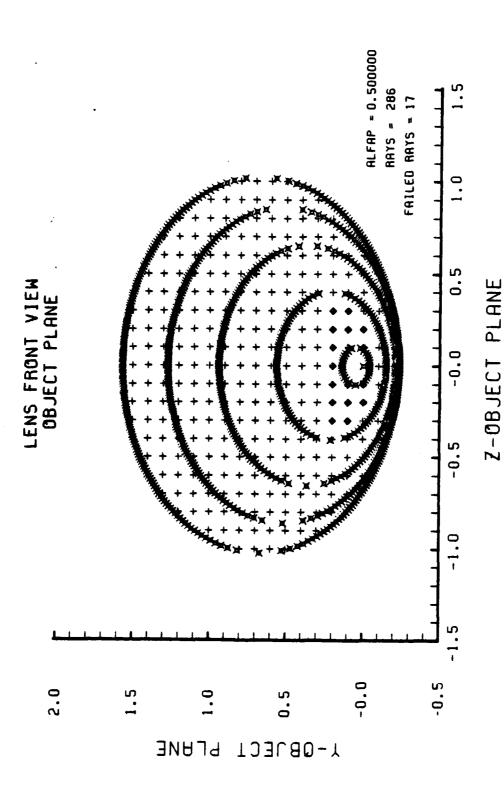


Figure D-14. Encircled Energy Plot at  $\alpha_p = 0.4$  Radians,  $N_2 = 1.5$ 

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HIN Lens Object Plane at  $\alpha_p$  = 0.5 Radians,  $N_2$  = 1.5 Figure D-15.



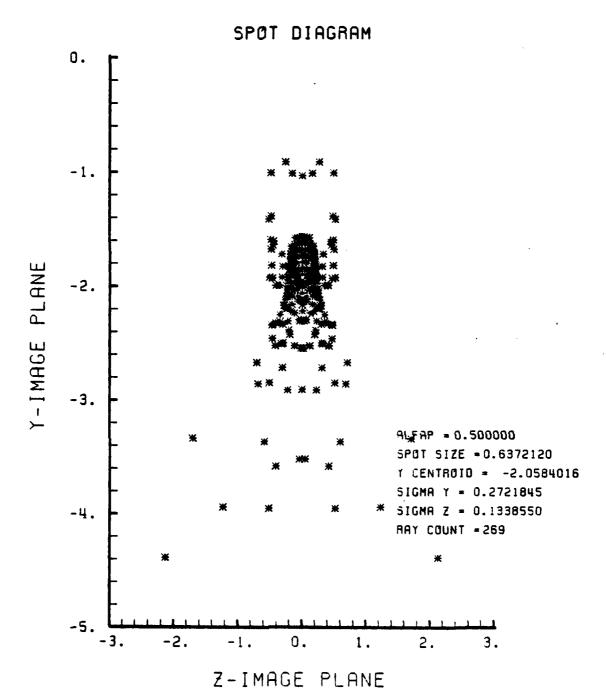


Figure D-16. Spot Diagram of HIN Lens at  $\alpha_p = 0.5$  Radians,  $N_2 = 1.5$ 

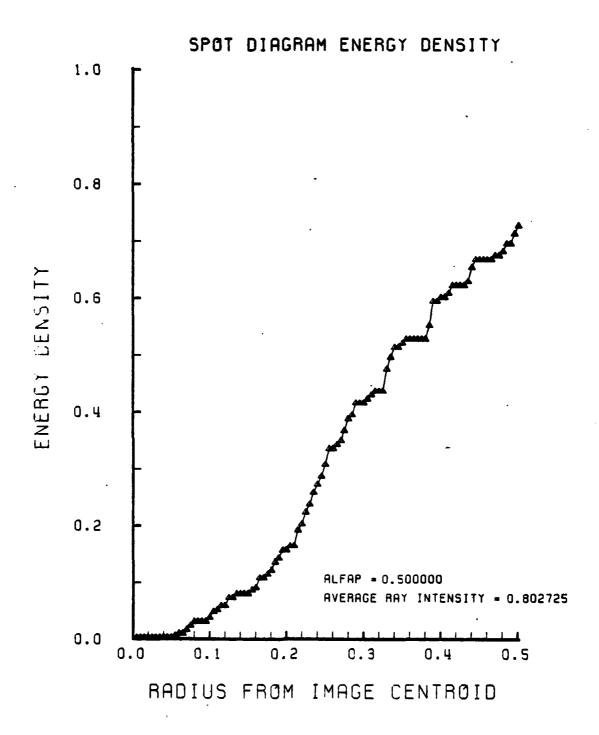
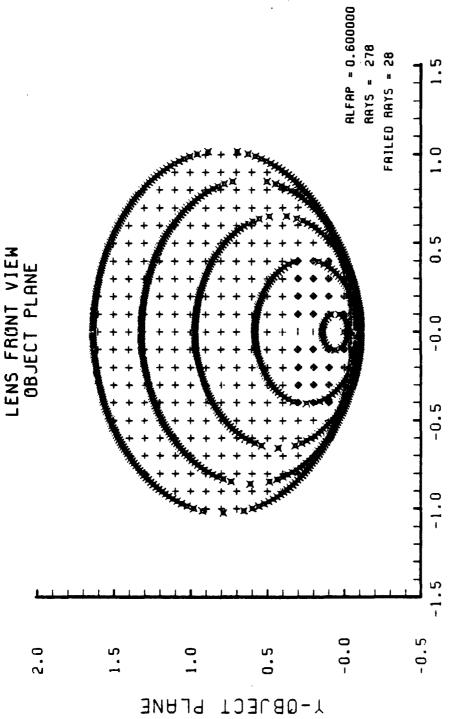


Figure D-17. Encircled Energy Plot at  $\alpha_p = 0.5$  Radians,  $N_2 = 1.5$ 



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HIN Object Plane at  $\alpha=0.6$  Radians,  $N_2$  Overlapping Portions Bf Inner Ellipses are hidden lines. Figure D-18.

Z-OBJECT PLANE

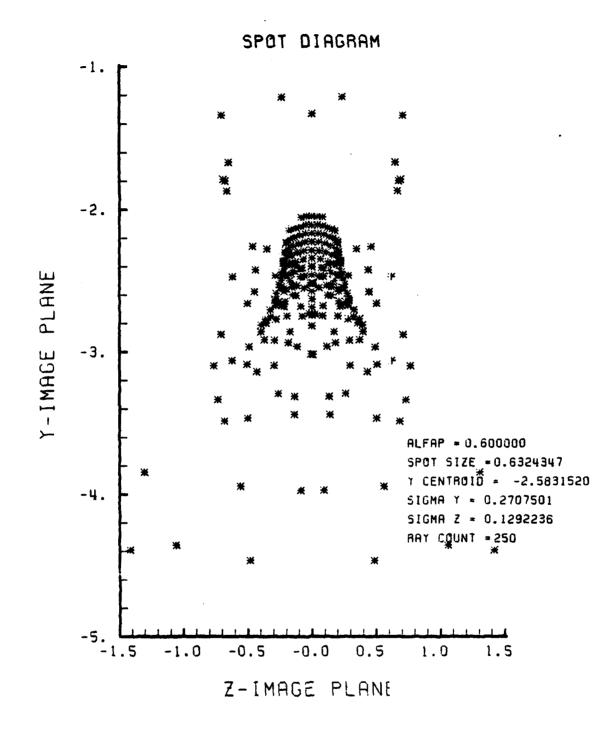


Figure D-19. Spot Diagram of HIN Lens at  $\alpha_p = 0.6$  Radians,  $N_2 = 1.5$ 

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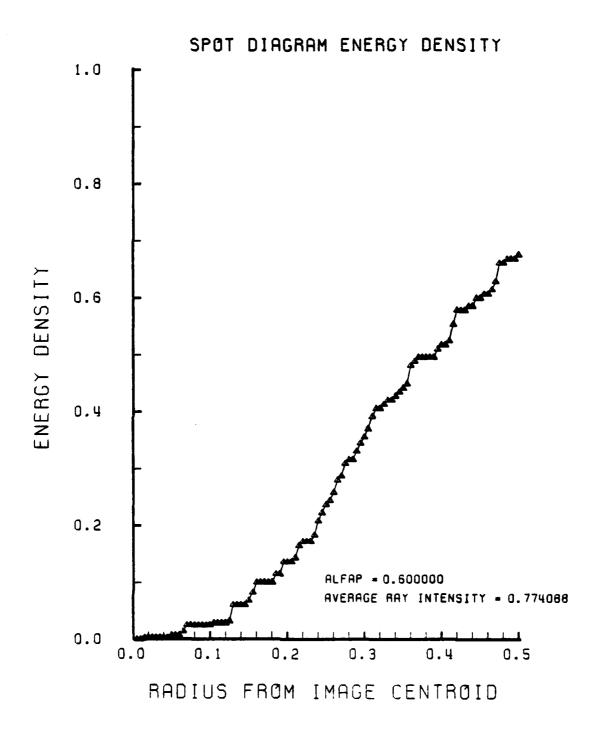
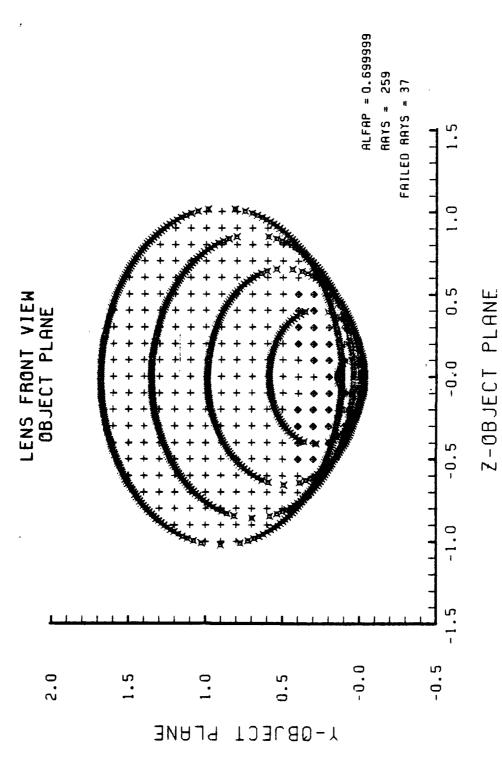


Figure D-20. Encircled Energy Plot at  $\alpha_p = 0.6$  Radians,  $N_2 = 1.5$ 





HIN Object Plane at  $\alpha=0.7$  Radians, N<sub>2</sub> Overlapping Portions Bf Inner Ellipses are Hidden Lines. Figure D-21.

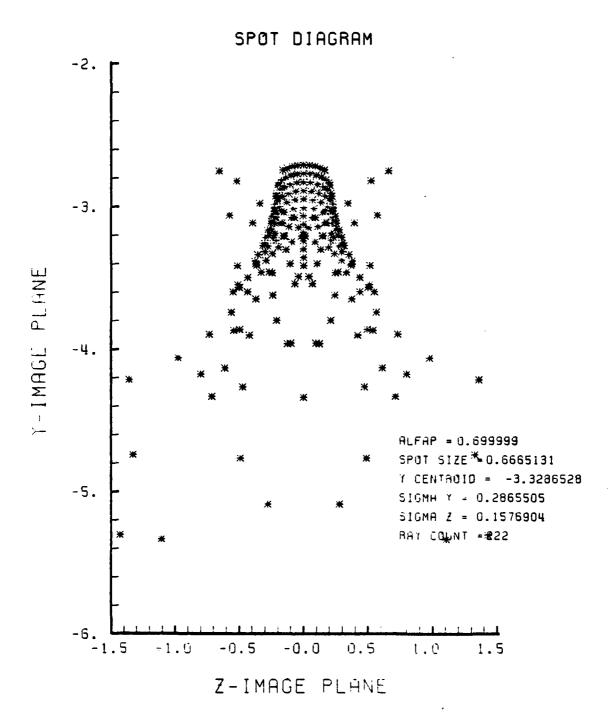


Figure D-22. Spot Diagram of HIN Lens at  $\alpha_p = 0.7$  Radians,  $N_2 = 1.5$ 

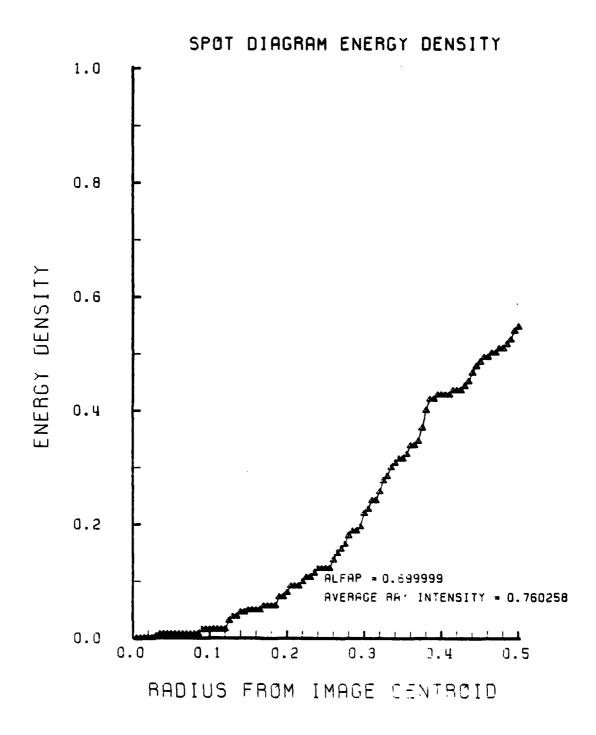


Figure D-23. Encircled Energy Plot at  $\alpha_p = 0.7$  Radians,  $N_2 = 1.5$ 

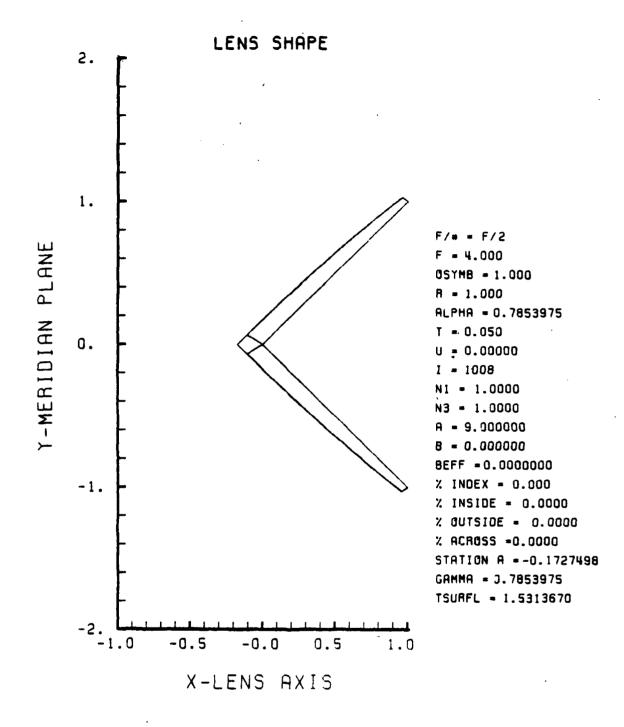


Figure D-24. HIN Lens Design for  $N_2 = 3.0$ 

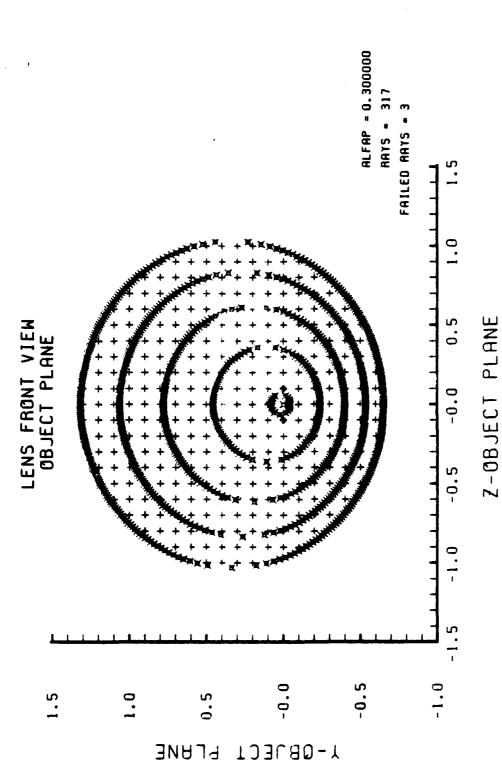


Figure D-25. Object Plane of HIN Lens in Figure D-24 at  $\alpha_p$  = 0.3 Radians

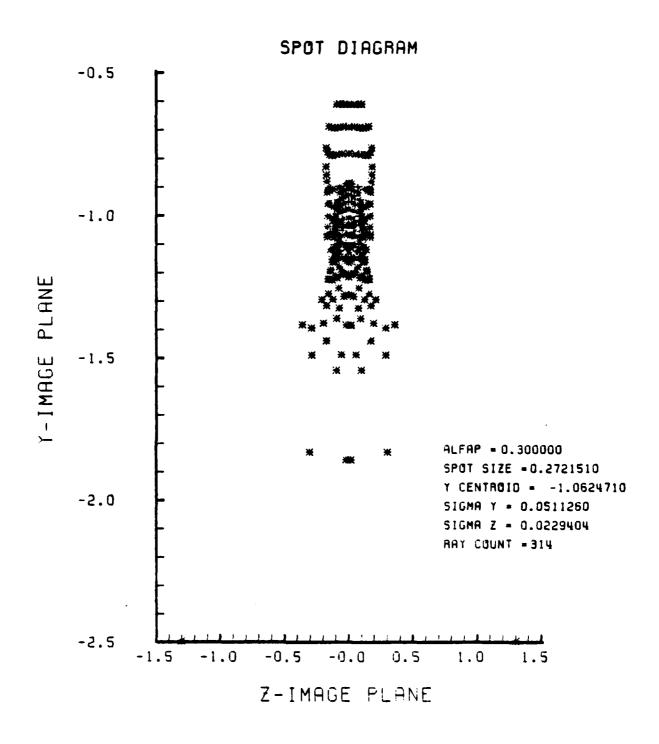


Figure D-26. Spot Diagram Corresponding to Figure D-25

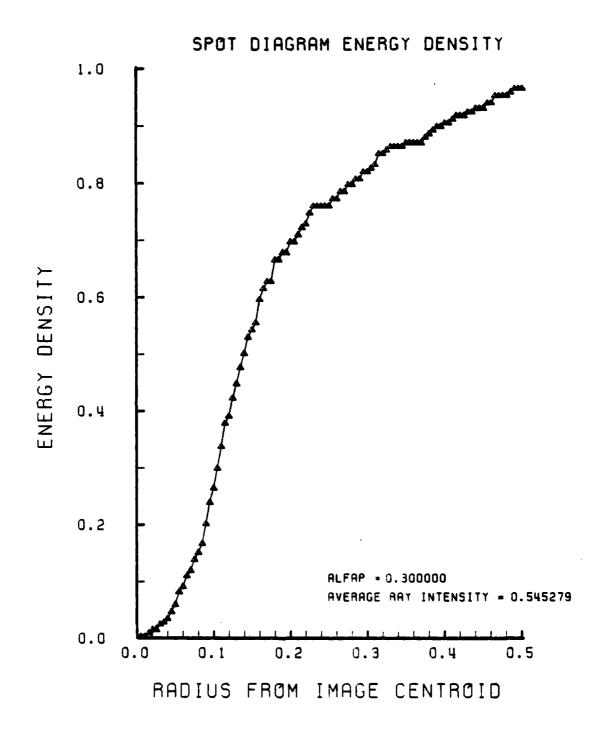


Figure D-27. Encircled Energy Plot for Spot Diagram of Figure D-26

APPENDIX E

## GRIN LENS PERFORMANCE PLOTS IN THE LOW RANGE OF INDICES OF REFRACTION (a = 2.25)

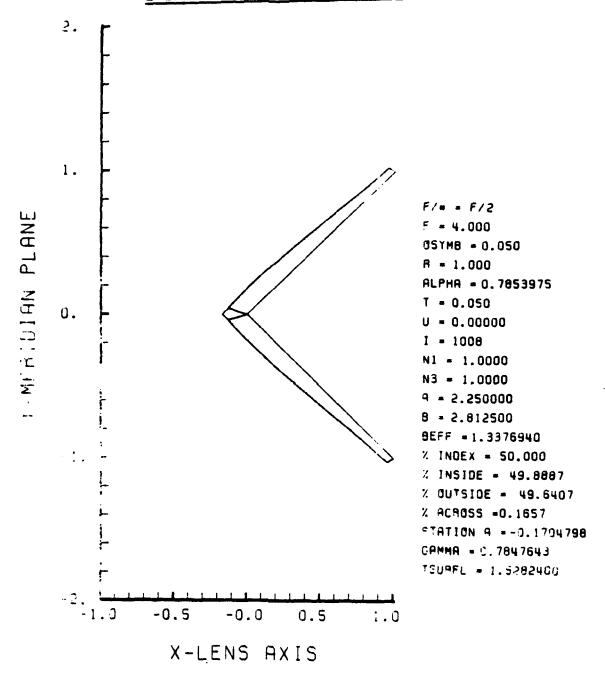
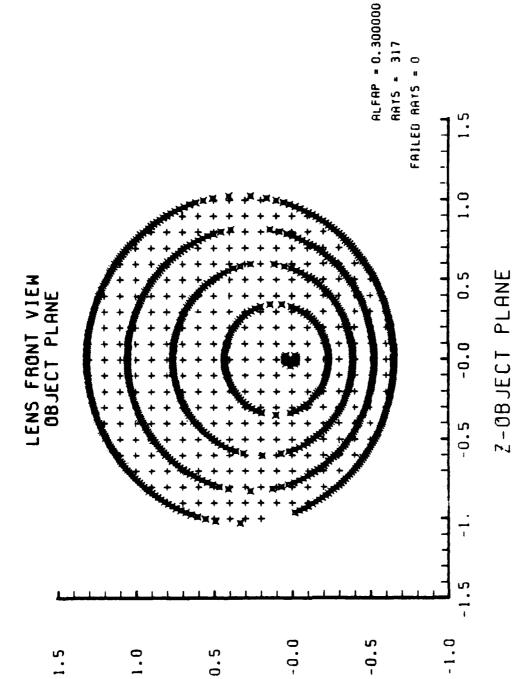


Figure E-1. GRIN Lens Shape at +50%, OB = 0.05, a = 2.25



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Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-1 Figure B-2.

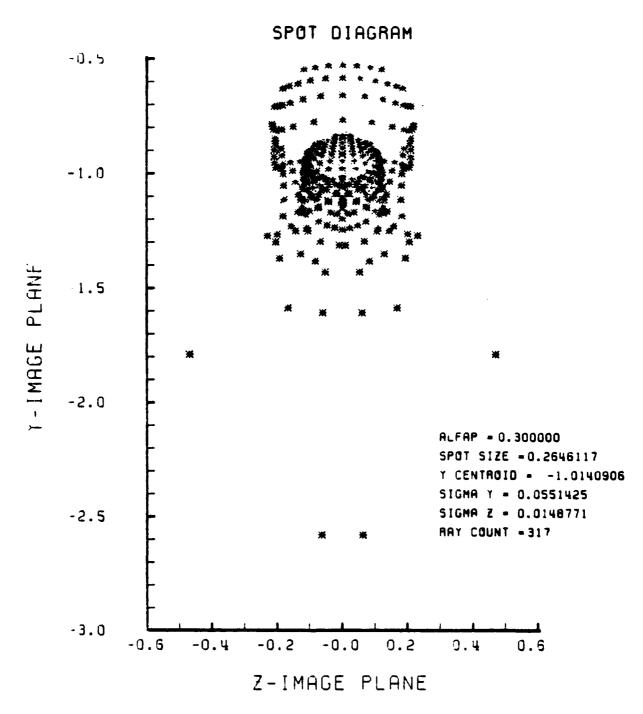


Figure E-3. Spot Diagram for Grid of Figure E-2

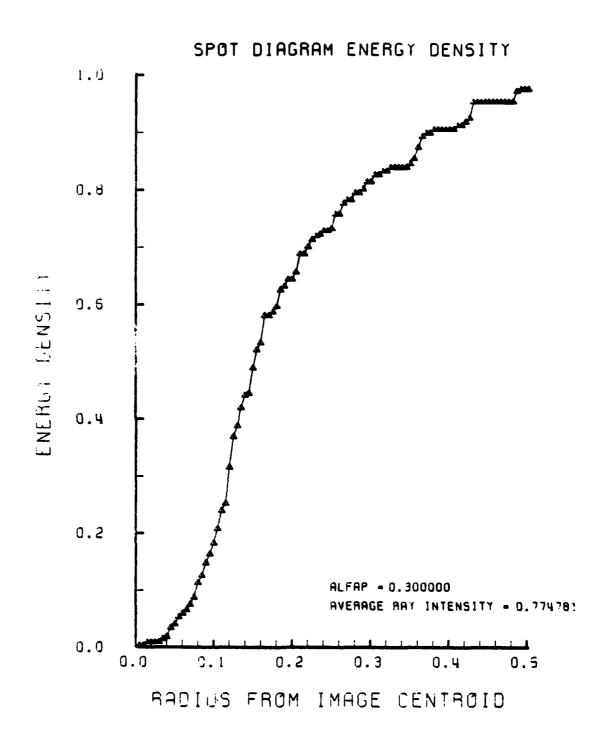


Figure E-4. Encircled Energy of Figure E-3

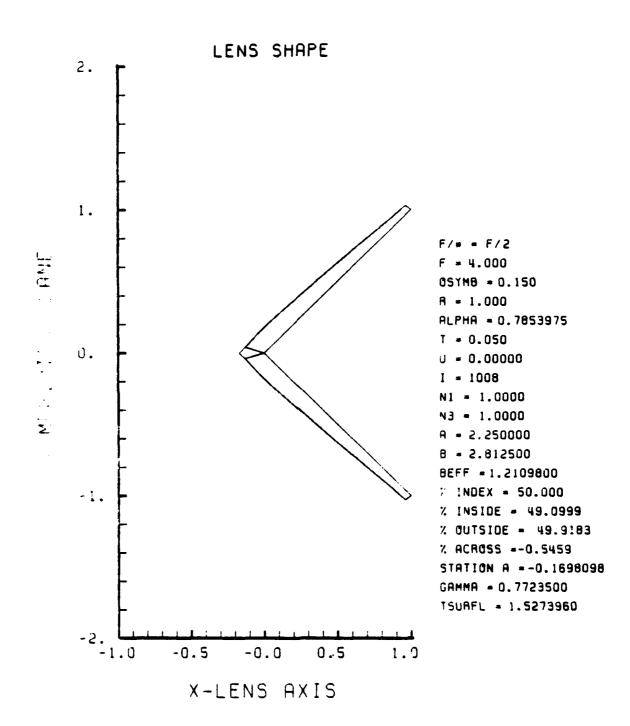
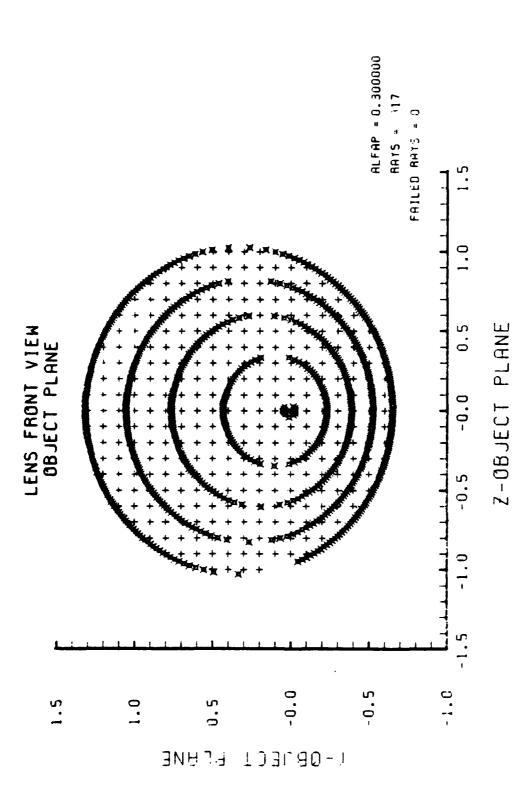


Figure E-5. GRIN Lens Shape at +50%, OB = 0.15, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-5 Figure E-6.

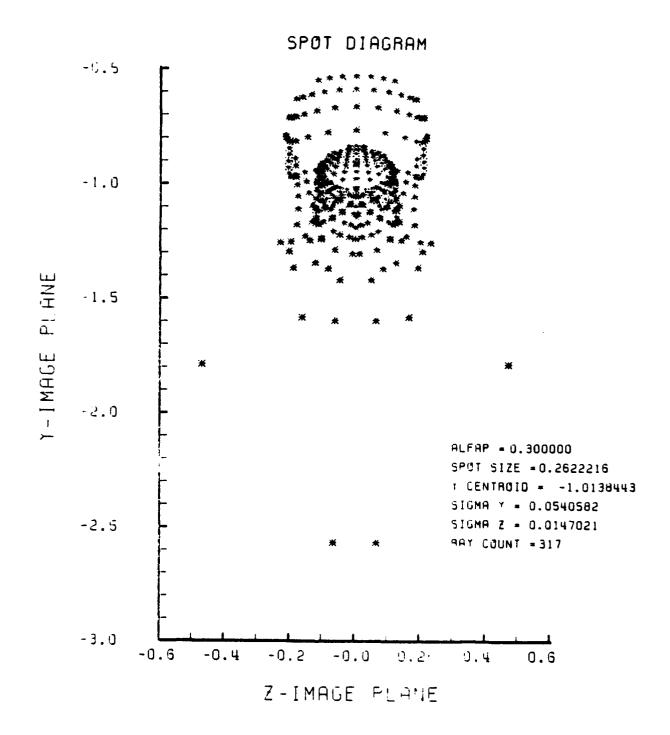


Figure E-7. Spot Diagram for Grid of Figure E-6

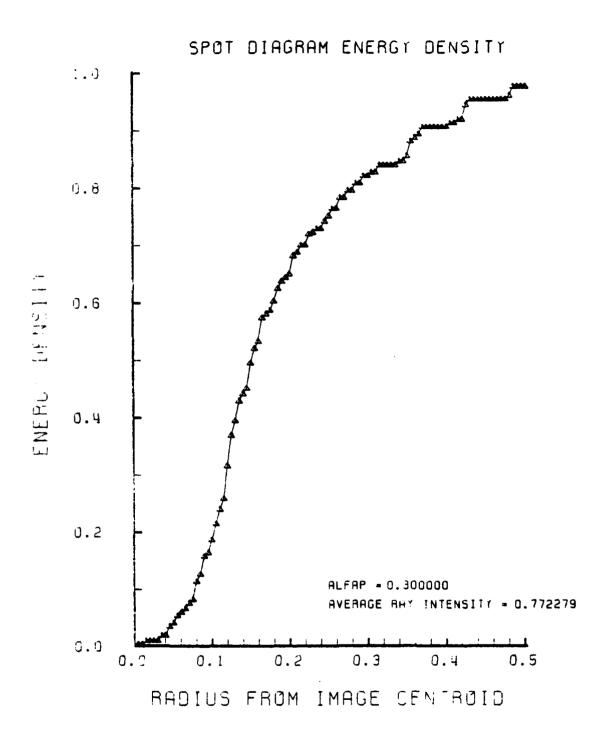


Figure E-8. Encircled Energy of Figure E-7

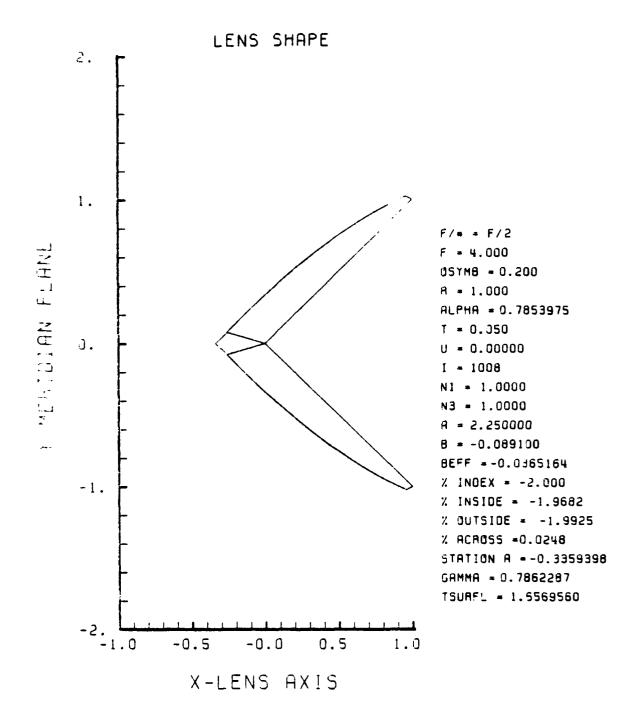
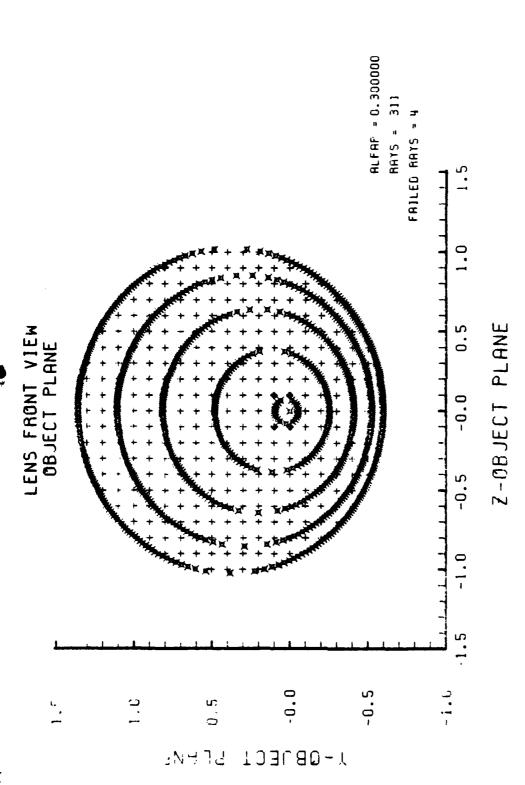


Figure E-9. GRIN Lens Shape at -2%, OB = 0.20, a = 2.25



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Grid Plane at  $\alpha$  = 0.3 for Lens of Figure E-9 Figure E-10.

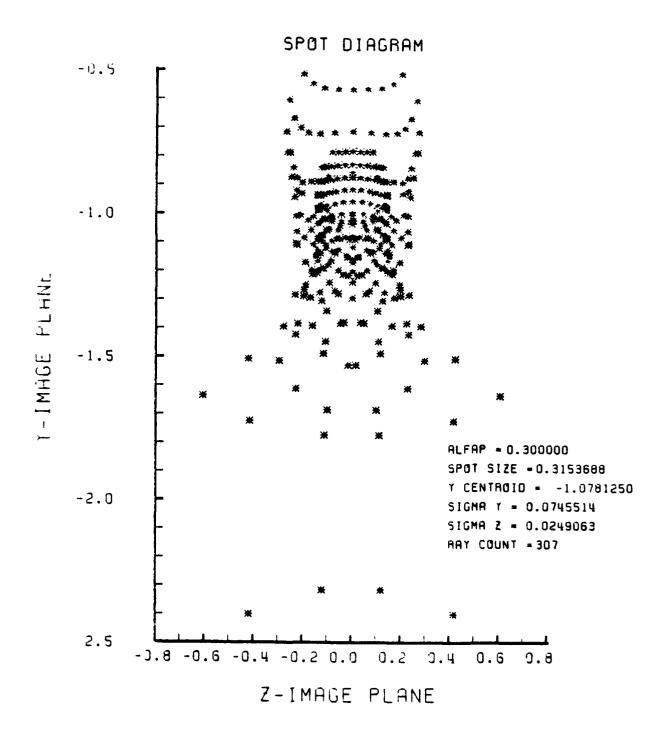


Figure E-11. Spot Diagram for Grid of Figure E-10

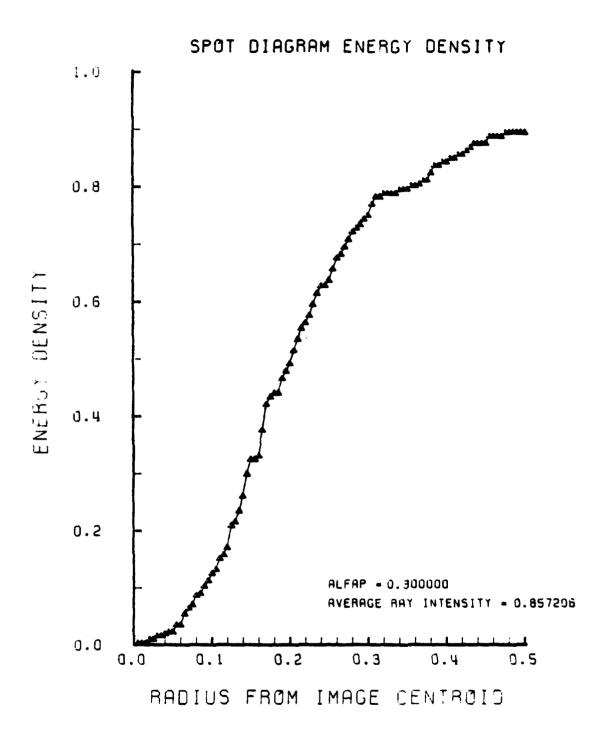


Figure E-12. Encircled Energy of Figure E-11

The second of th

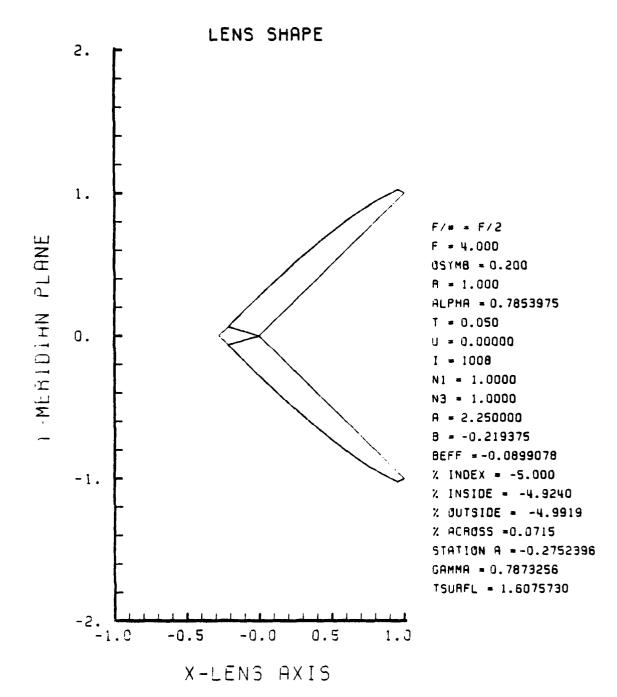
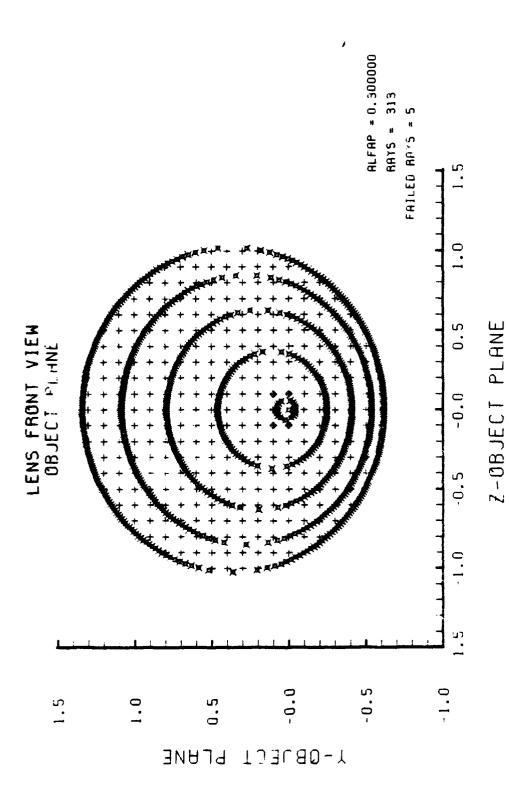


Figure E-13. GRIN Lens Shape at -5%, OB = 0.20, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-13 Figure E-14.

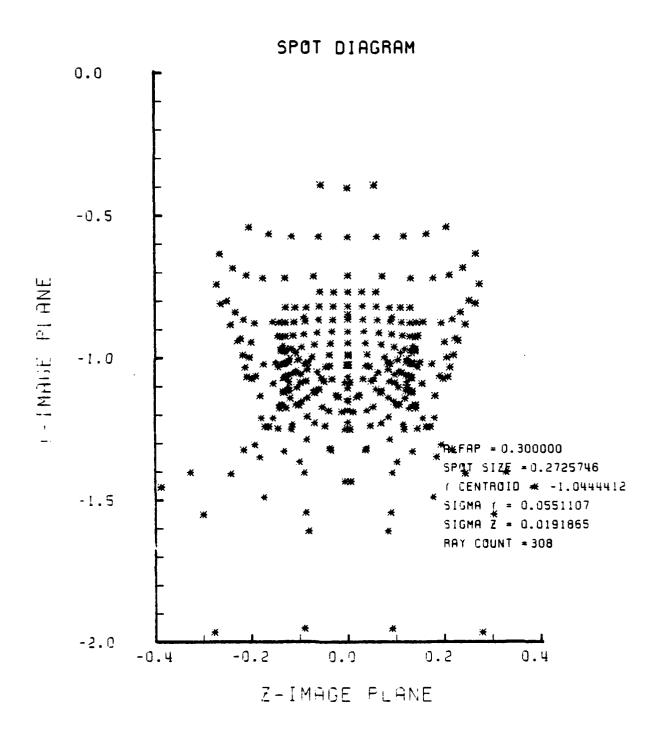


Figure E-15. Spot Diagram for Grid of Figure E-14

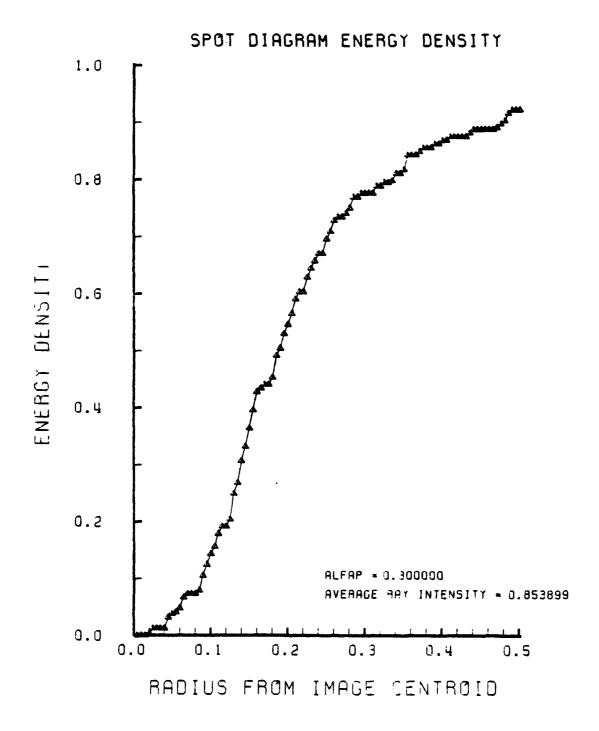


Figure E-16. Encircled Energy of Figure E-15

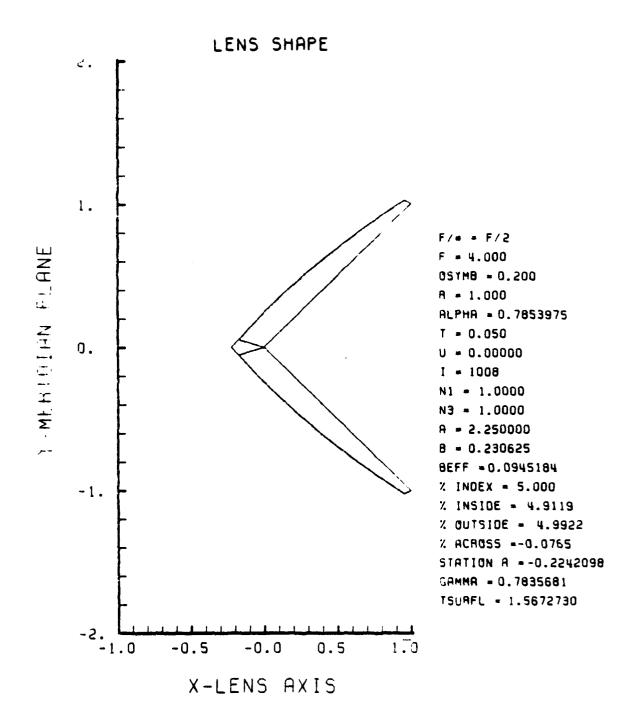
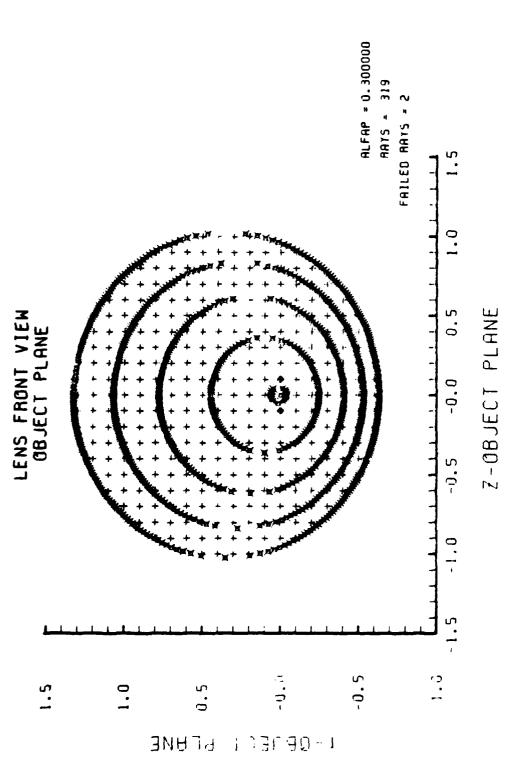


Figure E-17. GRIN Lens Shape at +5%, OB = 0.20, a = 2.25



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Grid Plane at  $\alpha_{\rm p}=0.3$  for Lens of Figure E-17 Figure E-18.

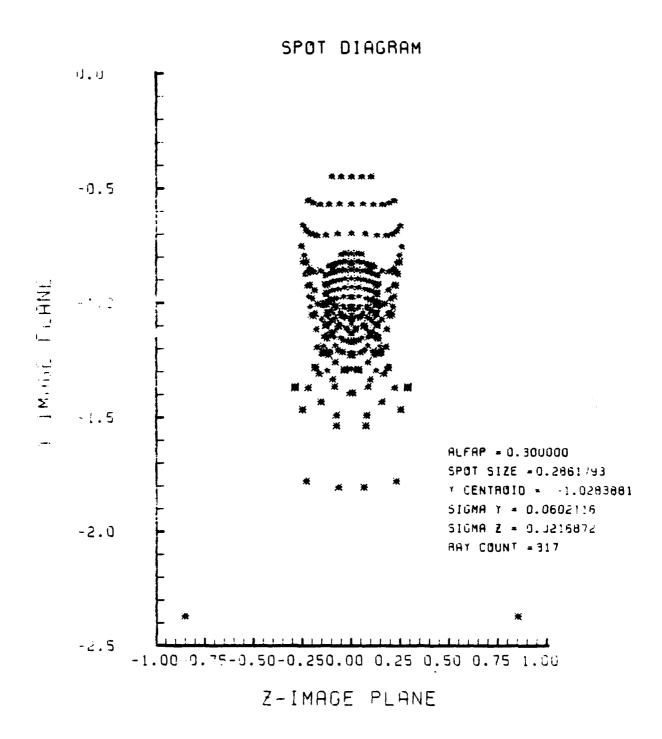


Figure E-19. Spot Diagram for Grid of Figure E-18

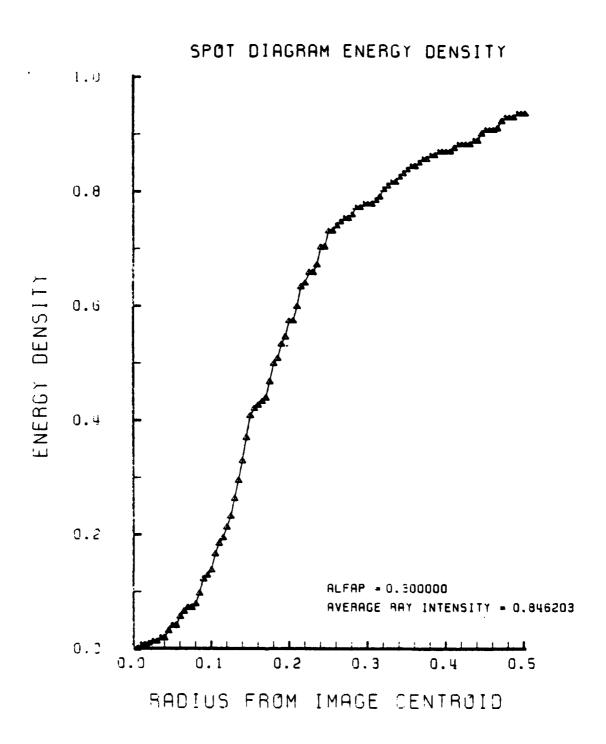


Figure E-20. Encircled Energy of Figure E-19

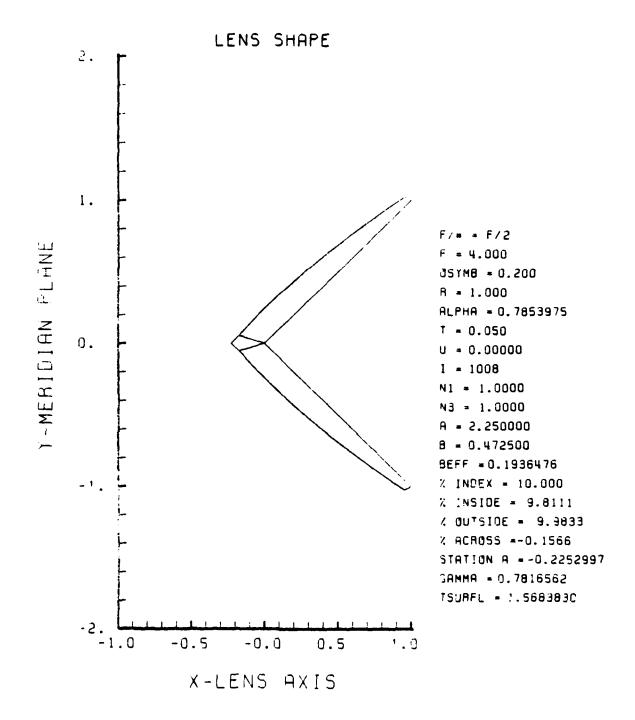
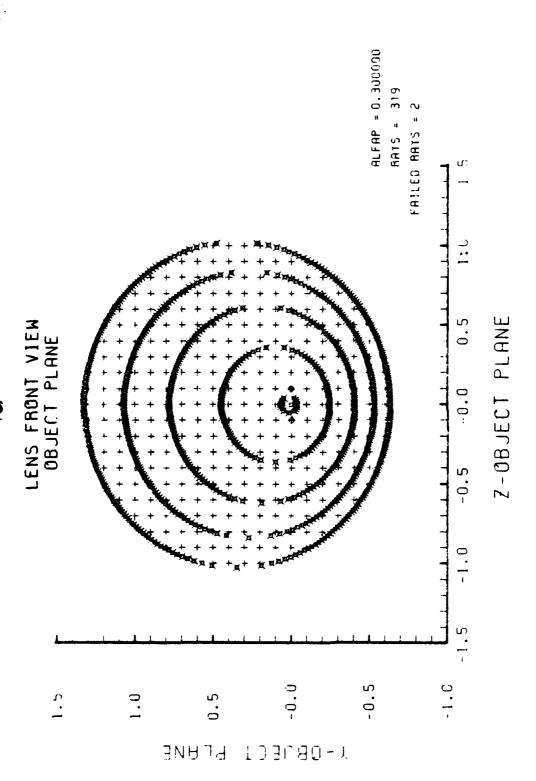


Figure E-21. GRIN Lens Shape at +10%, OB = 0.20, a = 2.25



Grid Plane at  $\alpha$  = 0.3 for Lens of Figure E-21 Figure E-22.

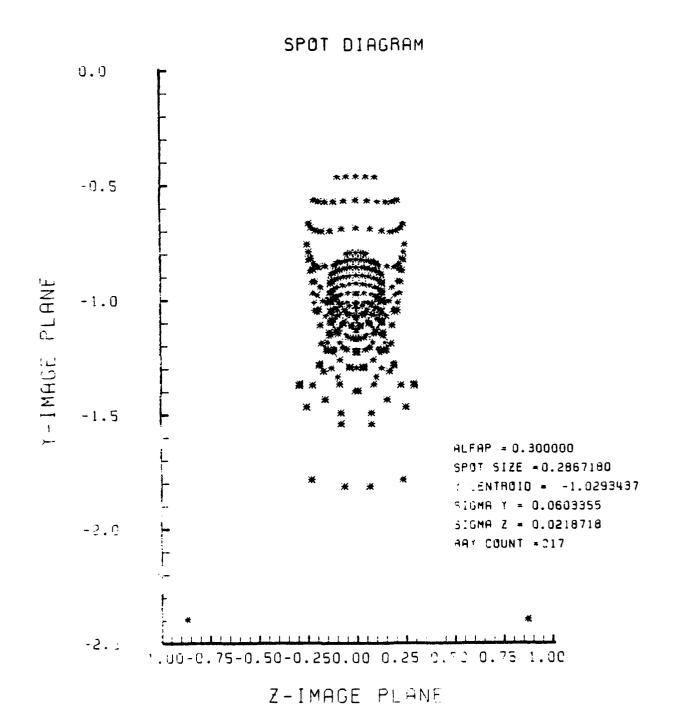


Figure E-23. Spot Diagram for Grid of Figure E-22

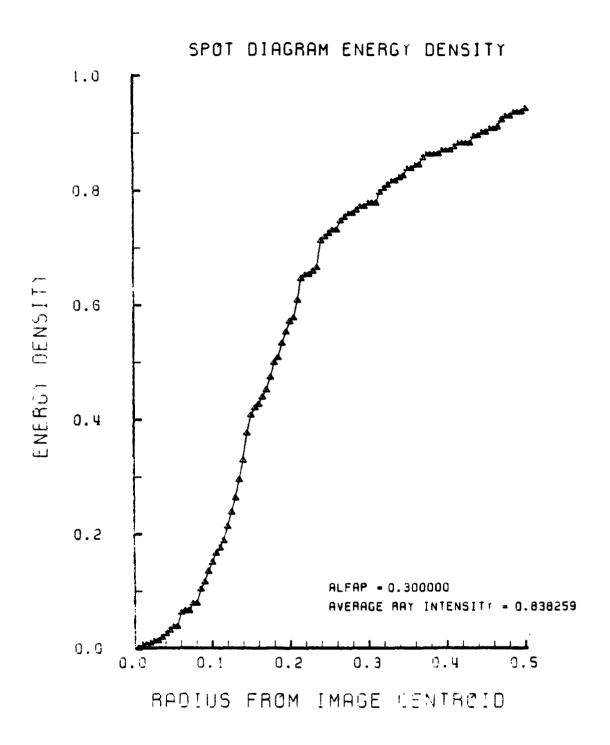


Figure E-24. Encircled Energy of Figure E-23

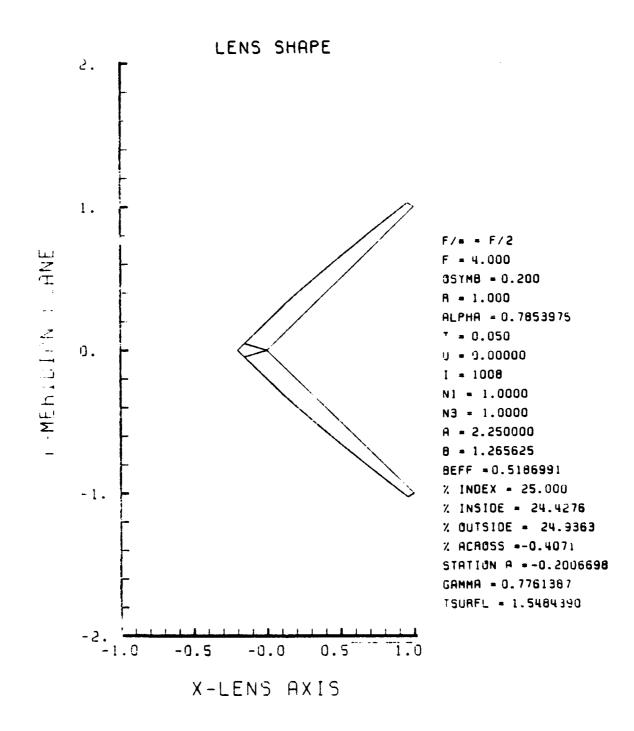


Figure E-25. GRIN Lens Shape at +25%, OB = 0.20, a = 2.25

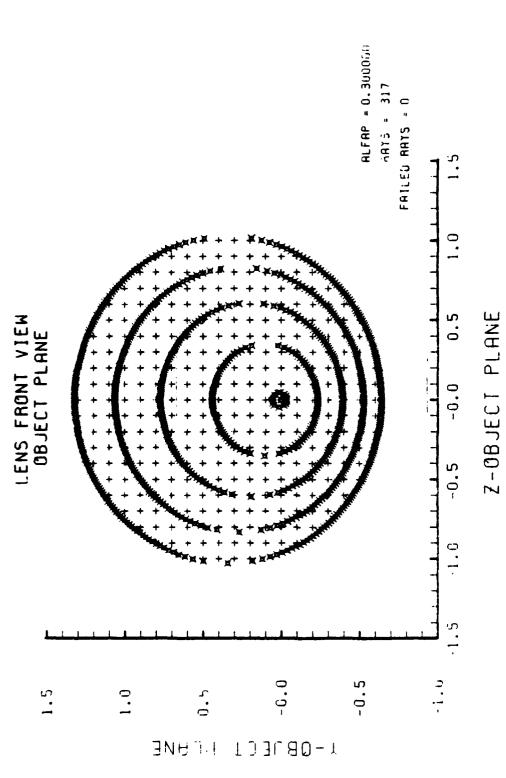


Figure E-26. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-25

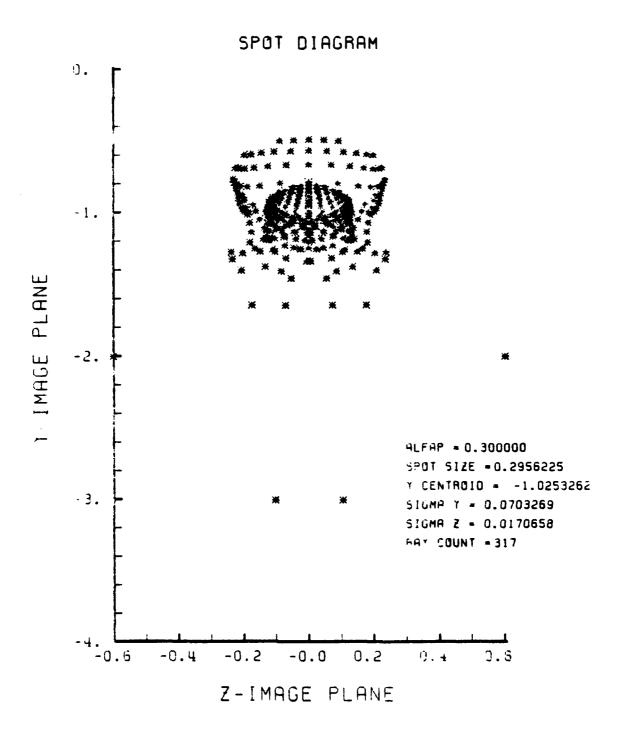


Figure E-27. Spot Diagram for Grid of Figure E-26

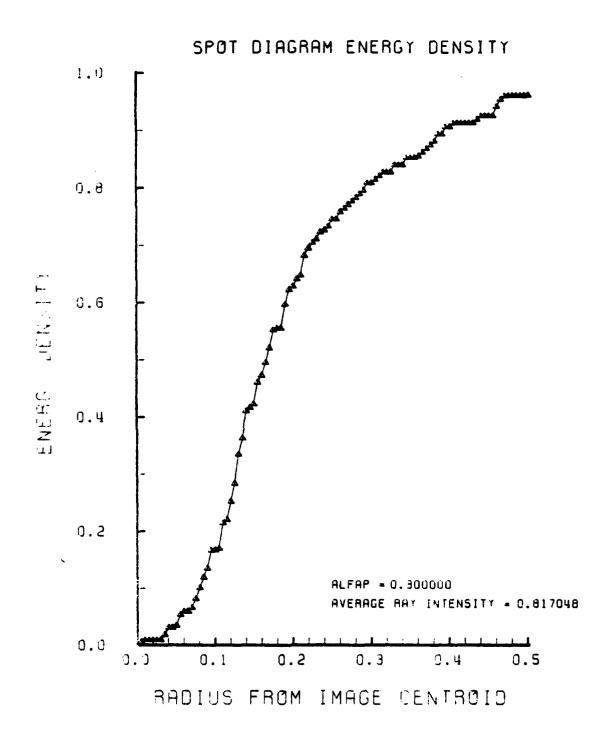


Figure E-28. Encircled Energy of Figure E-27

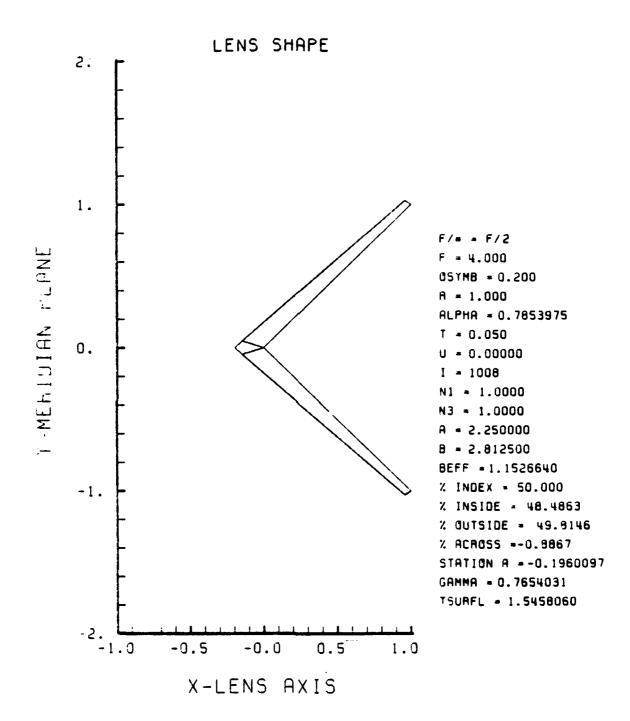
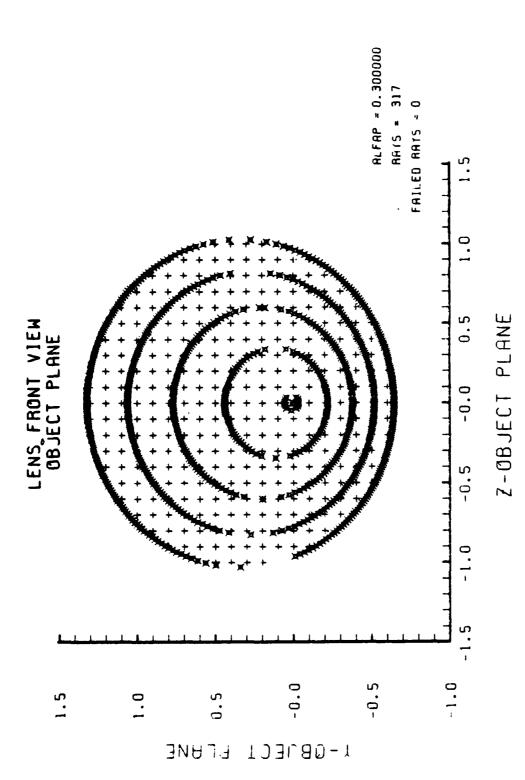


Figure E-29. GRIN Lens Shape at +50%, OB = 0.20, a = 2.25



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Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-29 Figure E-30.

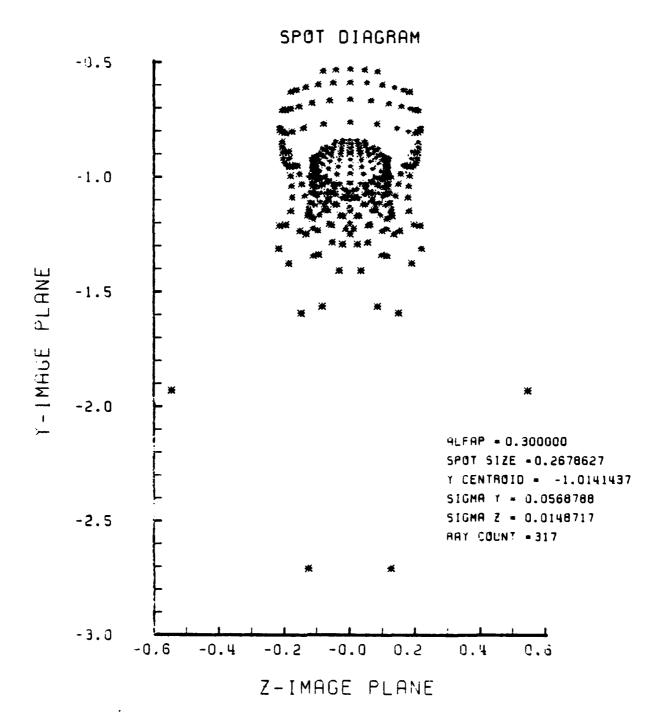


Figure E-31. Spot Diagram for Grid of Figure E-30

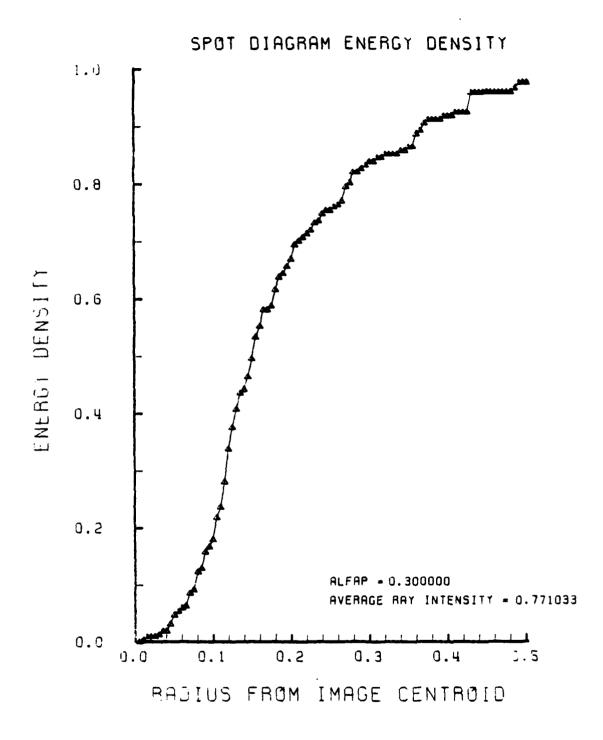


Figure E-32. Encircled Energy of Figure E-31

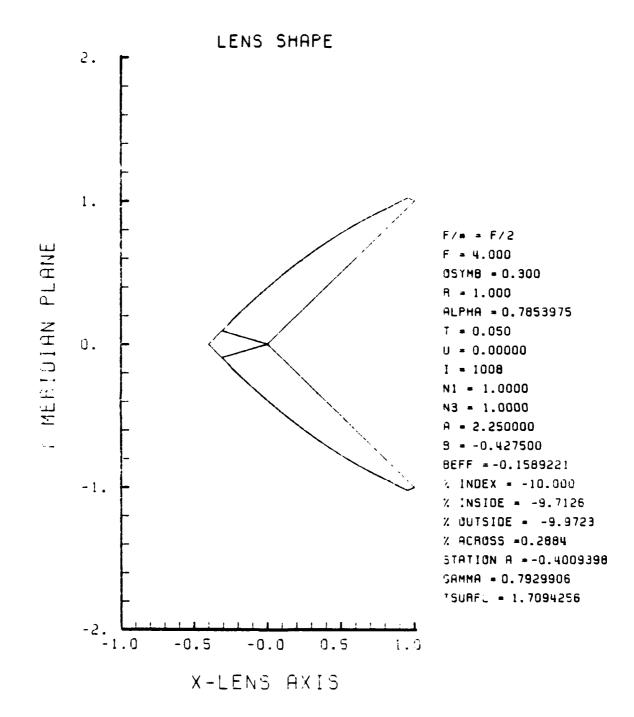
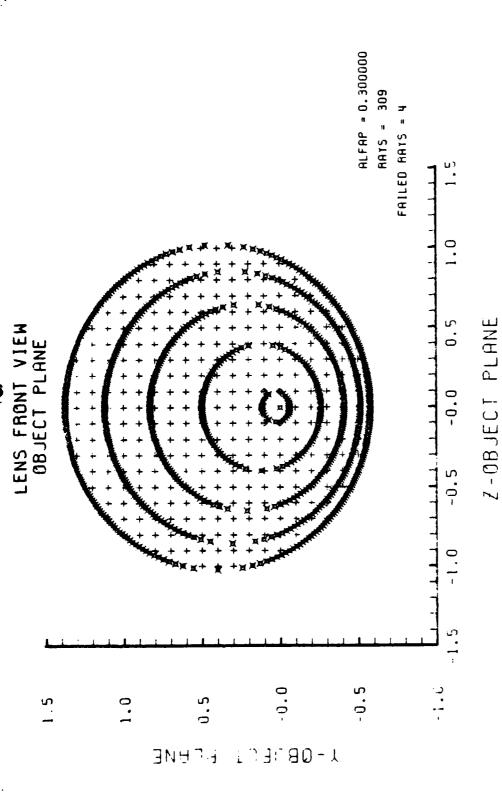


Figure E-33. GRIN Lens Shape at -10%, OB = 0.30, a = 2.25



Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure E-33 Figure E-34.

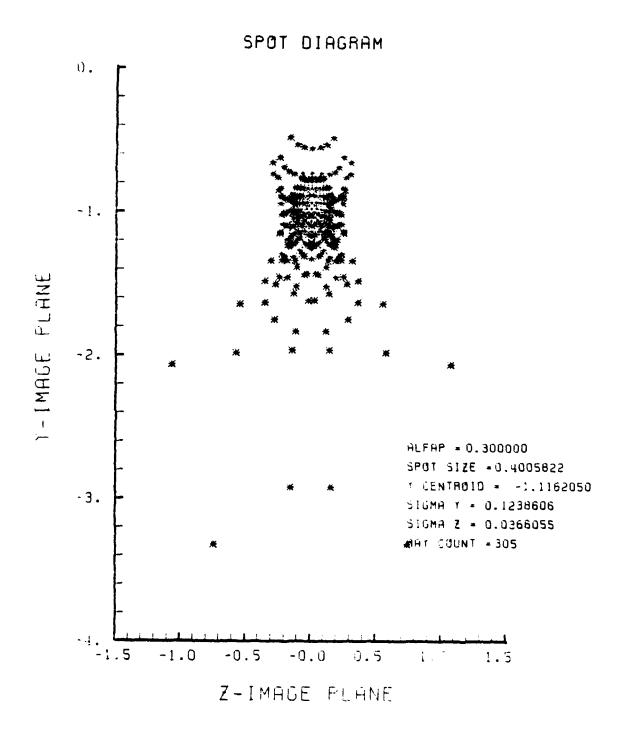


Figure E-35. Spot Diagram for Grid of Figure E-34

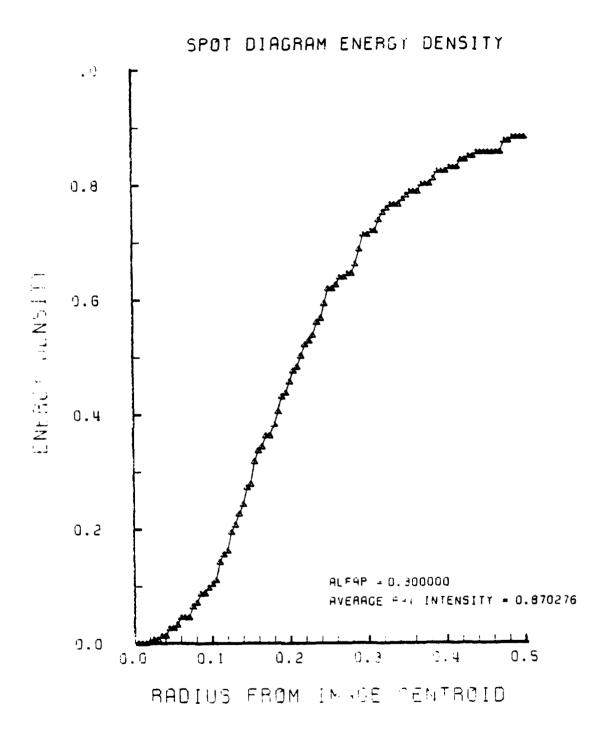


Figure E-36. Encircled Energy of Figure E-35

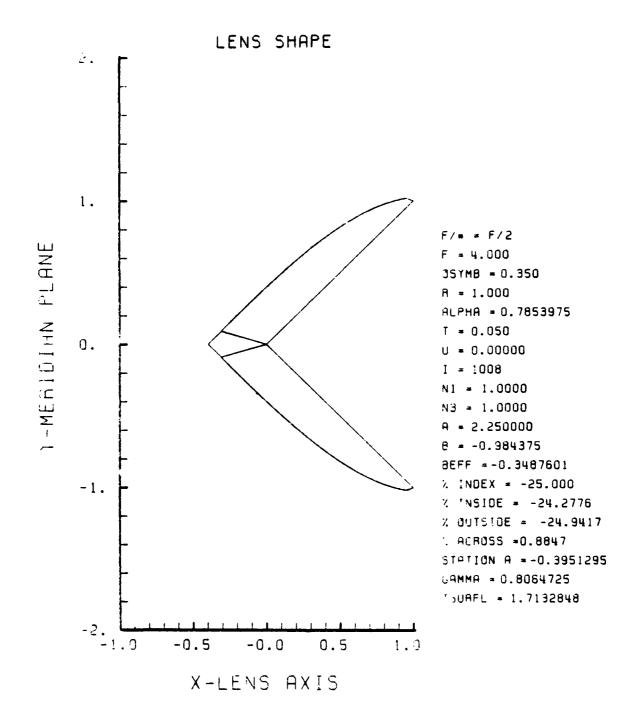
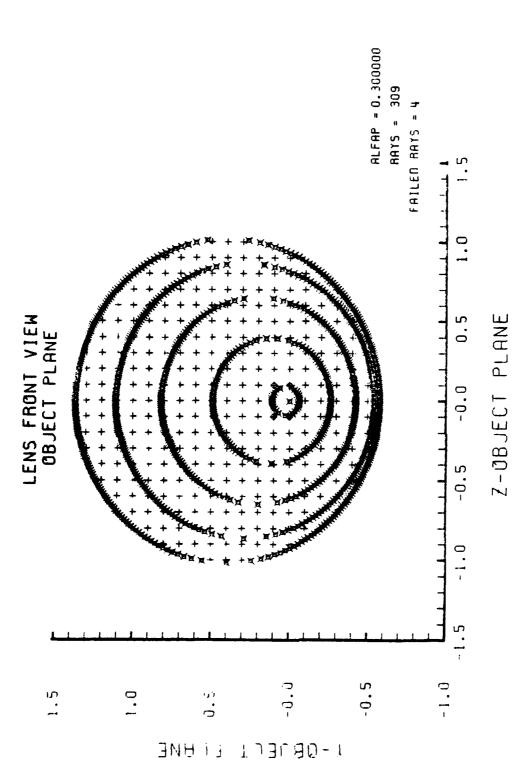


Figure E-37. GRIN Lens Shape at -25%, OB = 0.35, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-37 Figure E-38.

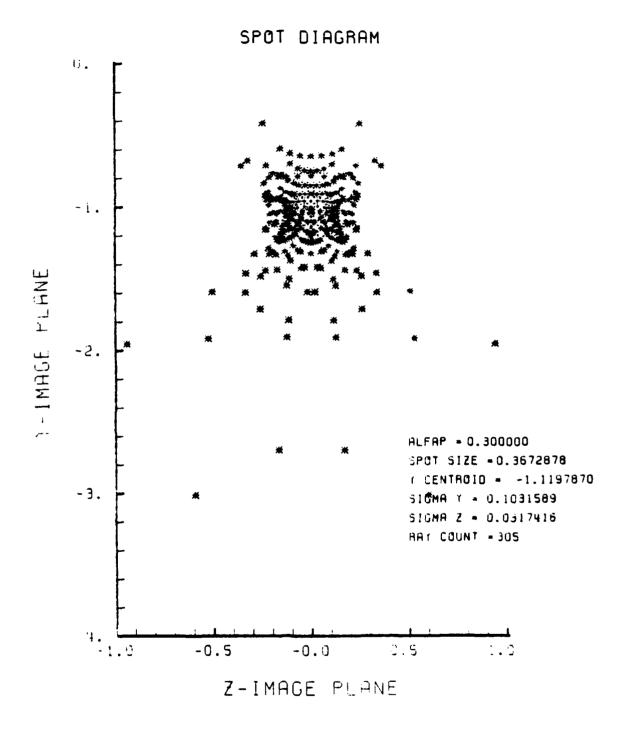


Figure E-39. Spot Diagram for Grid of Figure E-38

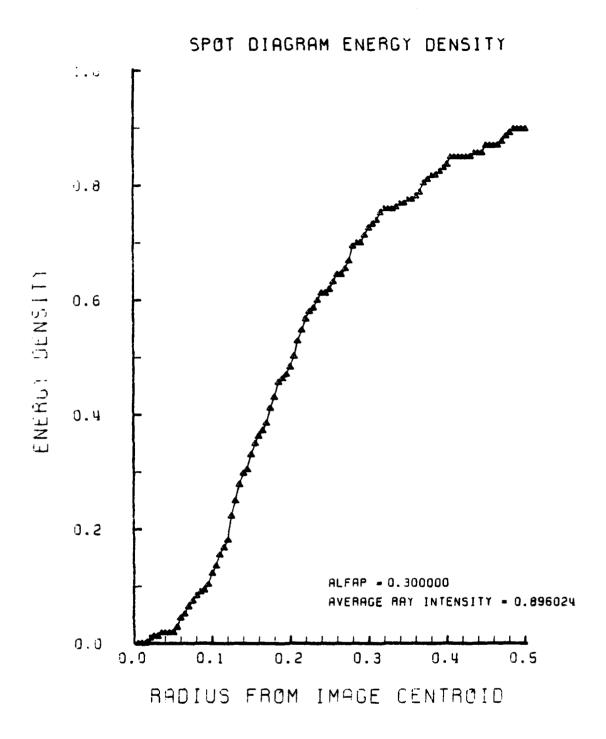


Figure E-40. Encircled Energy of Figure E-39

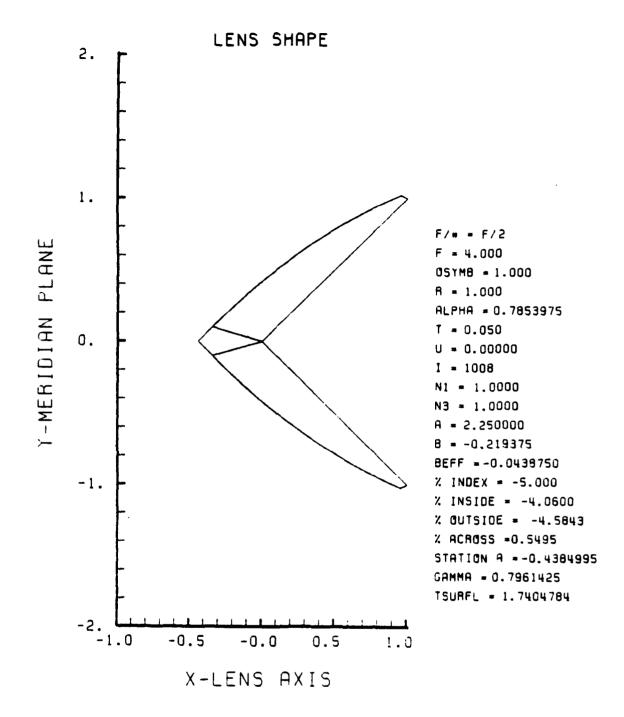


Figure E-41. GRIN Lens Shape at -5%, OB = 1.00, a = 2.25

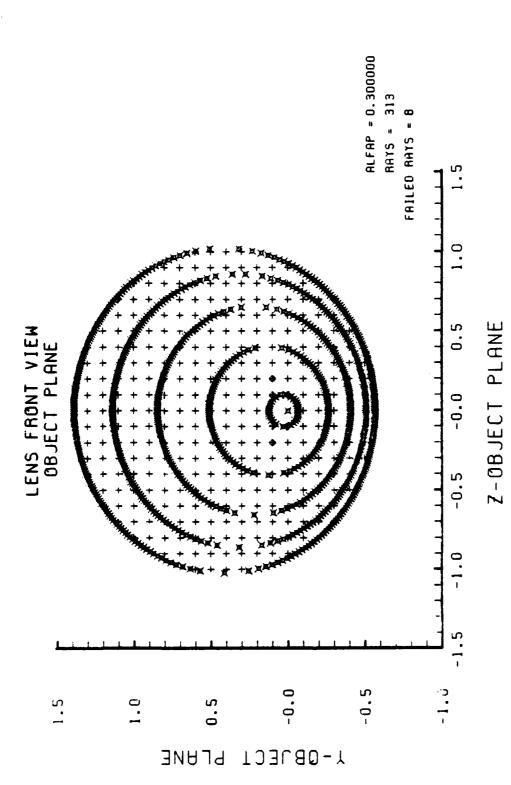


Figure E-42. Grid Plane at  $\alpha_p = 0.3$  for Lens of Figure E-41

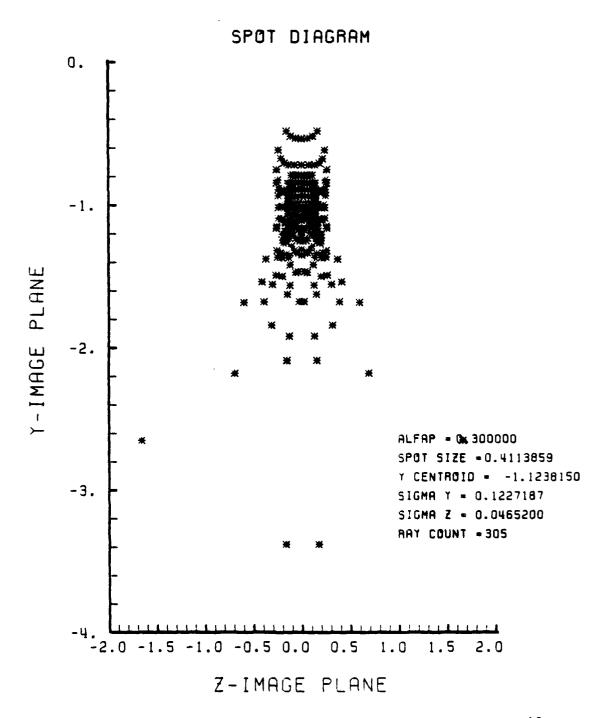


Figure E-43. Spot Diagram for Grid of Figure E-42

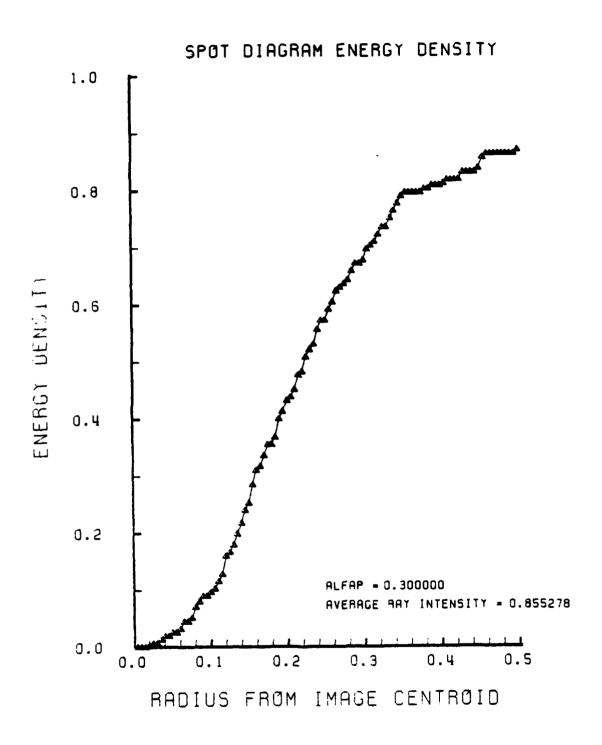


Figure E-44. Encircled Energy of Figure E-43

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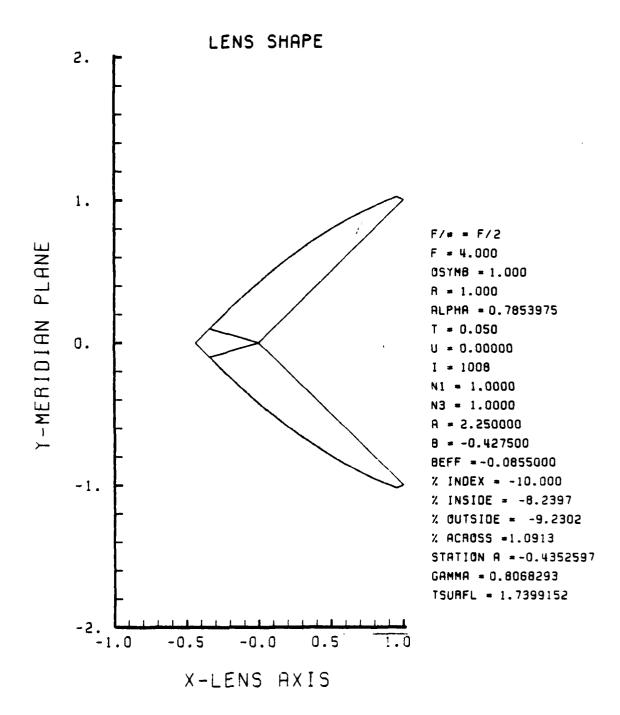
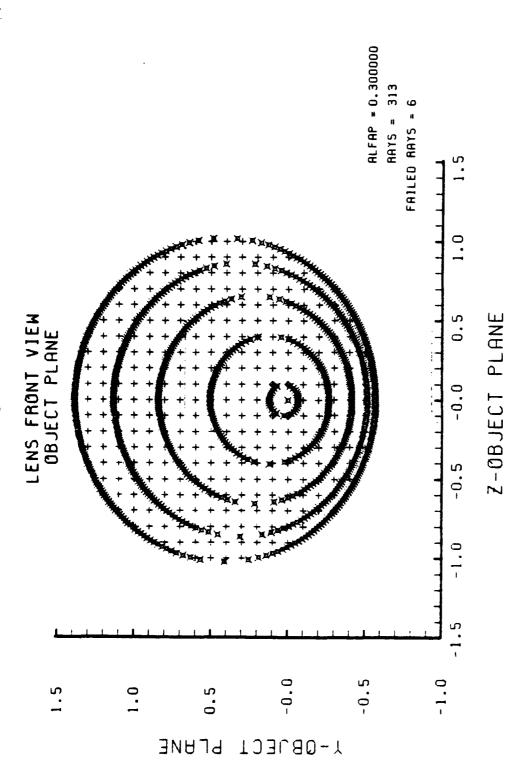


Figure E-45. GRIN Lens Shape at -10%, OB = 1.00 a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-45 Figure E-46.

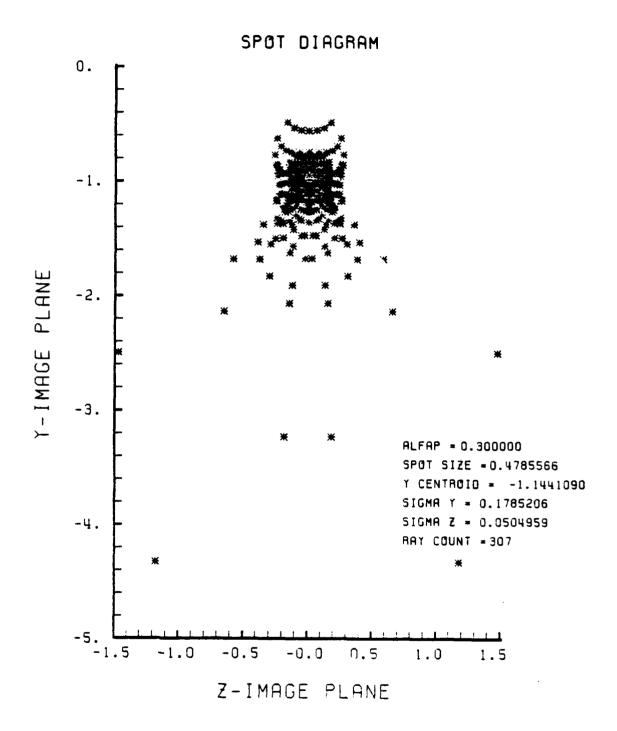


Figure E-47. Spot Diagram for Grid of Figure E-46

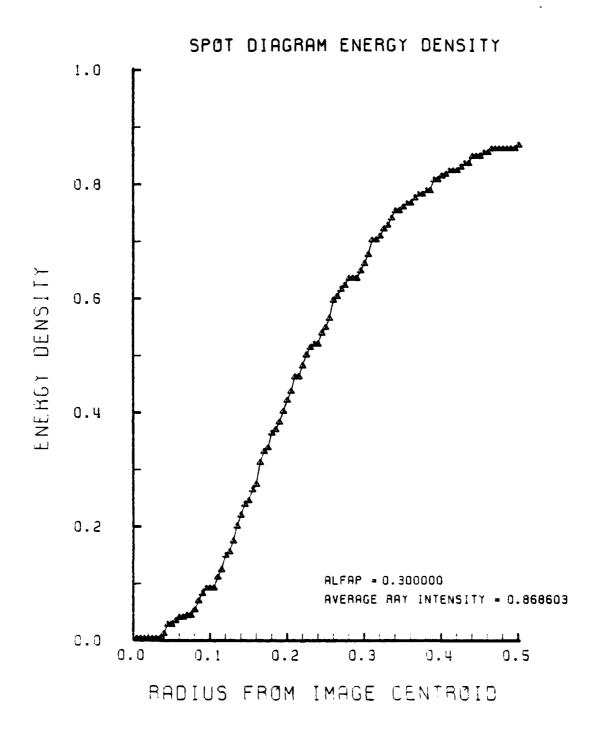


Figure E-48. Encircled Energy of Figure E-47

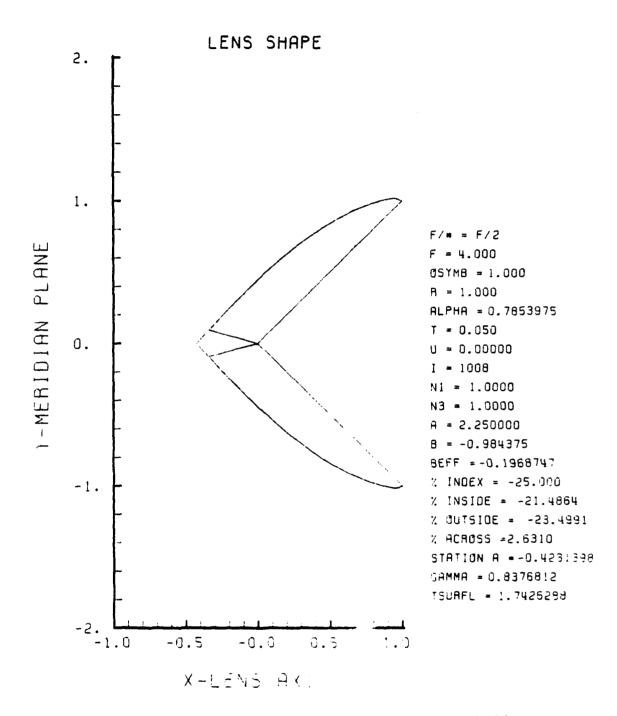
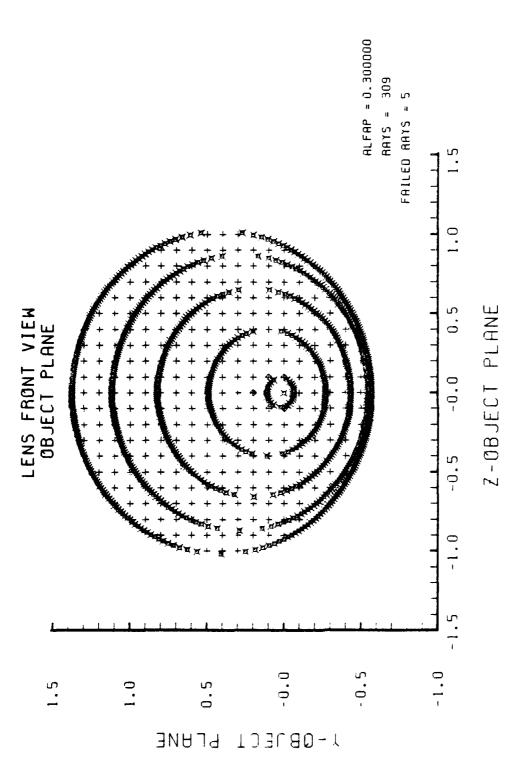


Figure E-49. GRIN Lens Shape at -25%, OB = 1.00, a = 2.25



Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure E-49 Figure E-50.

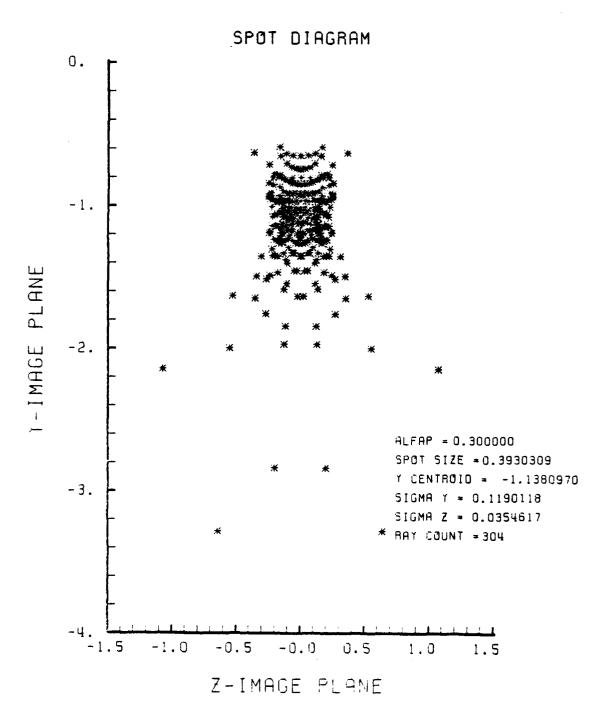


Figure E-51. Spot Diagram for Grid of Figure E-50

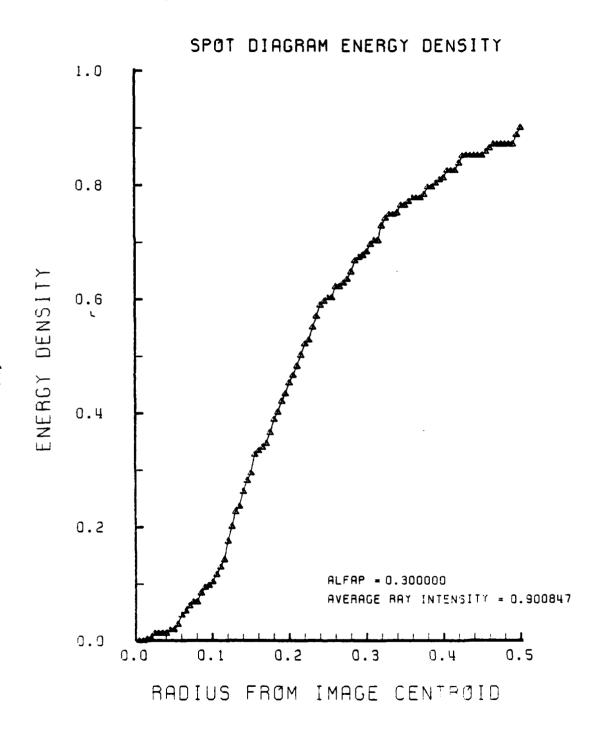


Figure E-52. Encircled Energy of Figure E-51

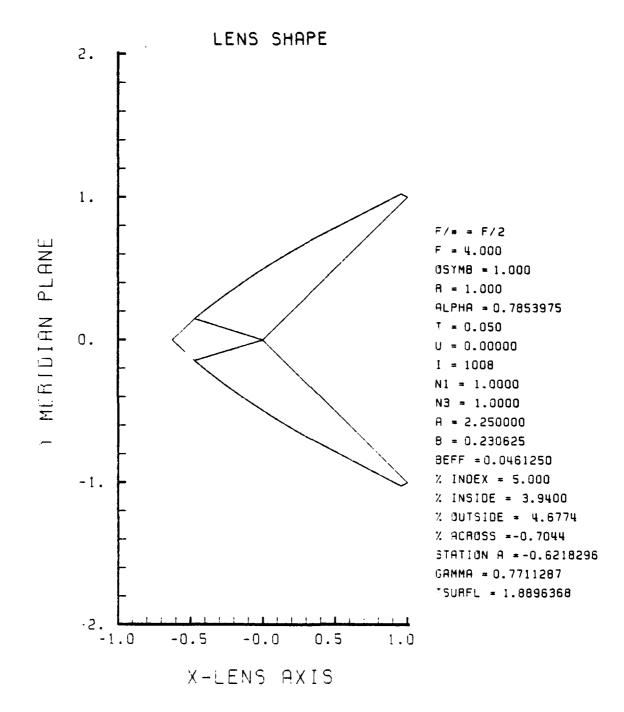
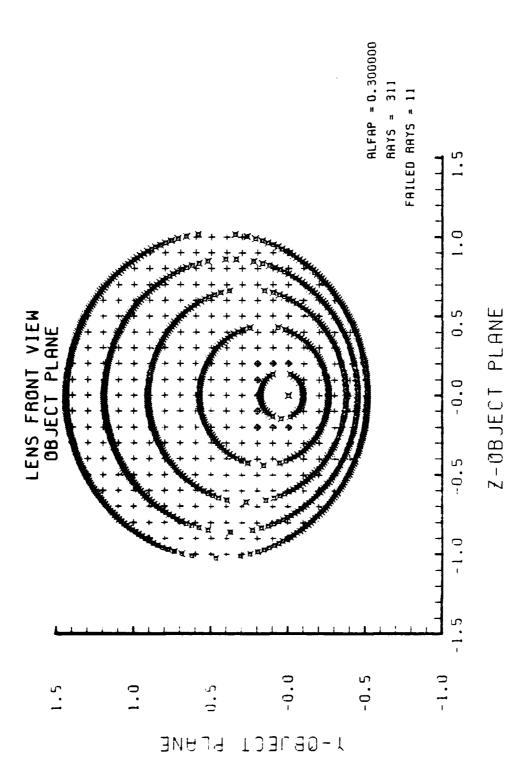


Figure E-53. GRIN Lens Shape at +5%, OB = 1.00, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-53 Figure E-54.

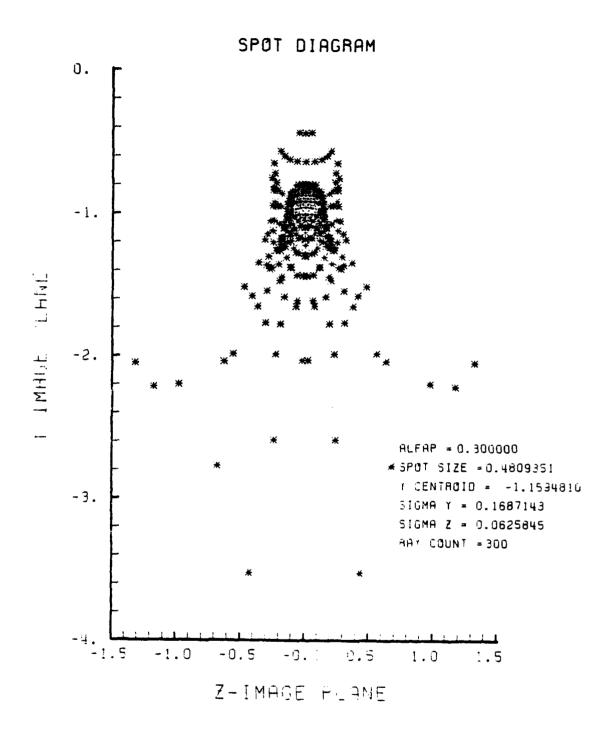


Figure E-55. Spot Diagram for Grid of Figure E-54

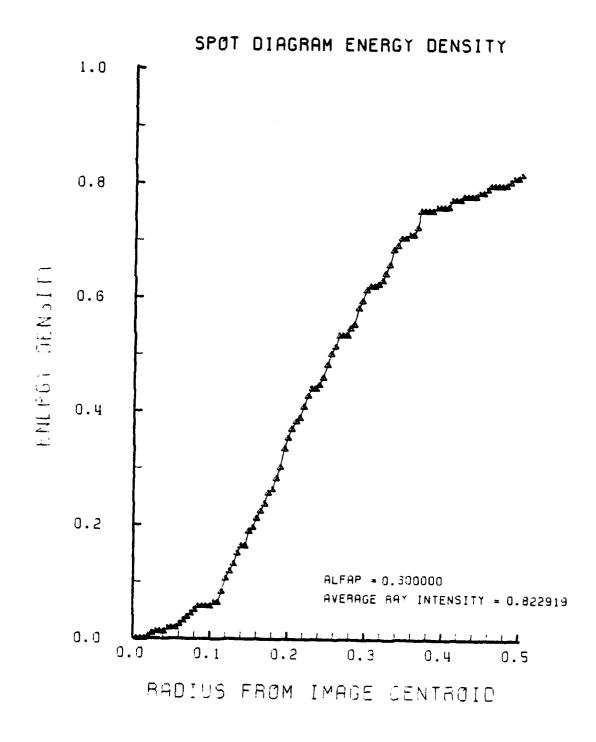


Figure E-56. Encircled Energy of Figure E-55

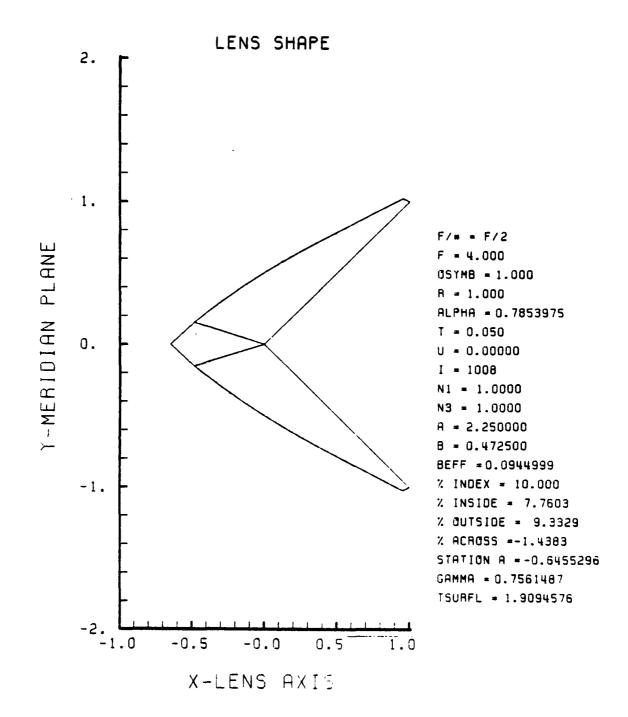
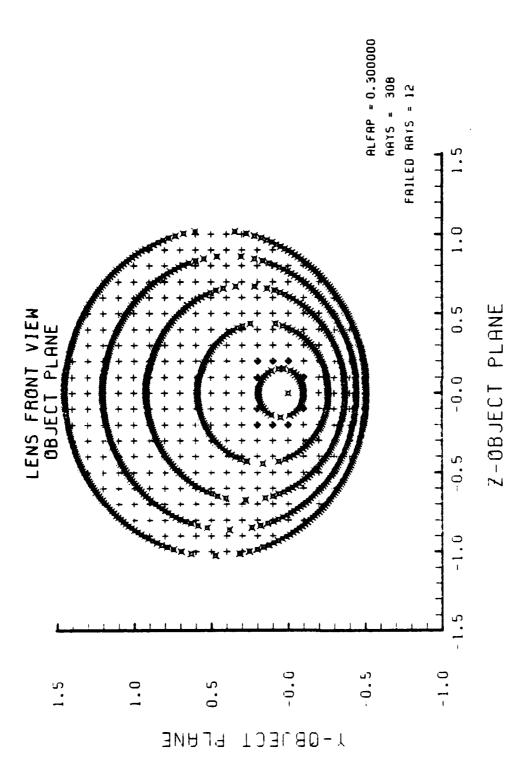


Figure E-57. GRIN Lens Shape at +10%, OB = 1.00, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-57 Figure E-58.

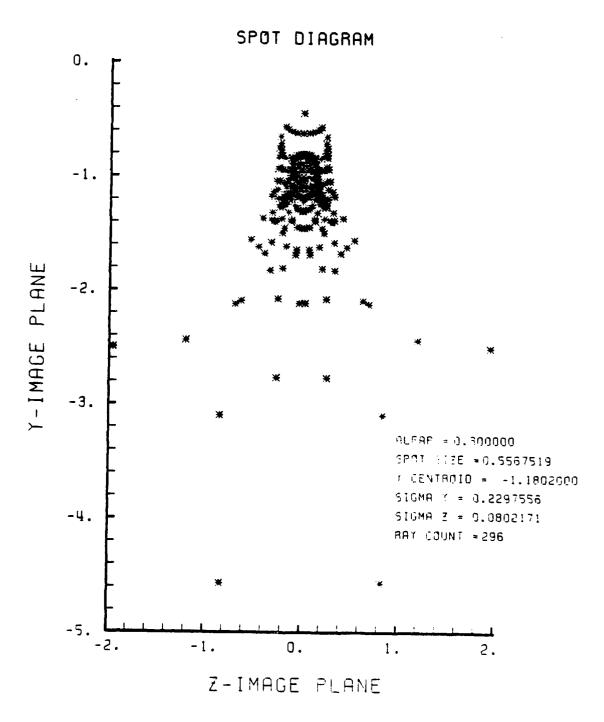


Figure E-59. Spot Diagram for Grid of Figure E-58

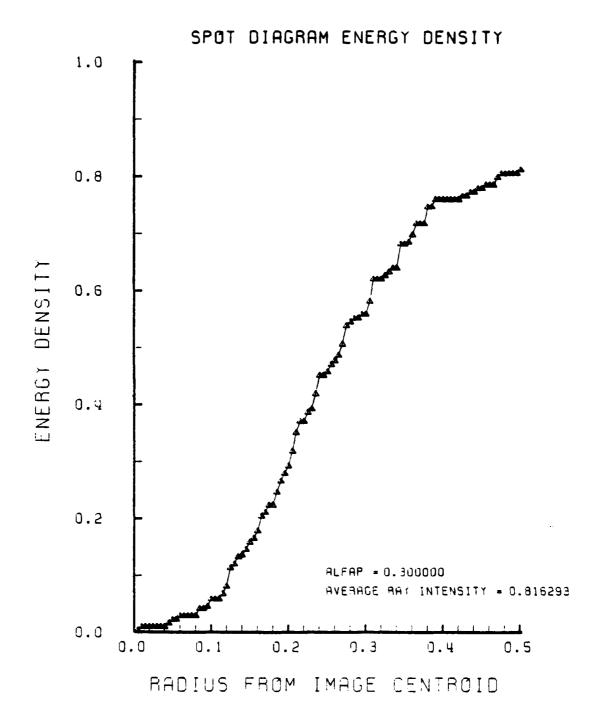


Figure E-60. Encircled Energy of Figure E-59

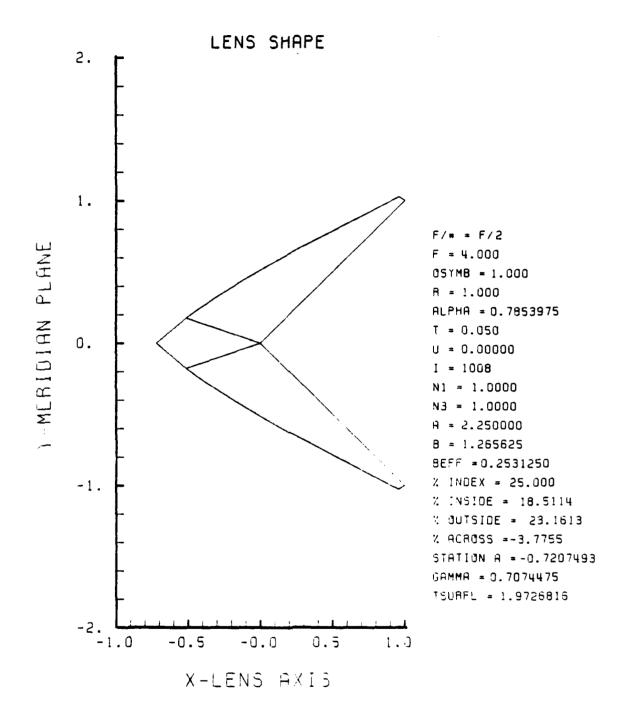
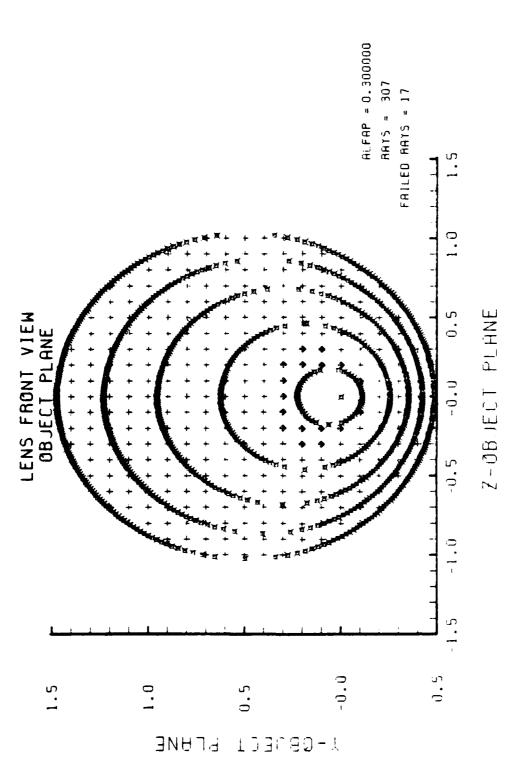


Figure E-61. GRIN Lens Shape at +25%, OB = 1.00, a = 2.25



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-61 Figure E-62.

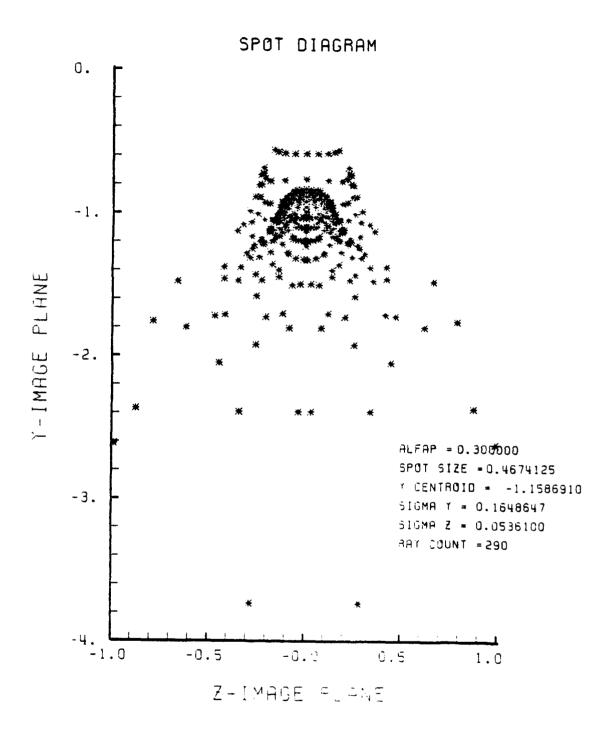


Figure E-63. Spot Diagram for Grid of Figure E-62

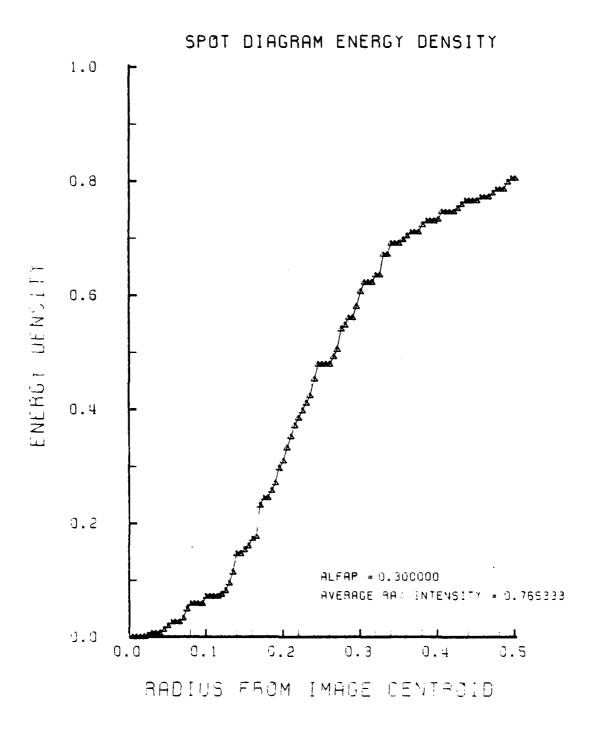


Figure E-64. Encircled Energy of Figure E-63

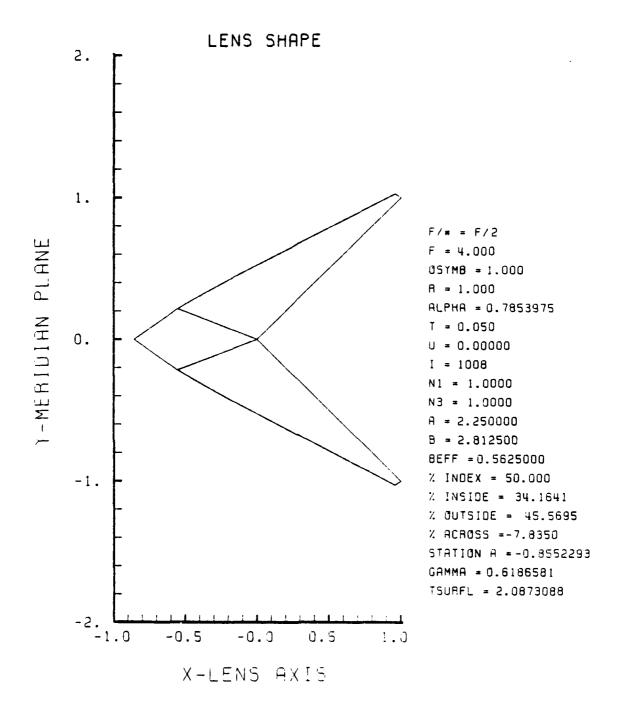


Figure E-65. GRIN Lens Shape at +50%, OB = 1.00, a = 2.25

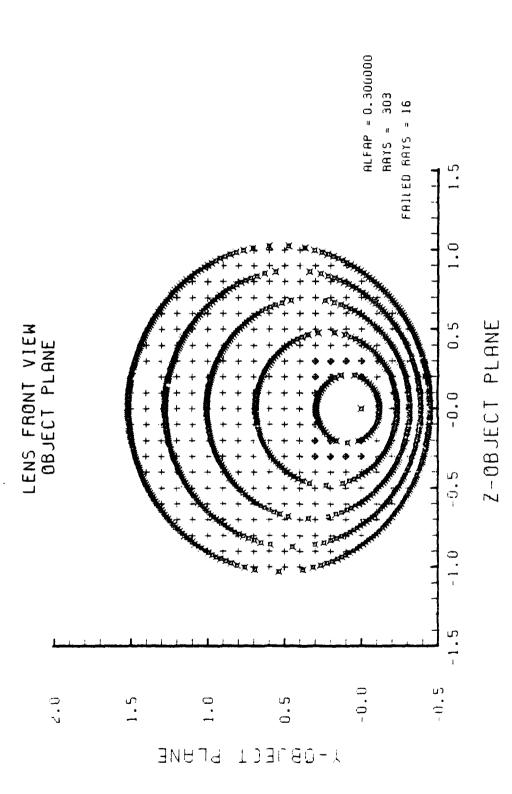


Figure E-66. Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure E-65

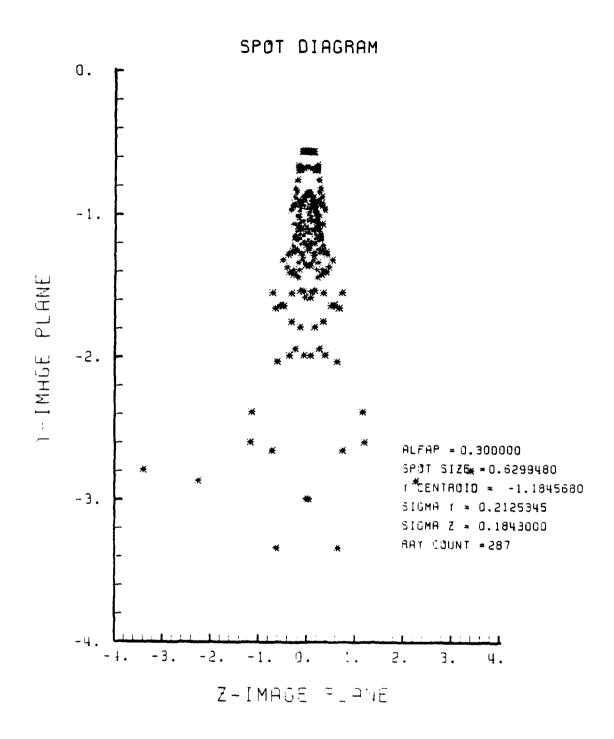


Figure E-67. Spot Diagram for Grid of Figure E-66

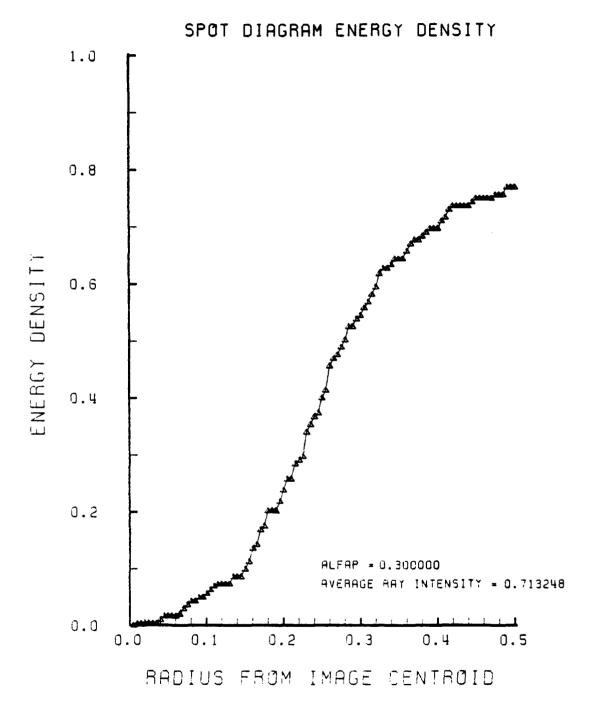
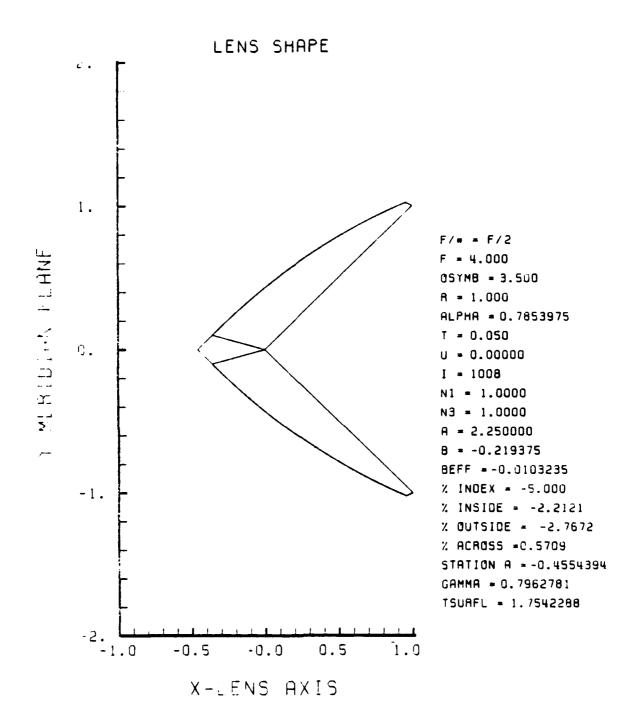


Figure E-68. Encircled Energy of Figure E-67



생활하게 되어 하면 보일에 가득하다가 나라가 하면 하는 사람들이 하는 것이 되었다. 그는 사람들이 되었다는 것이 없는 것이다.

Figure E-69. GRIN Lens Shape at -5%, OB = 3.50, a = 2.25

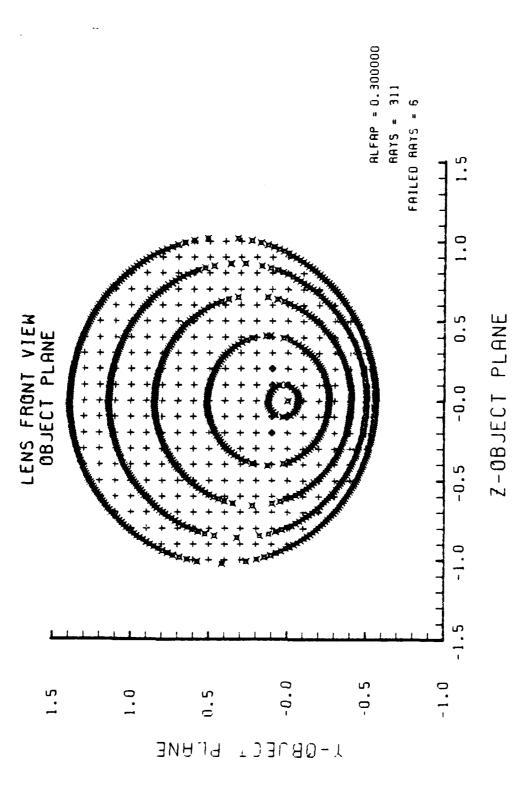


Figure E-70. Grid Plane at  $\alpha_p$   $\approx$  0.3 for Lens of Figure E-69

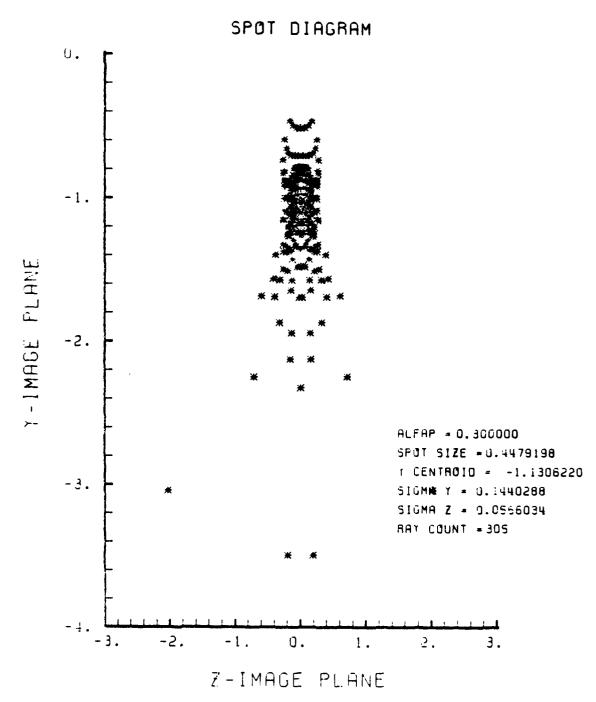


Figure E-71. Spot Diagram for Grid of Figure E-70

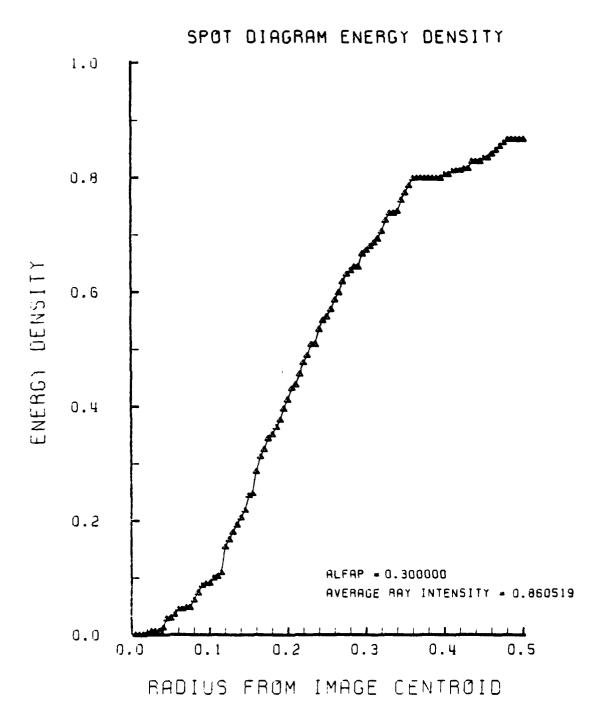


Figure E-72. Encircled Energy of Figure E-71

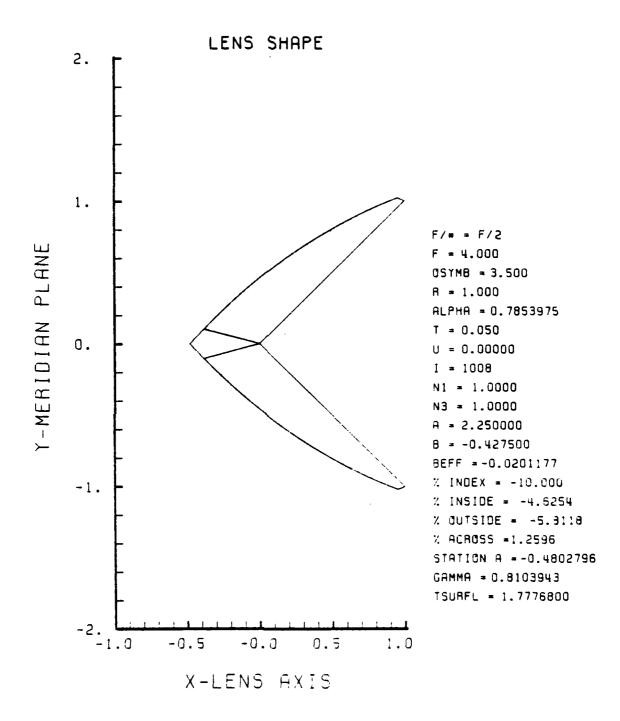
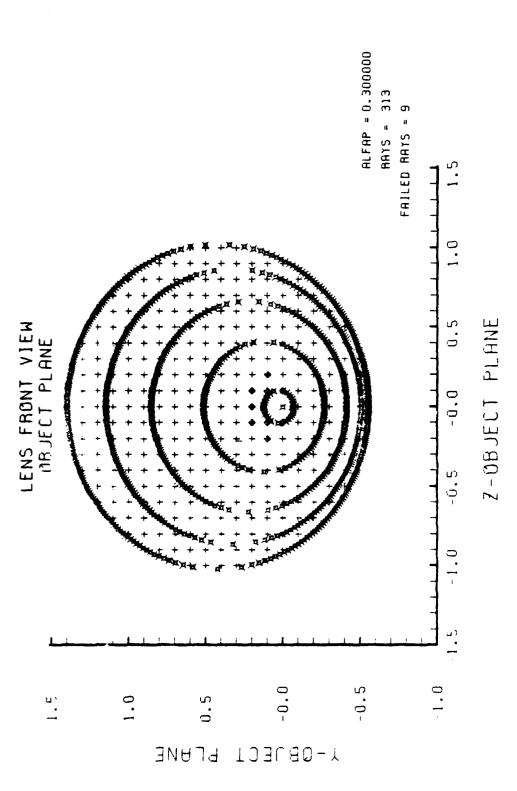


Figure E-73. GRIN Lens Shape at -10%, OB = 3.50, a = 2.25



Grid Plane at  $\alpha_{\rm p}=0.3$  for Lens of Figure E-73 Figure E-74.

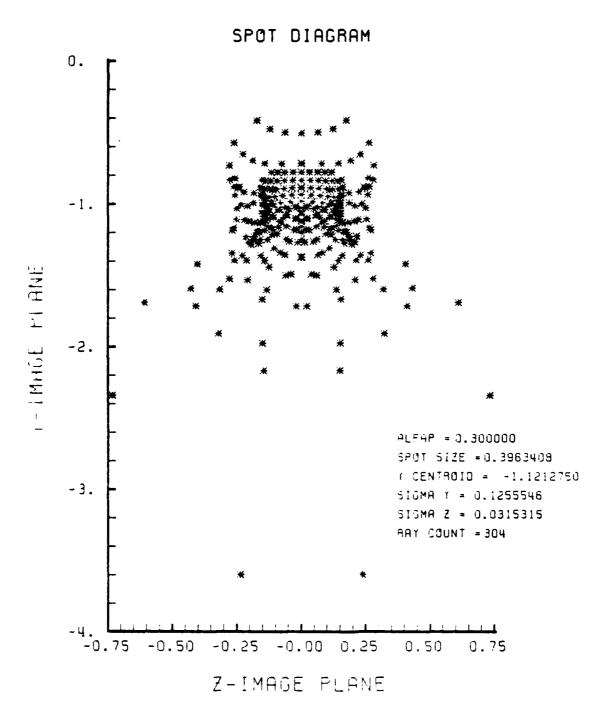


Figure E-75. Spot Diagram for Grid of Figure E-74

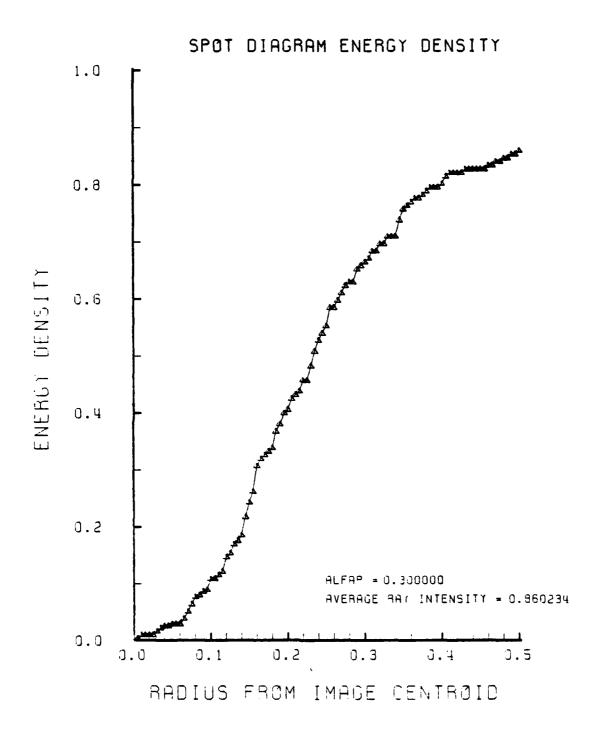


Figure E-76. Encircled Energy of Figure E-75

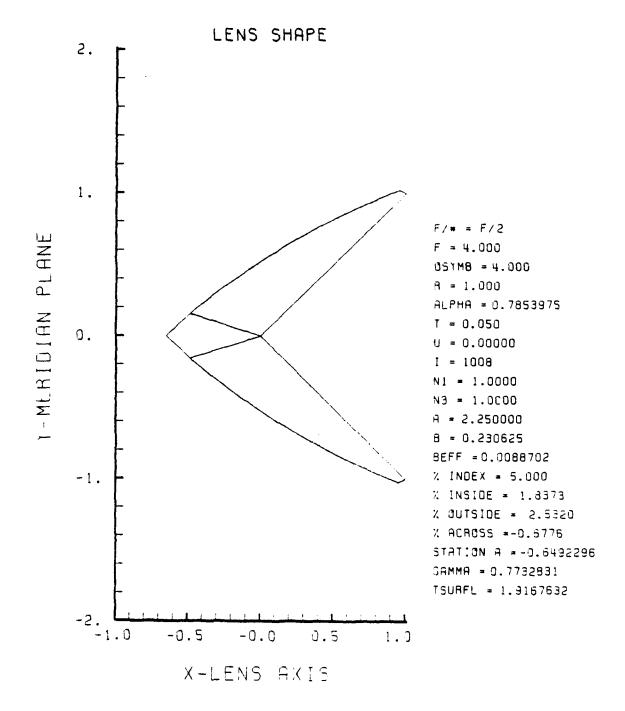
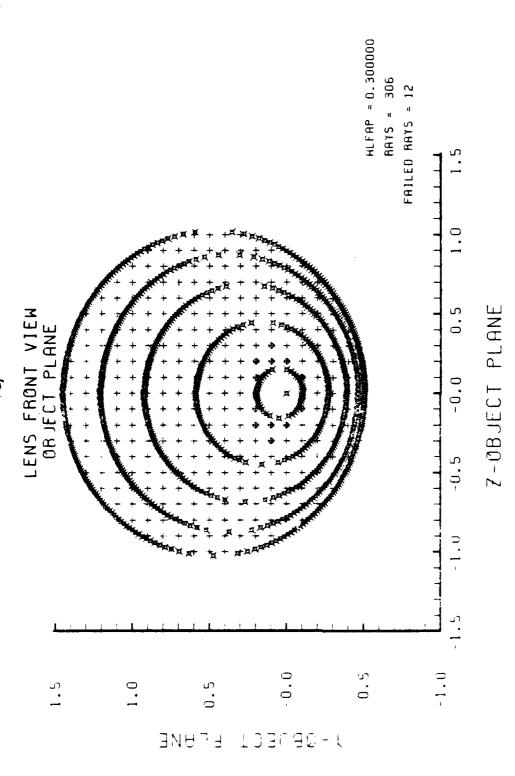
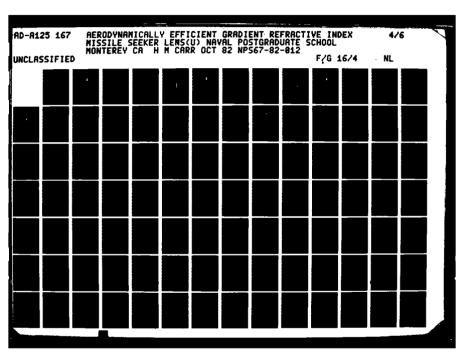
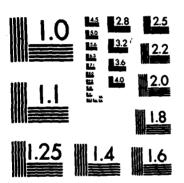


Figure E-77. GRIN Lens Shape at +5%, OB = 4.00, a = 2.25



Grid Plane at  $\alpha_{\rm p}=0.3$  for Lens of Figure E-77 Figure E-78.





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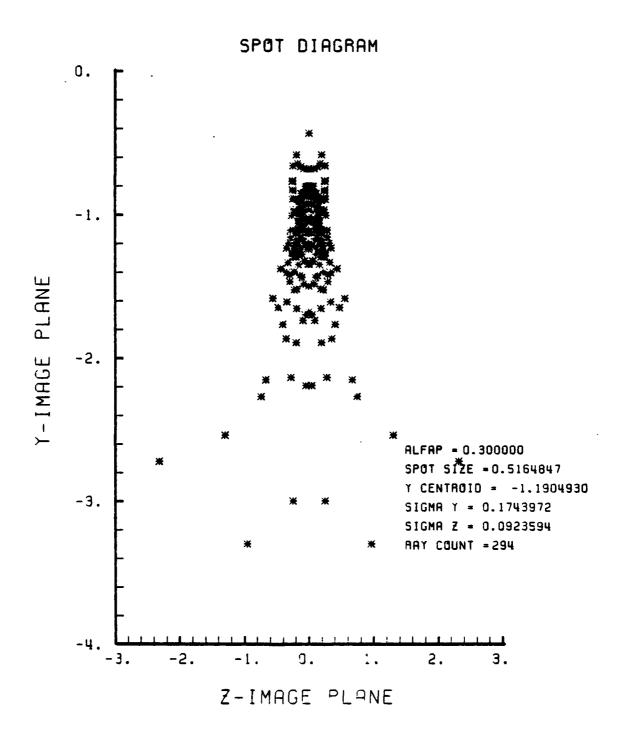


Figure E-79. Spot Diagram for Grid of Figure E-78

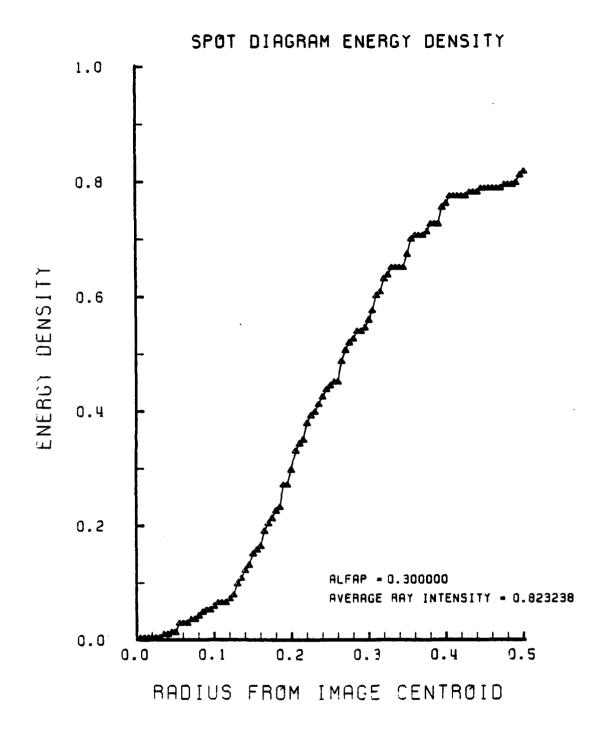


Figure E-80. Encircled Energy of Figure E-79

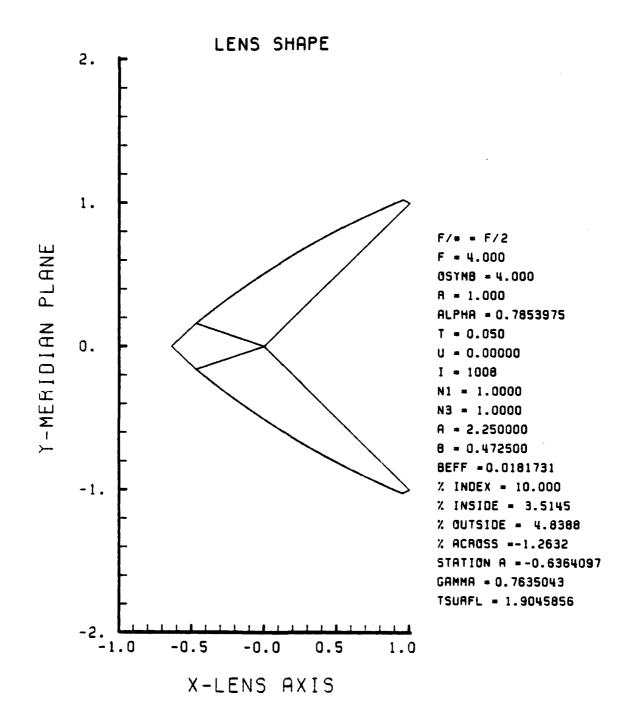
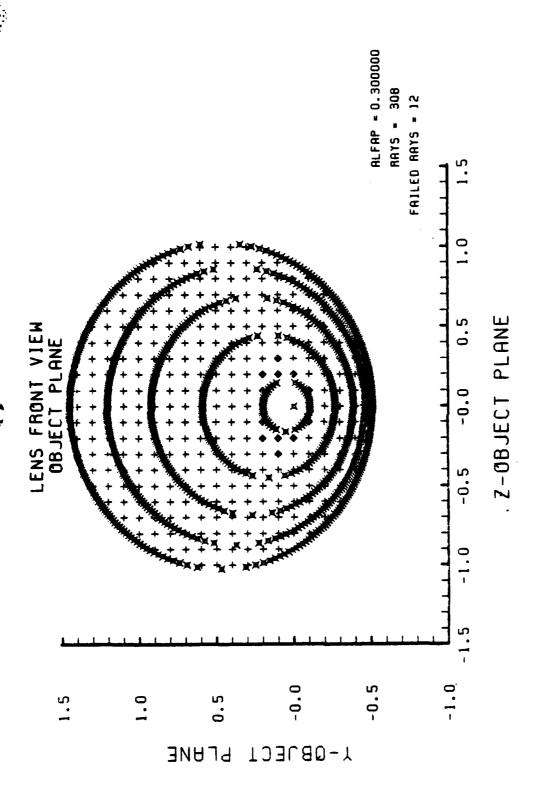


Figure E-81. GRIN Lens Shape at +10%, OB = 4.00, a = 2.25



Grid Plane at  $\alpha_p = 0.3$  for Lens of Figure E-81 Figure E-82.

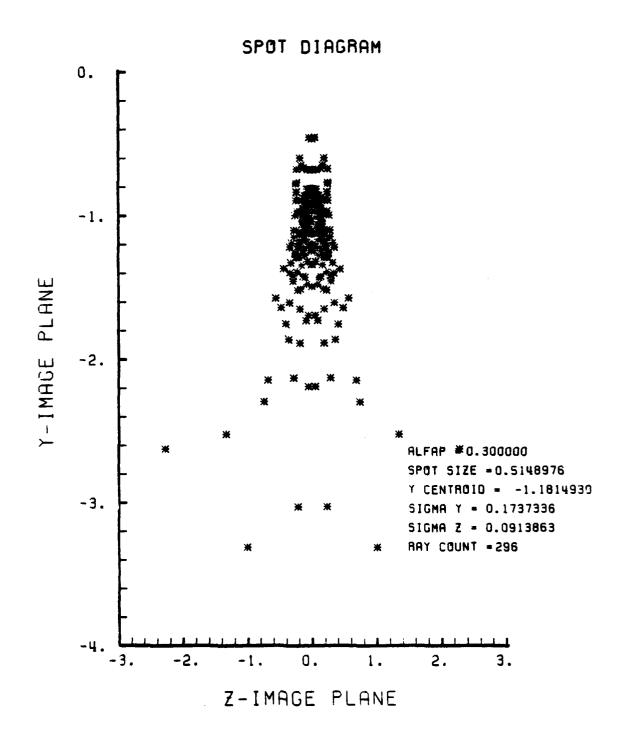


Figure E-83. Spot Diagram for Grid of Figure E-82

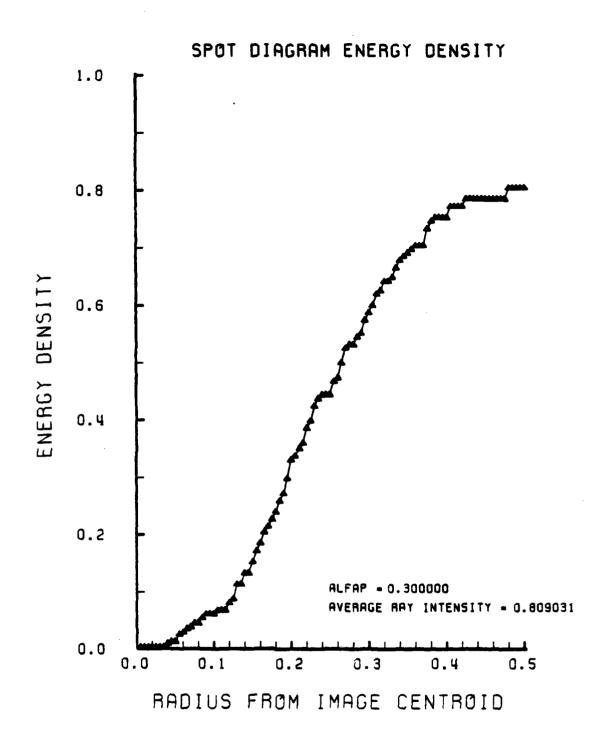


Figure E-84. Encircled Energy of Figure E-83

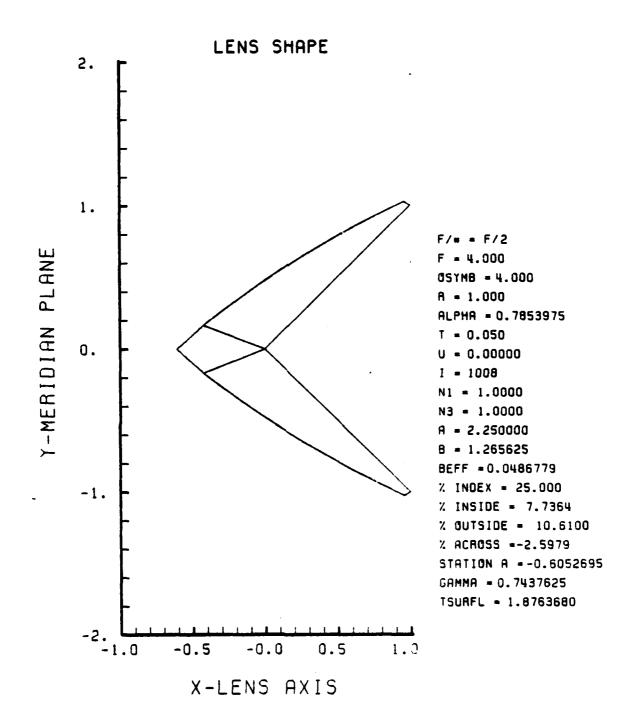
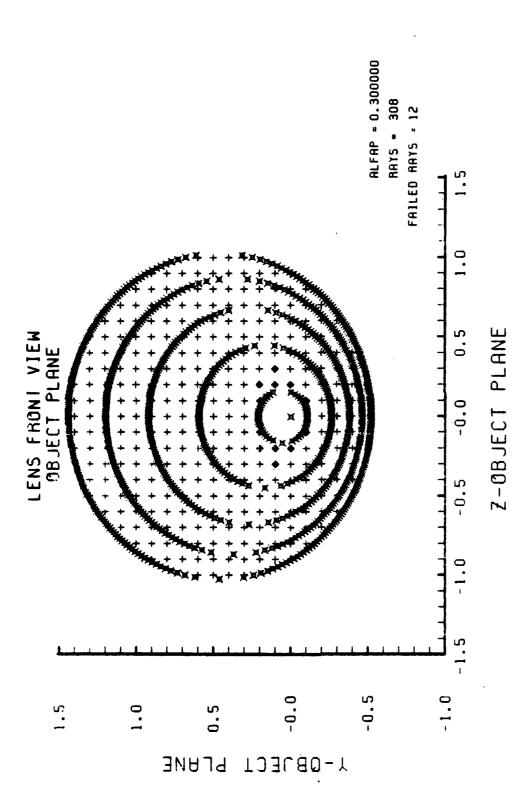


Figure E-85. GRIN Lens Shape at +25%, OB = 4.00, a = 2.25

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Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure E-85 Figure E-36.

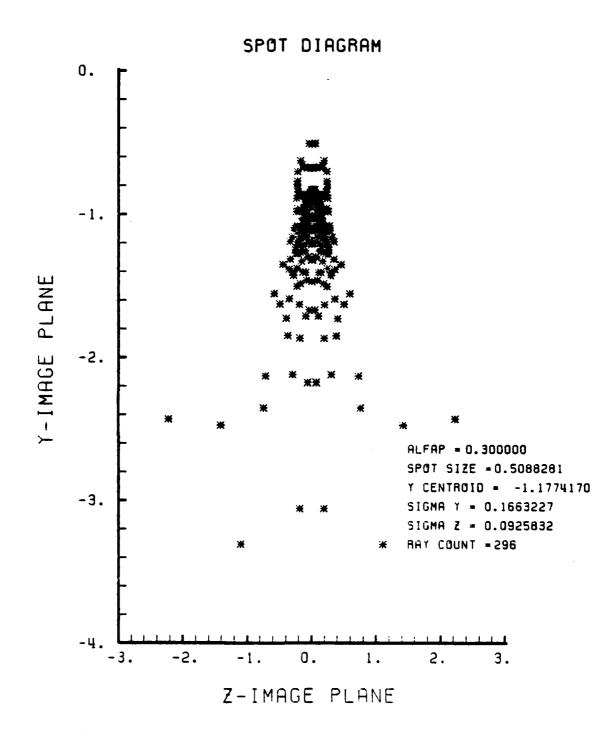


Figure E-87. Spot Diagram for Grid of Figure E-86

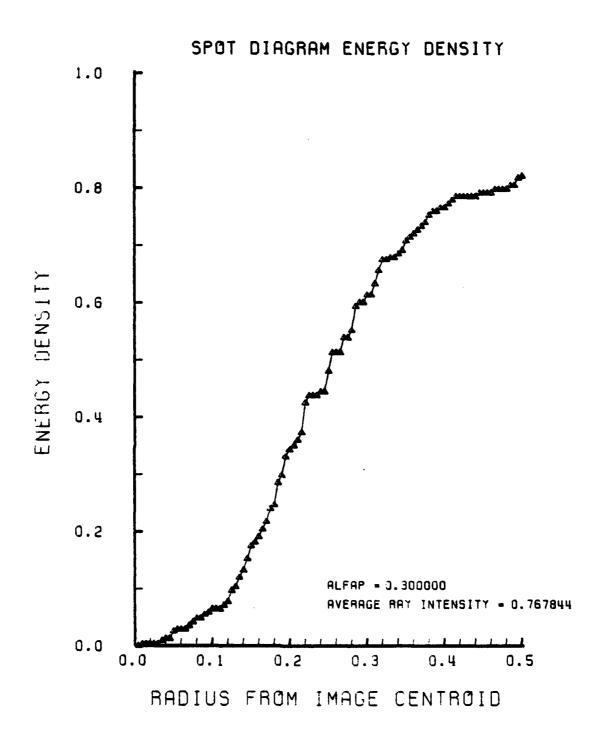


Figure E-88. Encircled Energy of Figure E-87

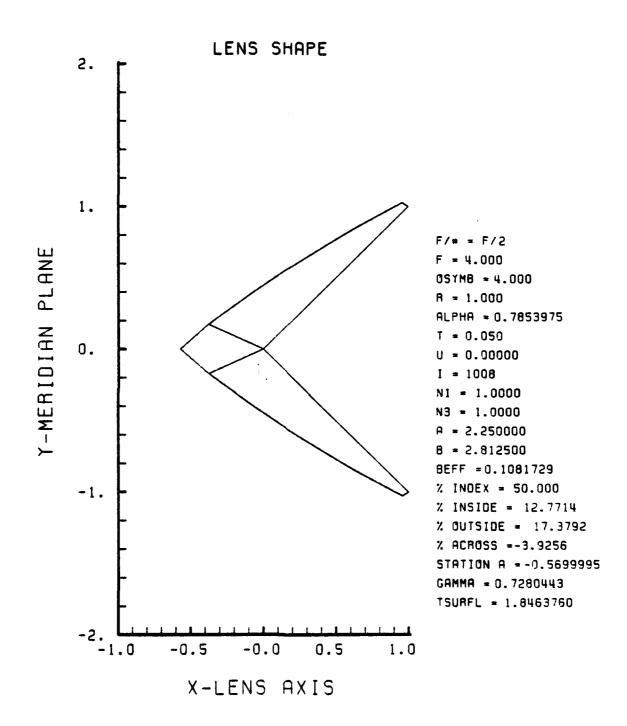
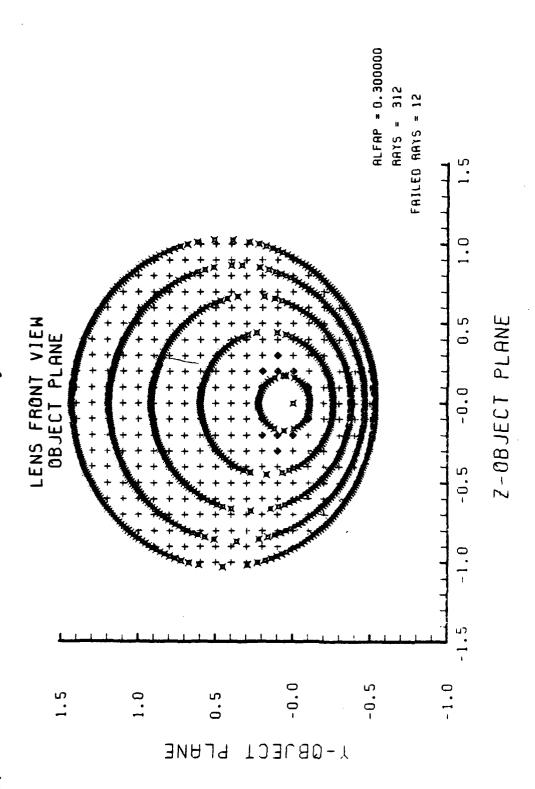


Figure E-89. GRIN Lens Shape at +50%, OB = 4.00, a = 2.25



Grid Plane at  $\alpha_{\rm p}=0.3$  for Lens of Figure E-89 Figure E-90.

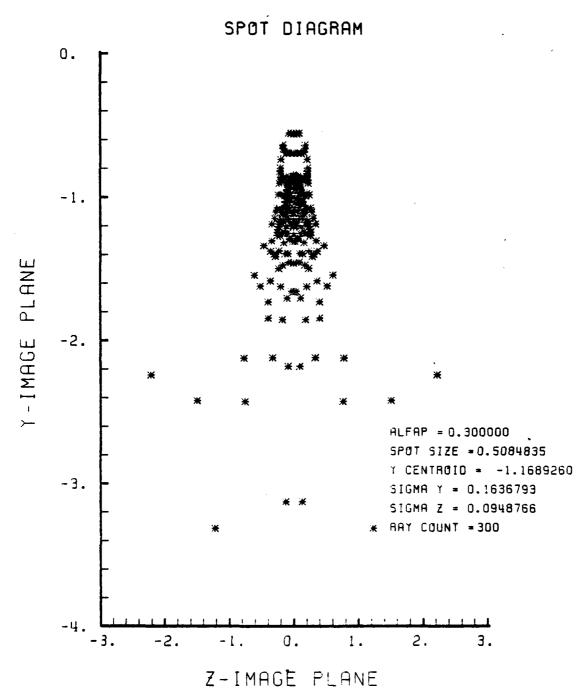


Figure E-91. Spot Diagram for Grid of Figure E-90

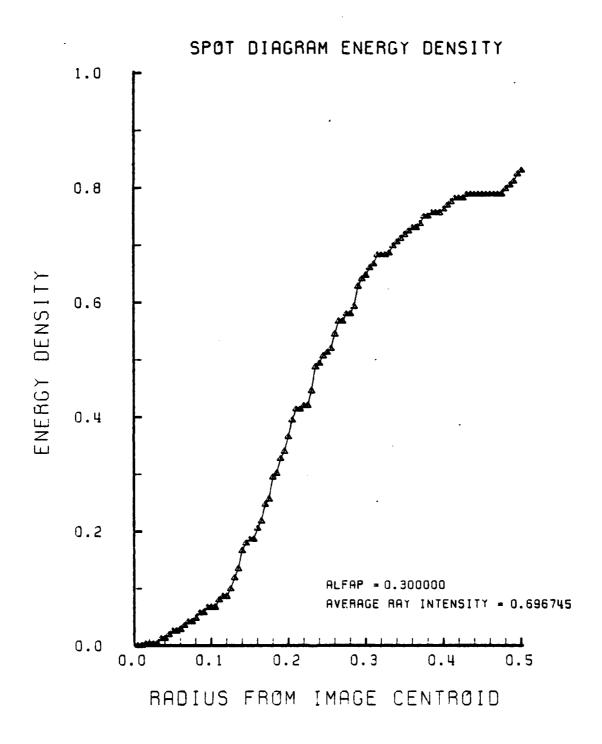


Figure E-92. Encircled Energy of Figure E-91

## APPENDIX F

## GRIN LENS PERFORMANCE PLOTS IN THE HIGH RANGE OF INDICES OF REFRACTION (a = 9.00)

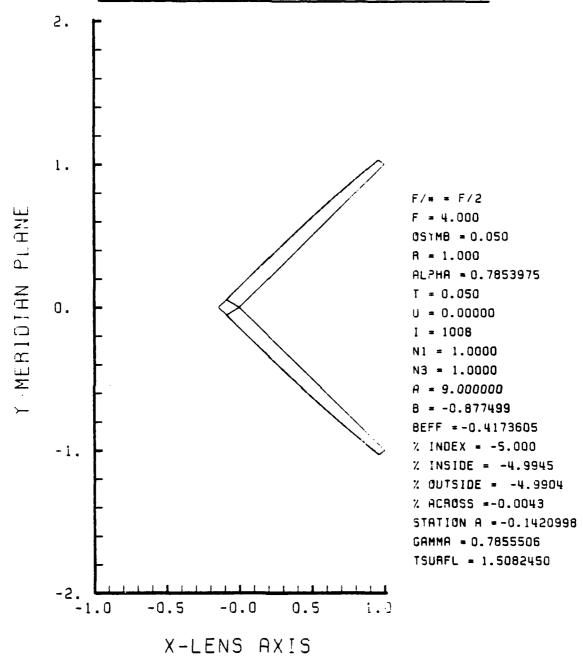
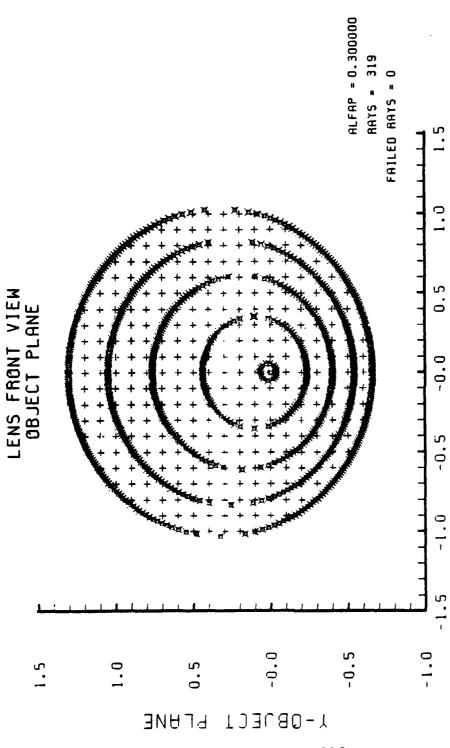


Figure F-1. GRIN Lens Shape at -5%, OB = 0.05, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-1 Figure F-2.

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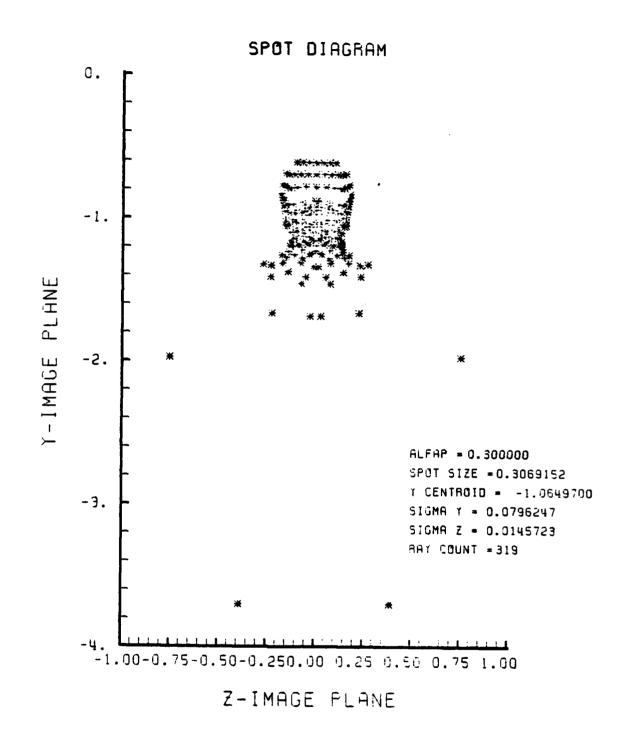


Figure F-3. Spot Diagram for Grid of Figure F-2

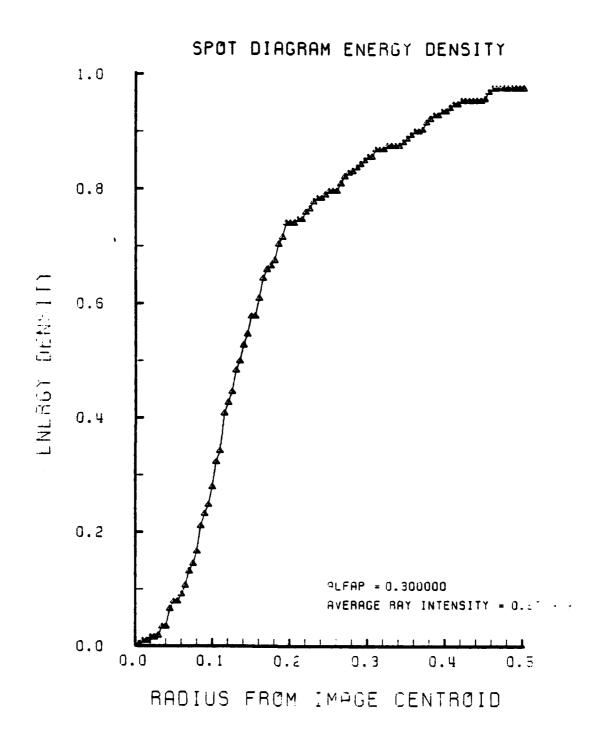


Figure F-4. Encircled Energy of Figure F-3

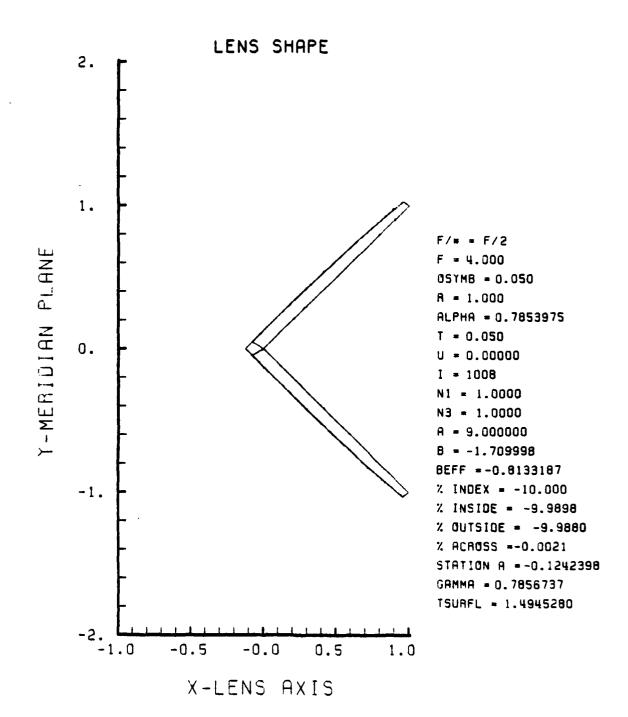


Figure F-5. GRIN Lens Shape at -10%, OB = 0.05, a = 9.00

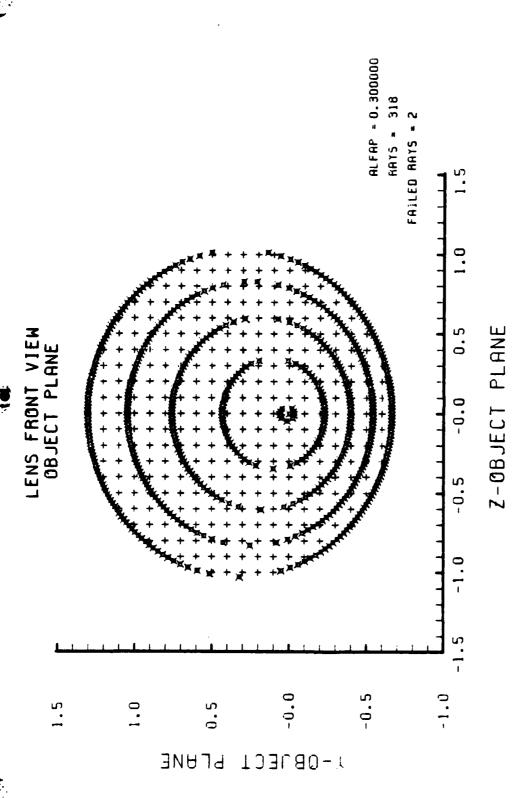


Figure F-6. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-5

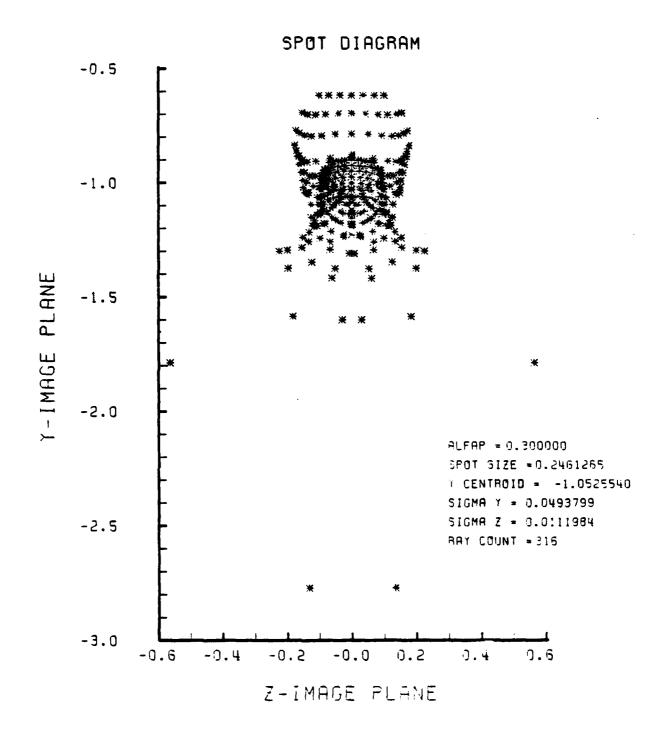


Figure F-7. Spot Diagram for Grid of Figure F-6

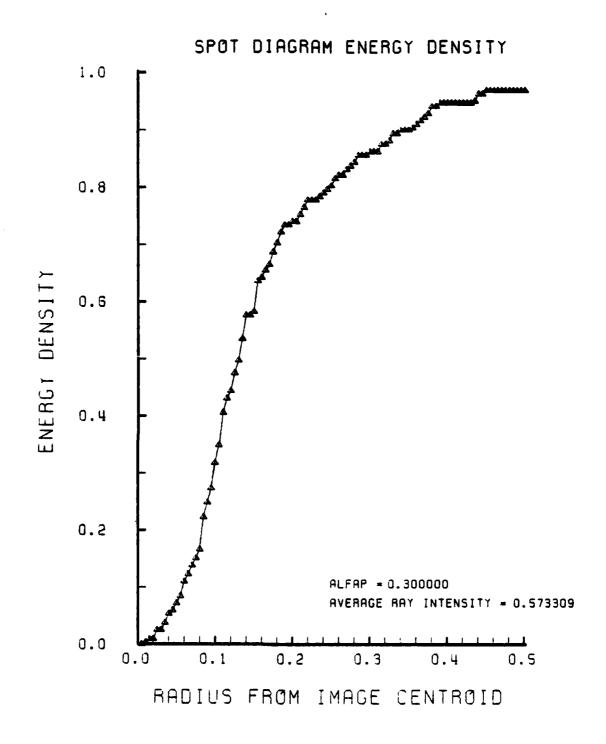


Figure F-8. Encircled Energy of Figure F-7

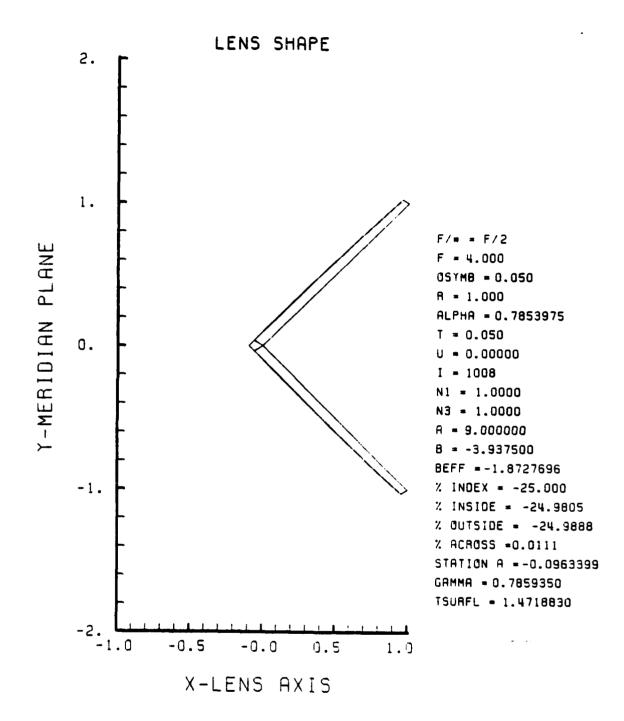
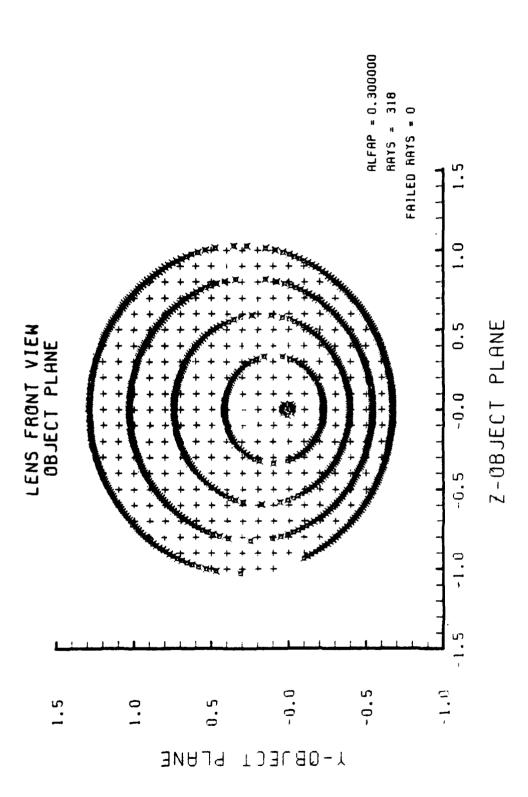


Figure F-9. GRIN Lens Shape at -25%, OB = 0.05, a = 9.00



Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure F-9 Figure F-10.

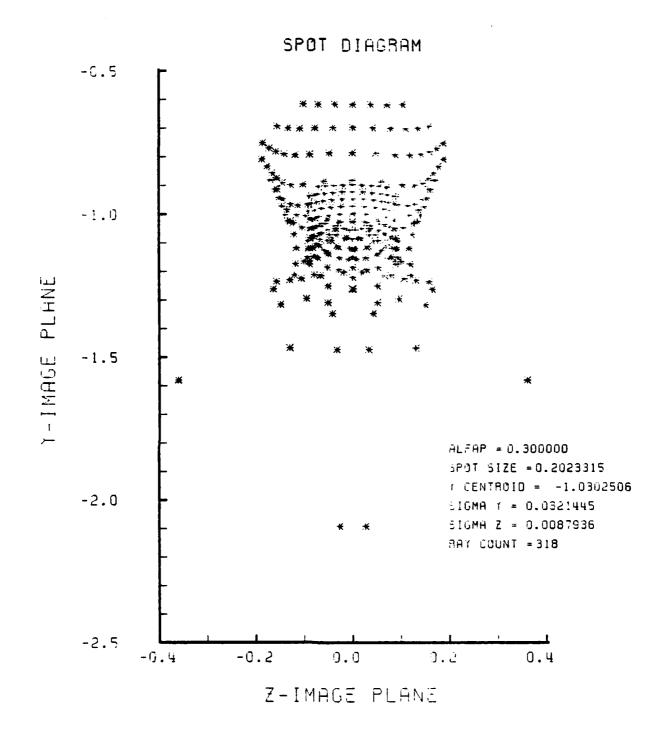


Figure F-11. Spot Diagram for Grid of Figure F-10

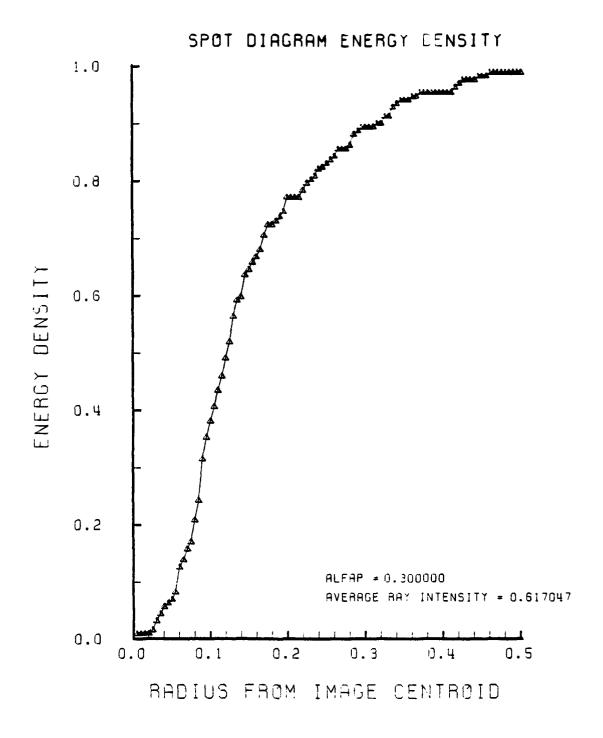


Figure F-12. Encircled Energy of Figure F-11

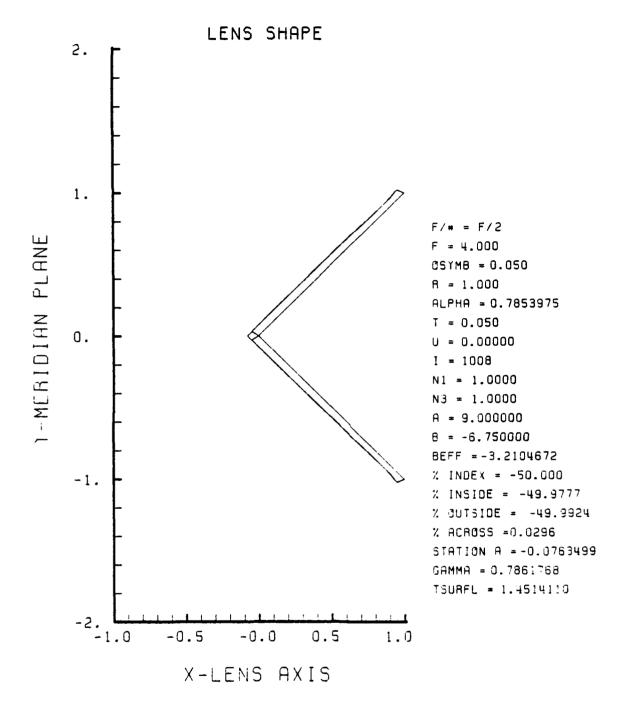
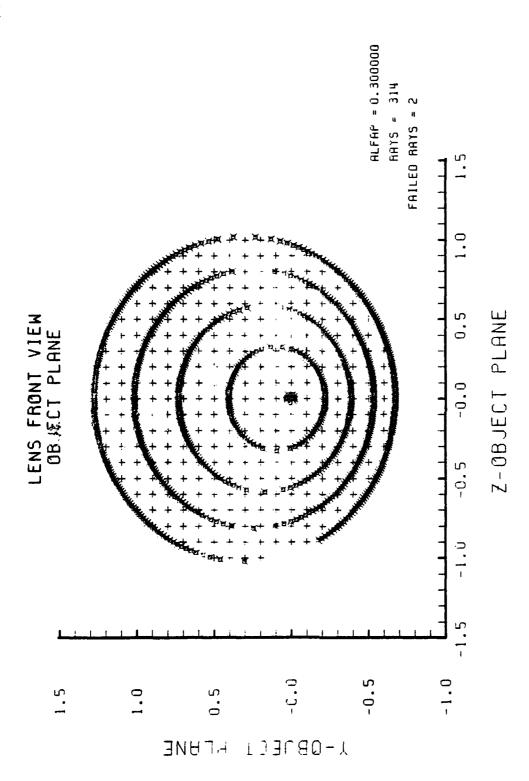


Figure F-13. GRIN Lens Shape at -50%, OB = 0.05, a = 9.00



Grid Plane at  $\alpha_p = 0.3$  for Lens of Figure F-13 Figure F-14.

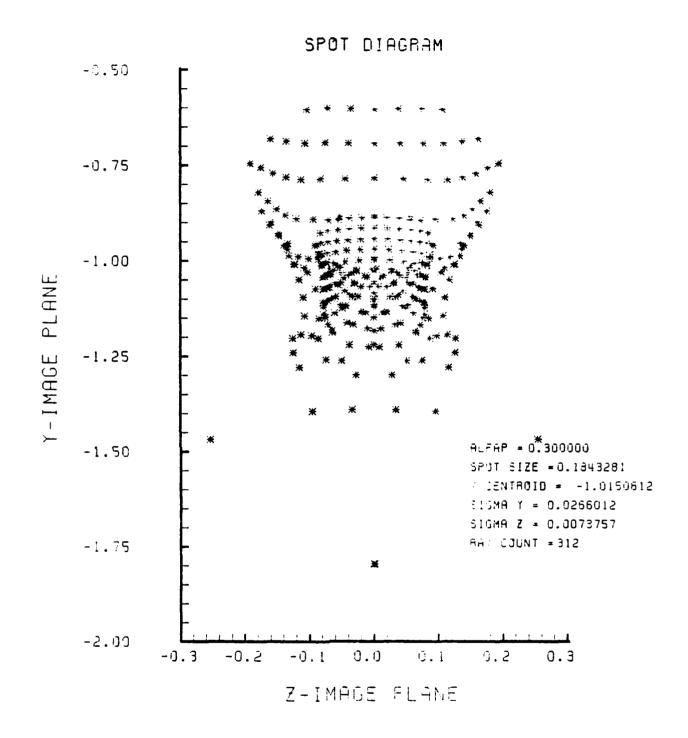


Figure F-15. Spot Diagram for Grid of Figure F-14

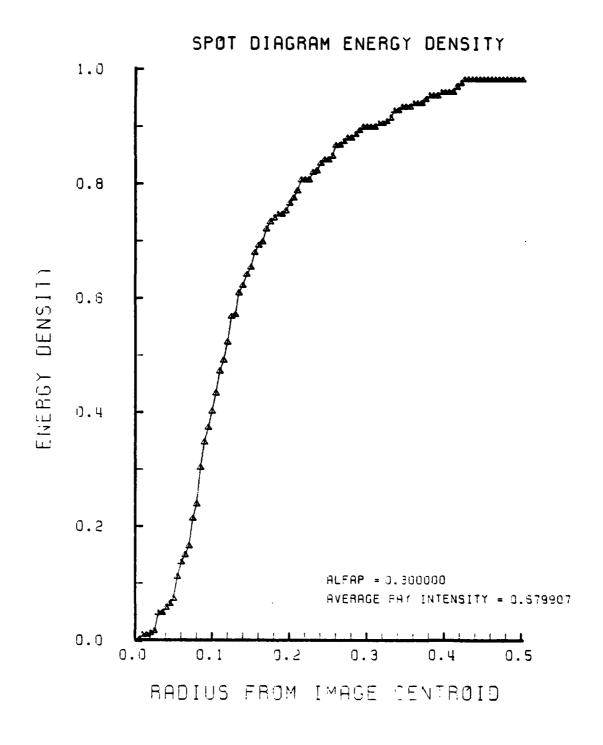


Figure F-16. Encircled Energy of Figure F-15

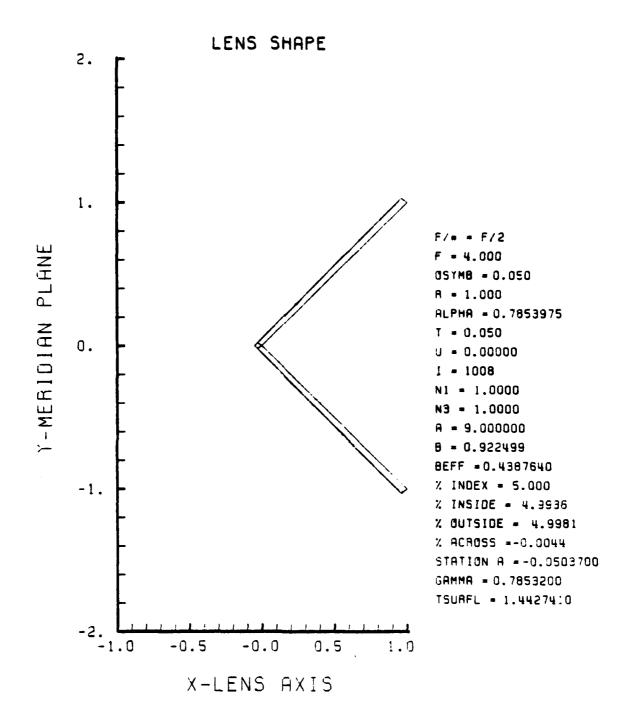
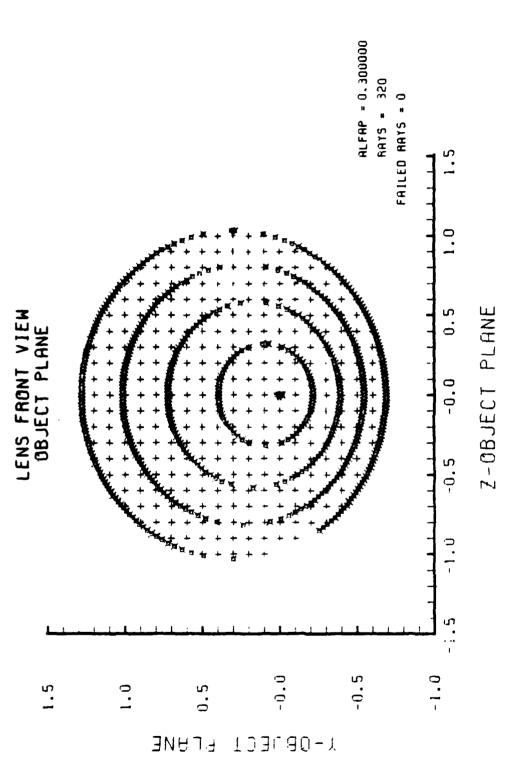


Figure F-17. GRIN Lens Shape at +5%, OB = 0.05, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-17 Figure F-18.

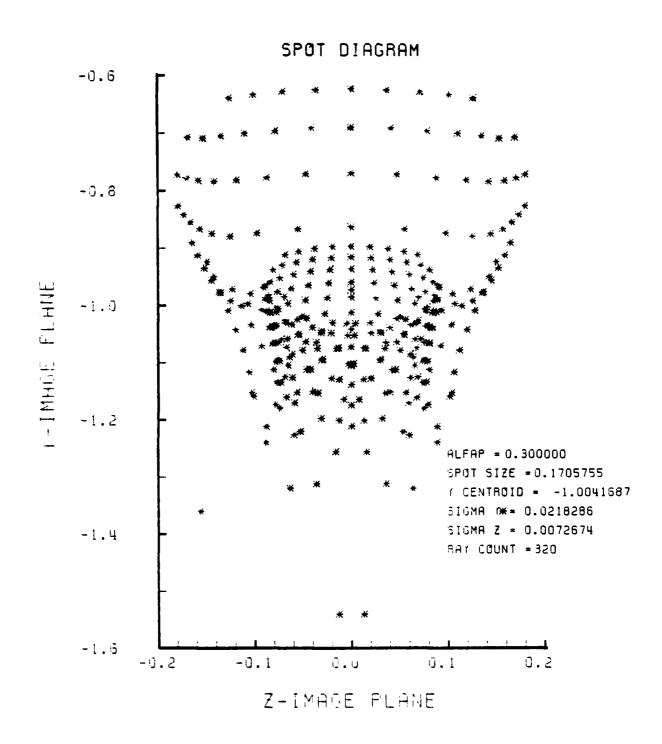


Figure F-19. Spot Diagram for Grid of Figure F-18

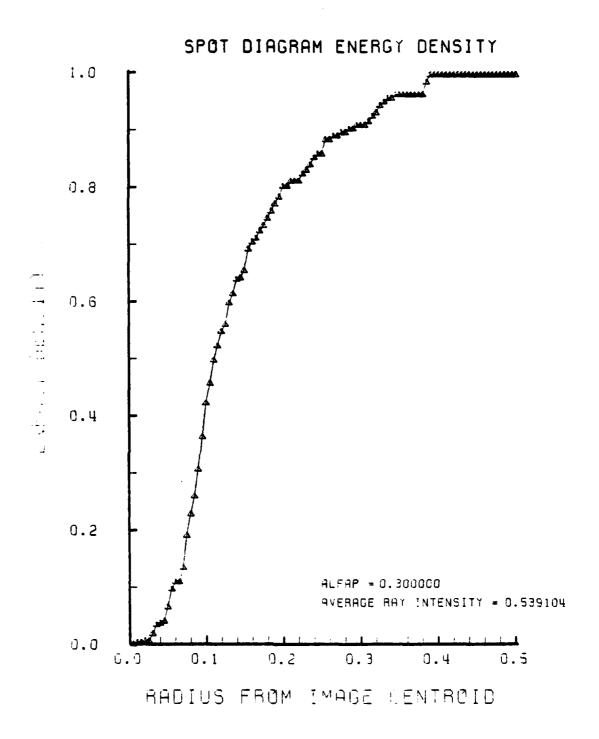


Figure F-20. Encircled Energy of Figure F-19

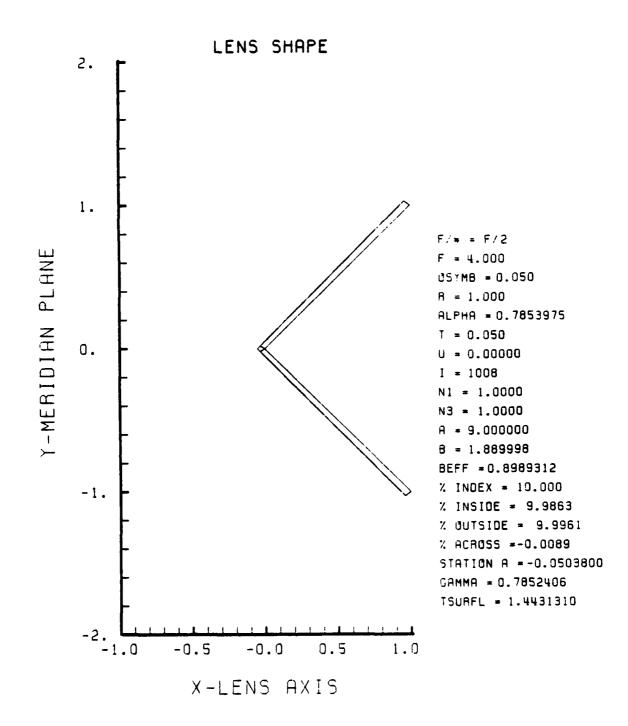
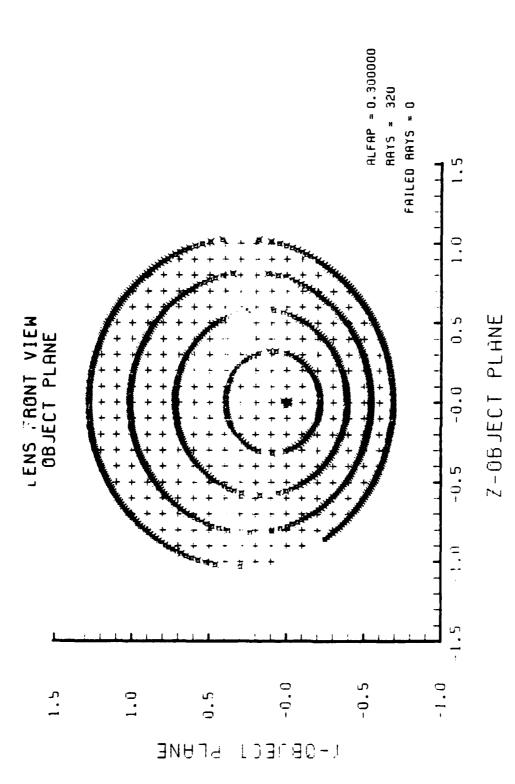


Figure F-21. GRIN Lens Shape for +10%, OB = 0.05, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 Radians for Lens of Figure F-21 Figure F-22.

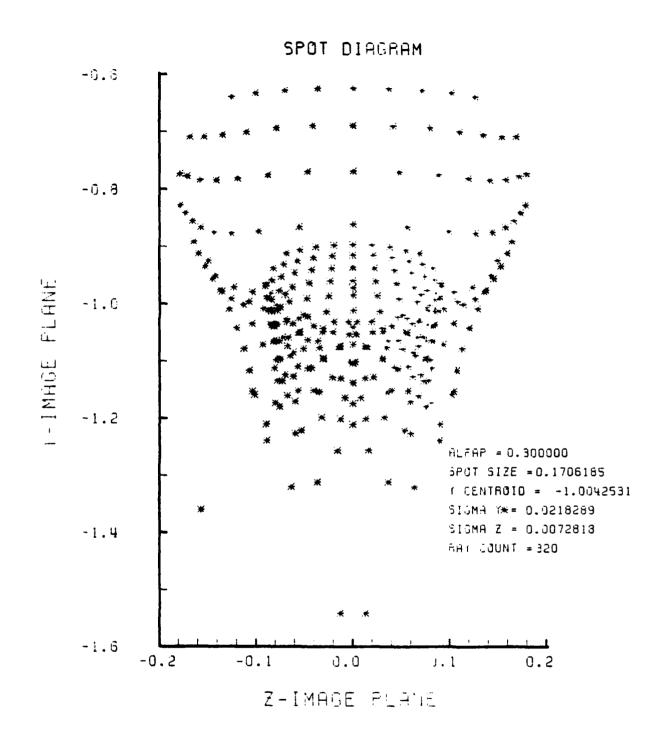


Figure F-23. Spot Diagram for Grid of Figure F-22

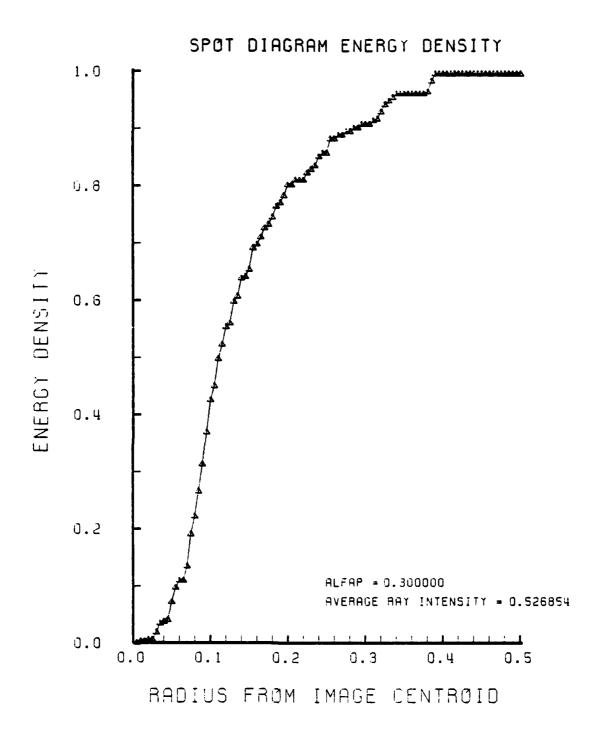


Figure F-24. Encircled Energy of Figure F-23

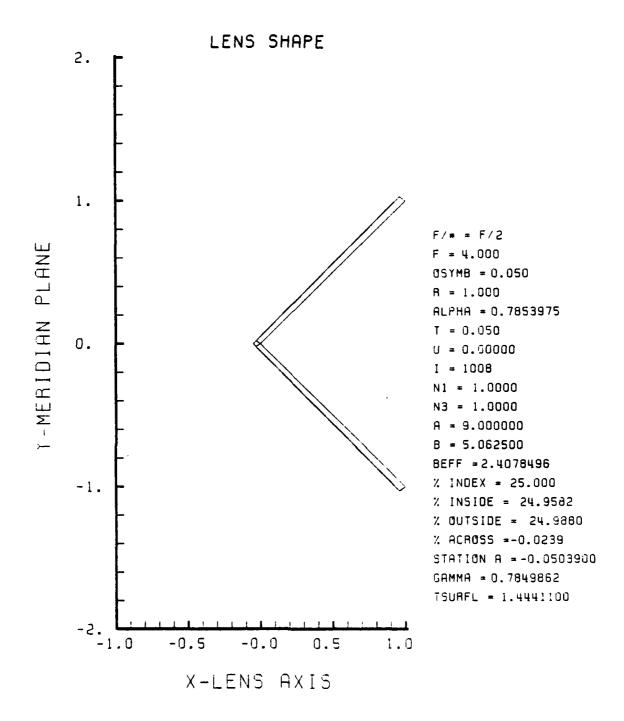


Figure F-25. GRIN Lens Shape for +25%, OB = 0.05, a = 9.00

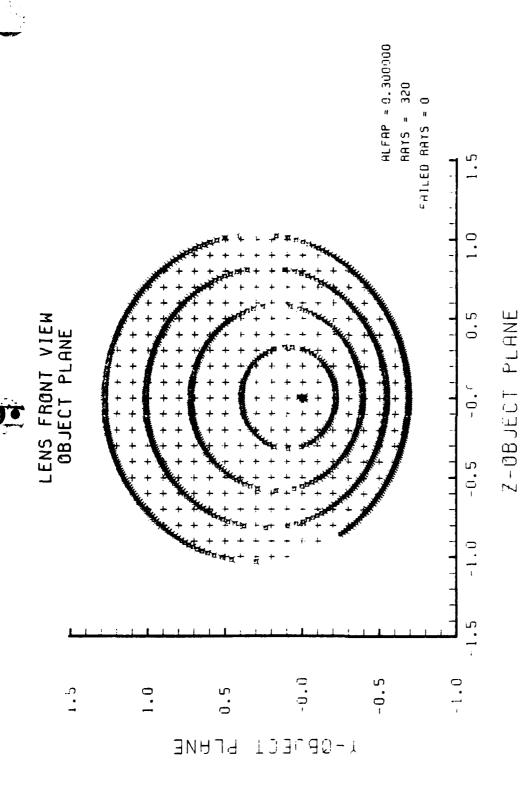


Figure F-26. Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure F-25

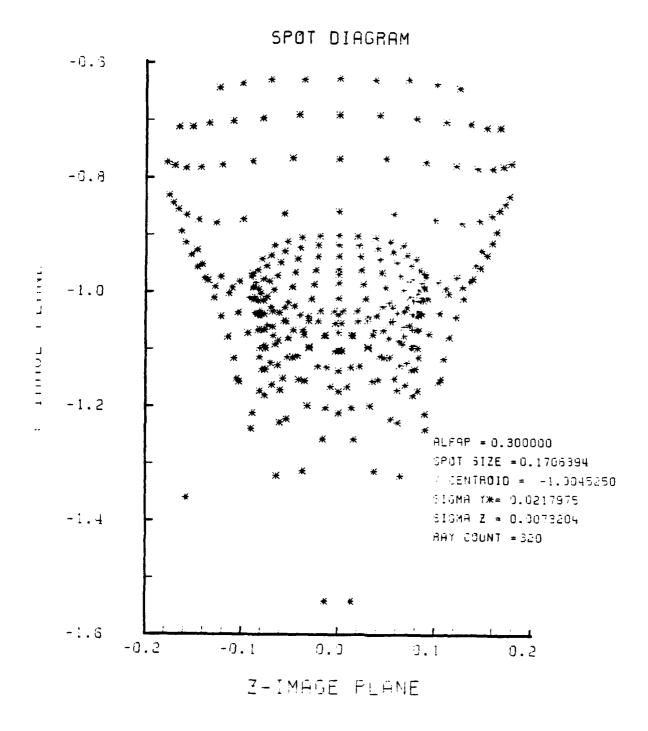


Figure F-27. Spot Diagram for Grid of Figure F-26

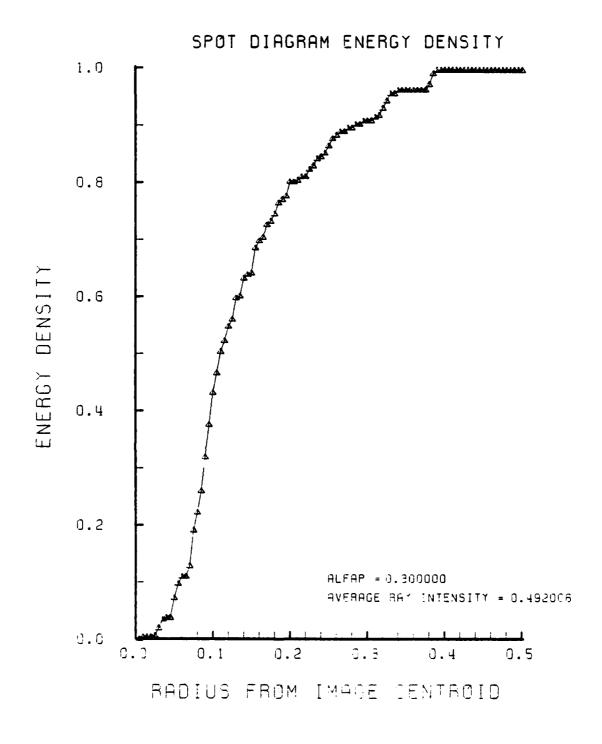


Figure F-28. Encircled Energy of Figure F-27

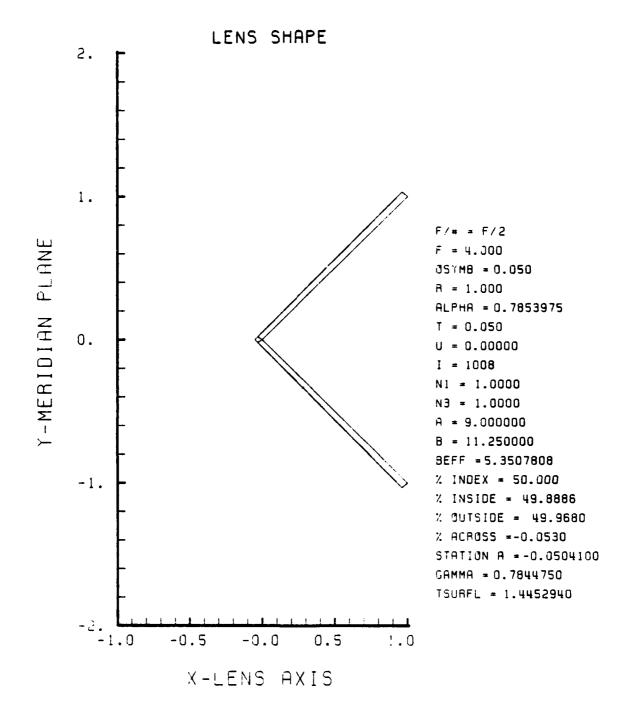


Figure F-29. GRIN Lens Shape at +50%, OB = 0.05, a = 9.00

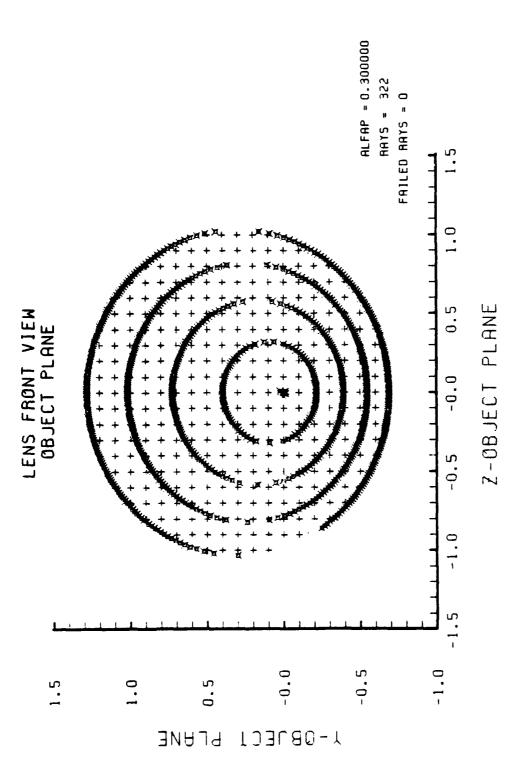


Figure F-30. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-29

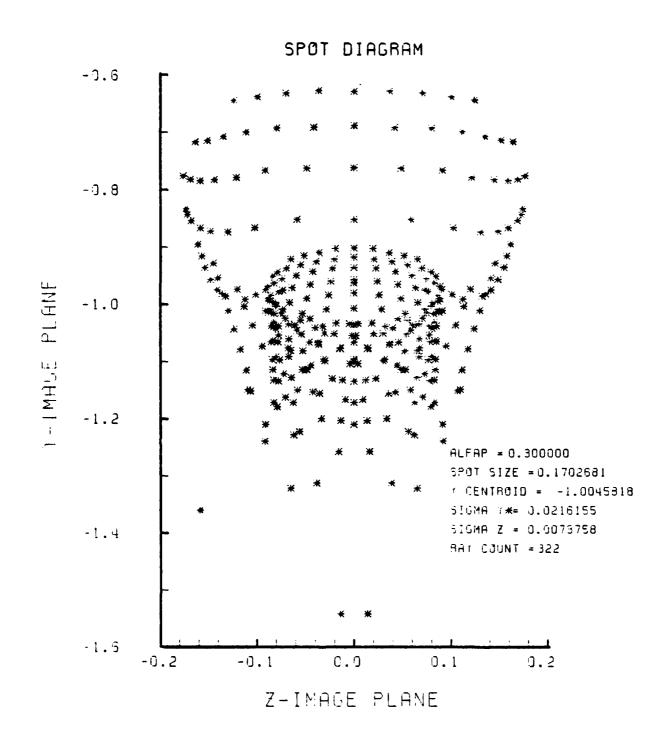


Figure F-31. Spot Diagram for Grid of Figure F-30

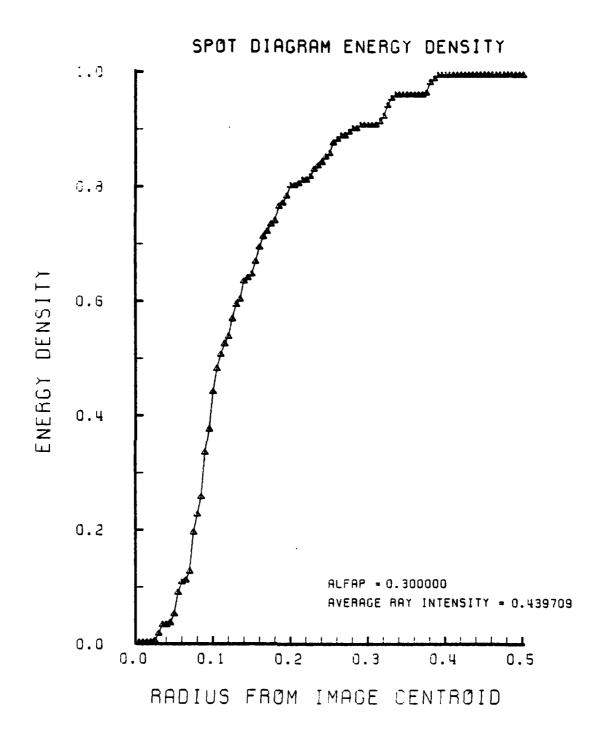


Figure F-32. Encircled Energy of Figure F-31

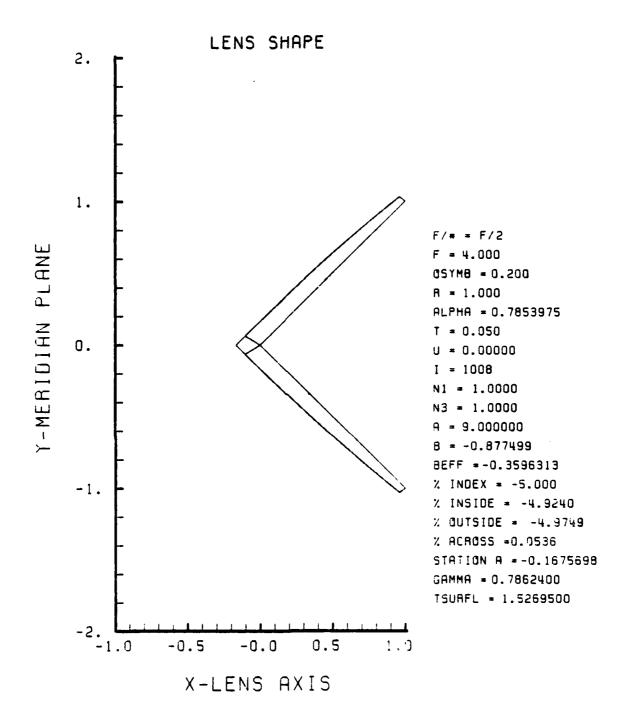
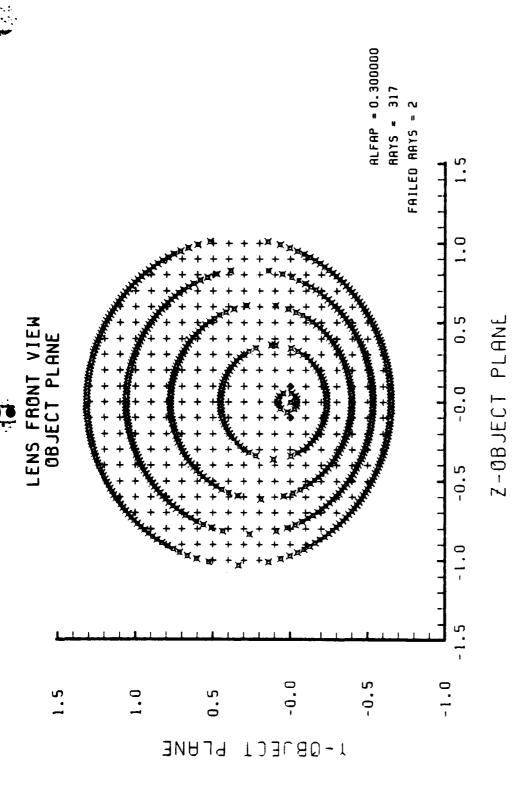


Figure F-33. GRIN Lens Shape at -5%, OB = 0.20, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-33 Figure F-34.

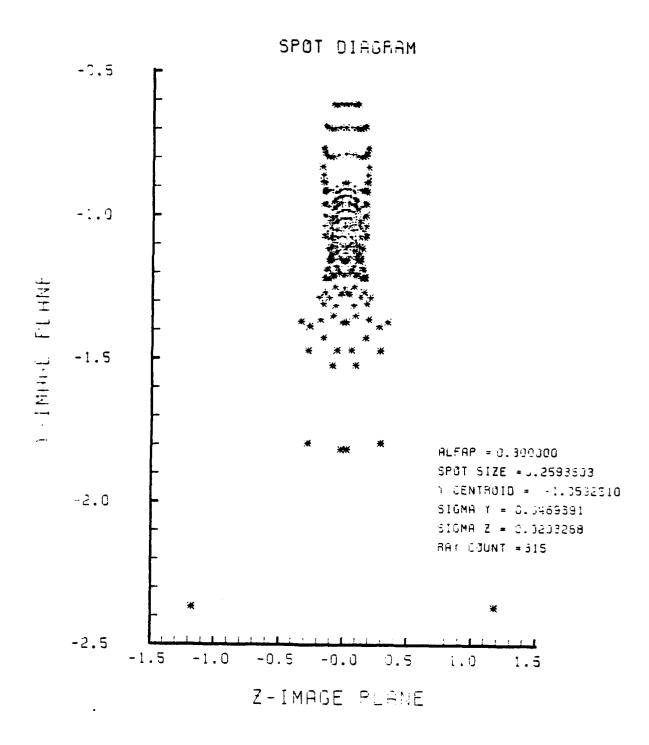


Figure F-35. Spot Diagram for Grid of Figure F-34

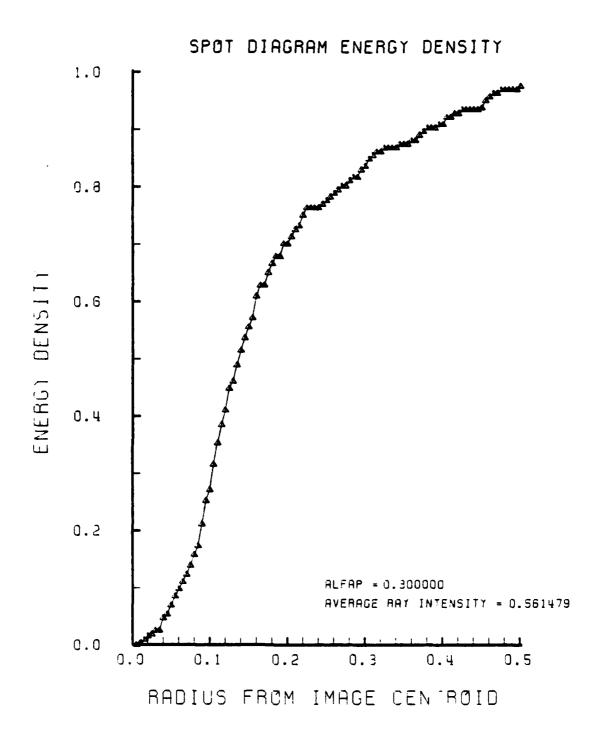


Figure F-36. Encircled Energy of Figure F-35

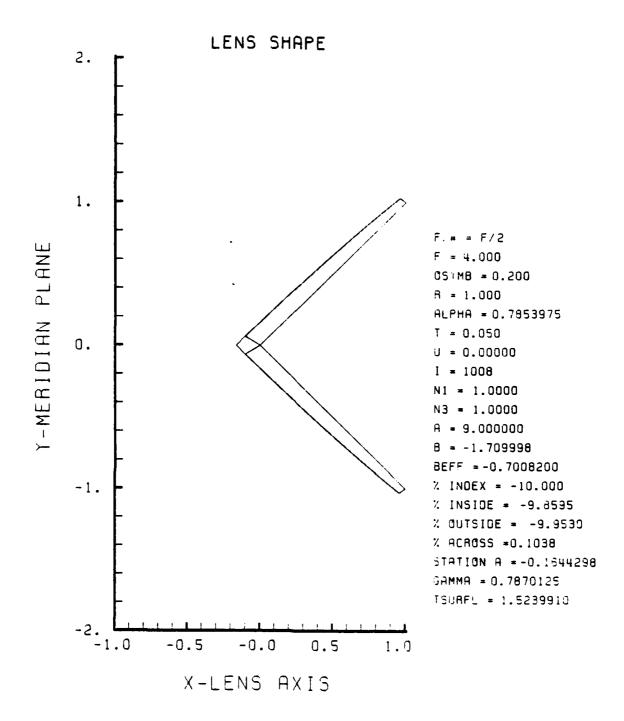
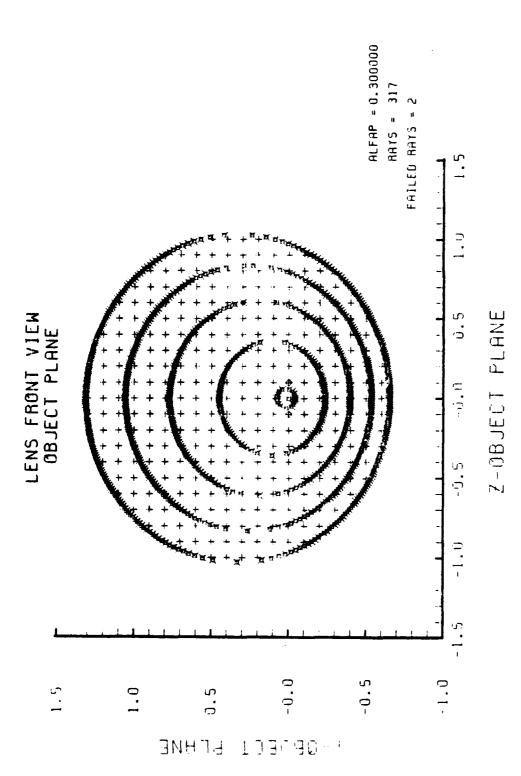


Figure F-37. GRIN Lens Shape at -10%, OB = 0.20, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-37 Figure F-38.

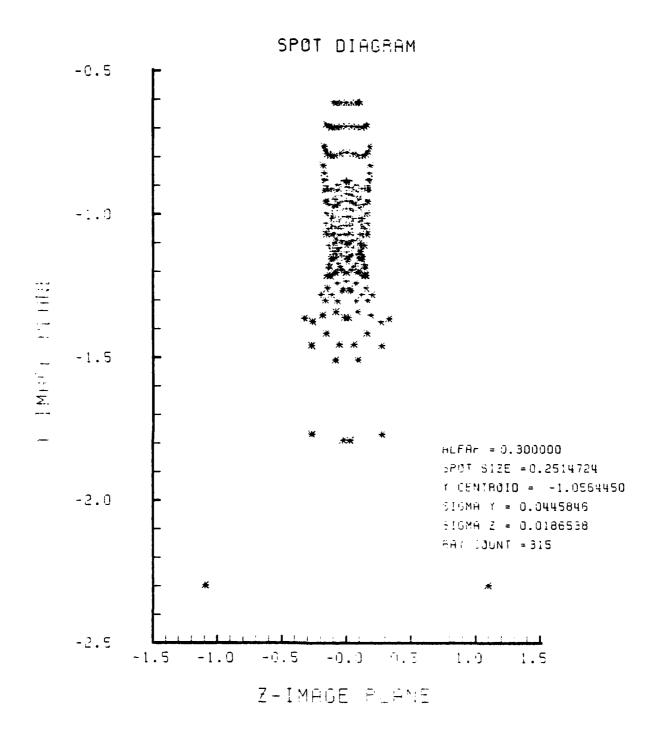


Figure F-39. Spot Diagram for Grid of Figure F-38

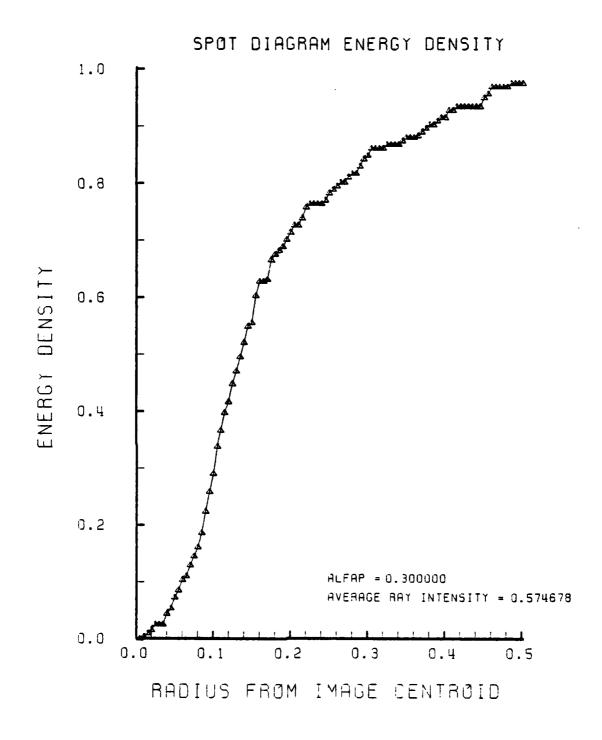


Figure F-40. Encircled Energy of Figure F-39

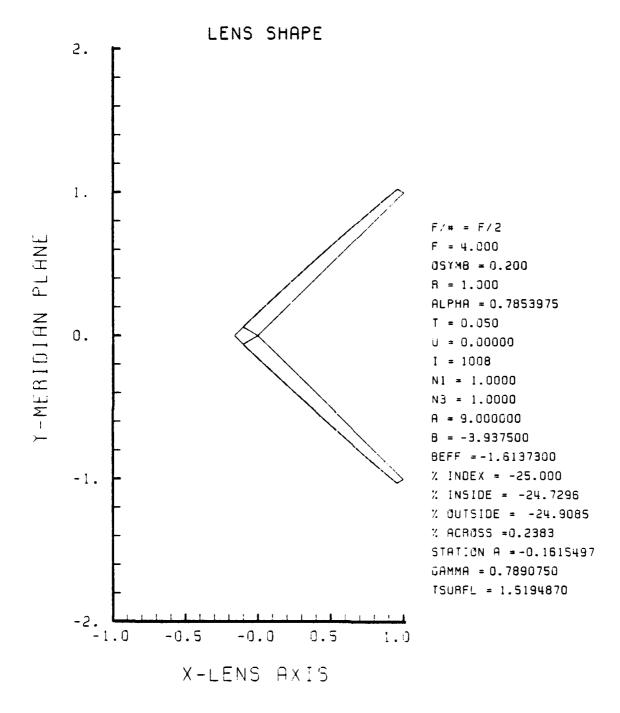


Figure F-41. GRIN Lens Shape at -25%, OB = 0.20, a = 9.00

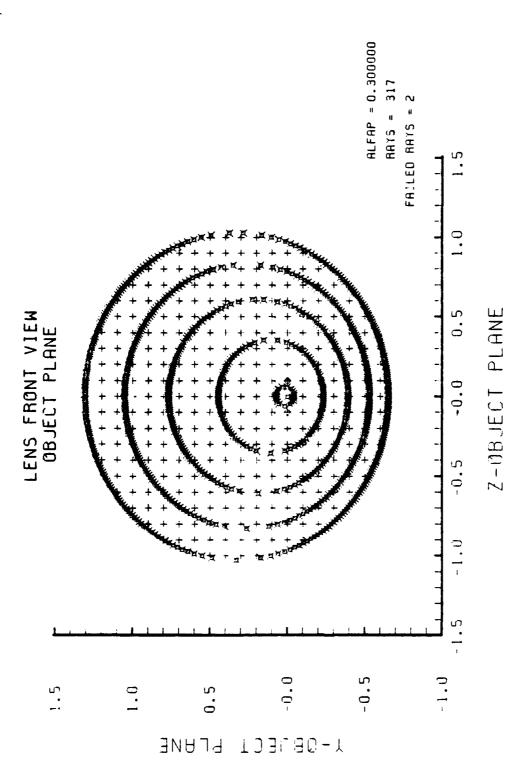


Figure F-42. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-41

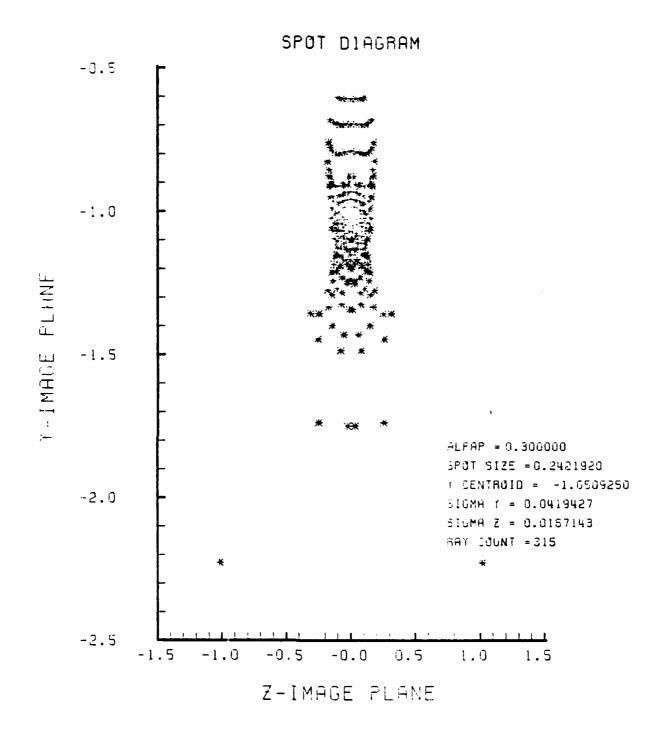


Figure F-43. Spot Diagram for Grid of Figure F-42

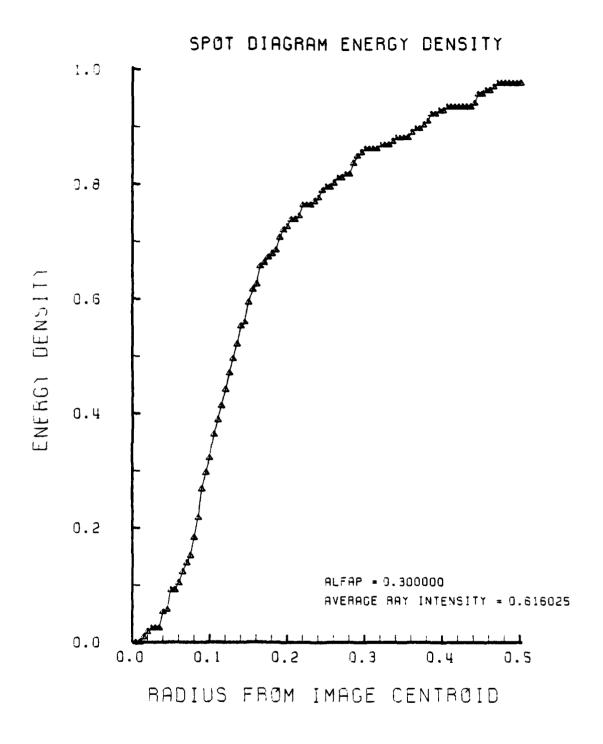


Figure F-44. Encircled Energy of Figure F-43

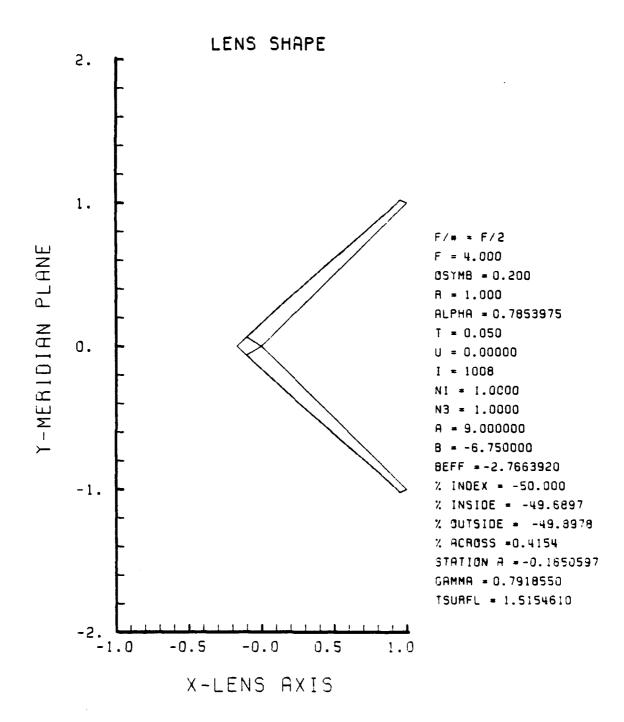


Figure F-45. GRIN Lens Shape for -50%, OB = 0.20, a = 9.00

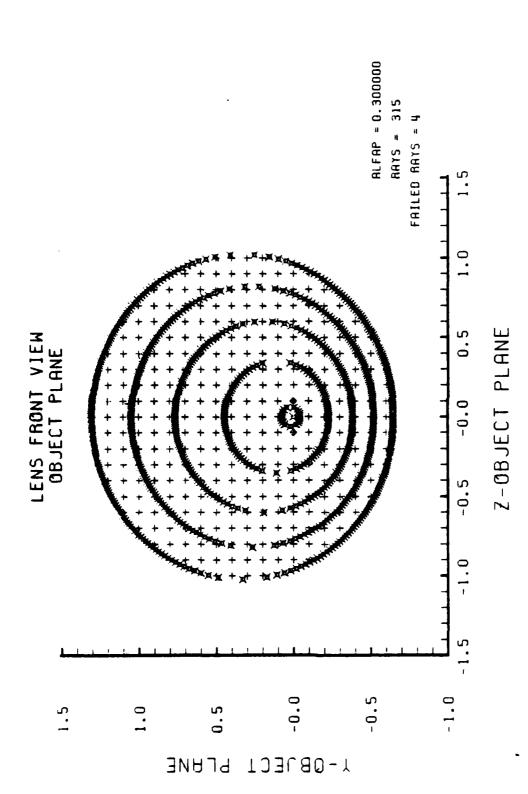


Figure F-46. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-45

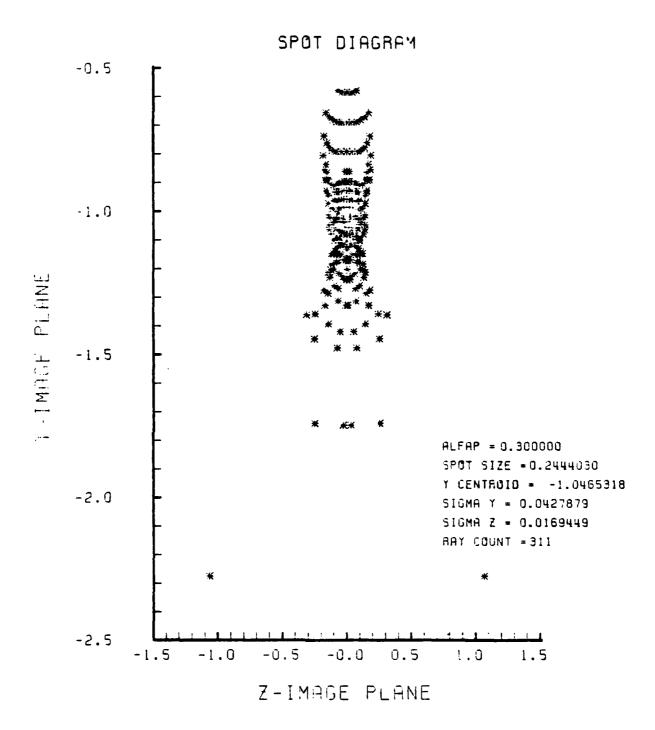


Figure F-47. Spot Diagram for Grid of Figure F-46

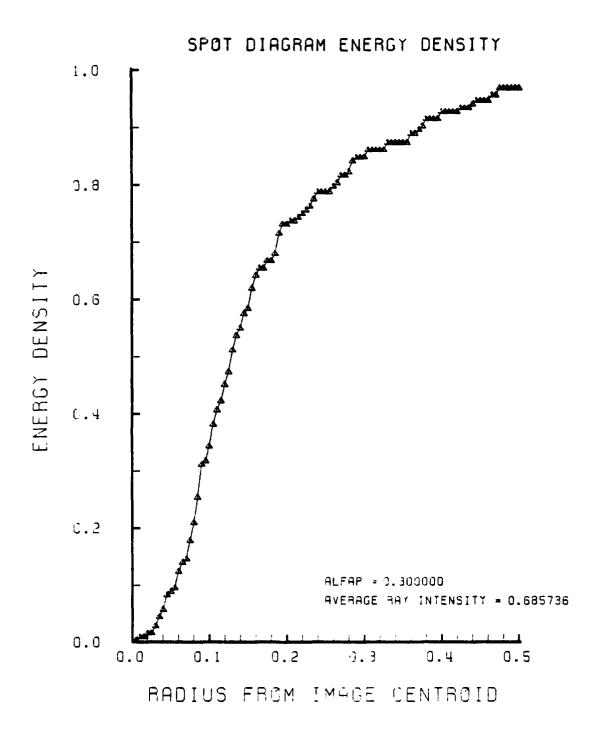


Figure F-48. Encircled Energy of Figure F-47

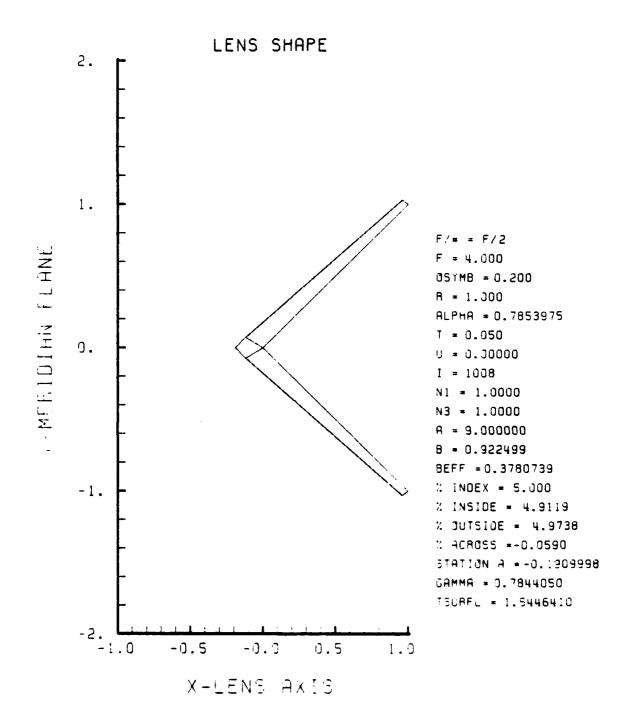
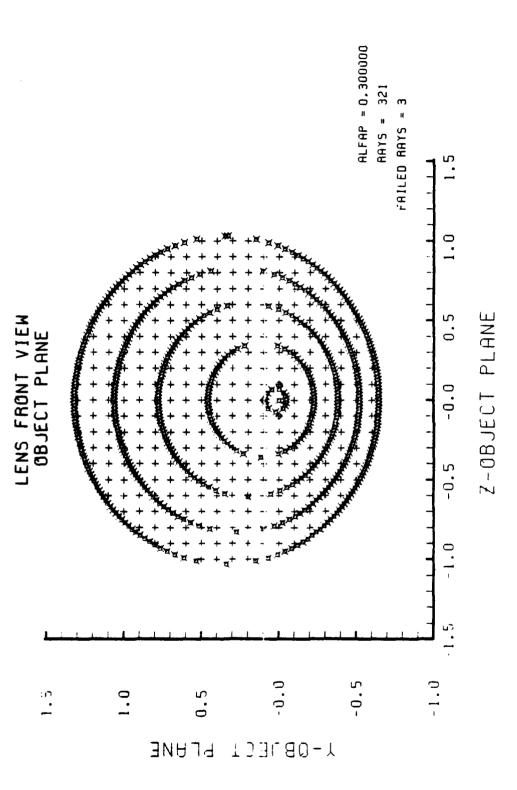


Figure F-49. GRIN Lens Shape for +5%, OB = 0.20, a = 9.00



Grid Plane at  $\alpha_{p}$  = 0.3 for Lens of Figure F-49 Figure F-50.

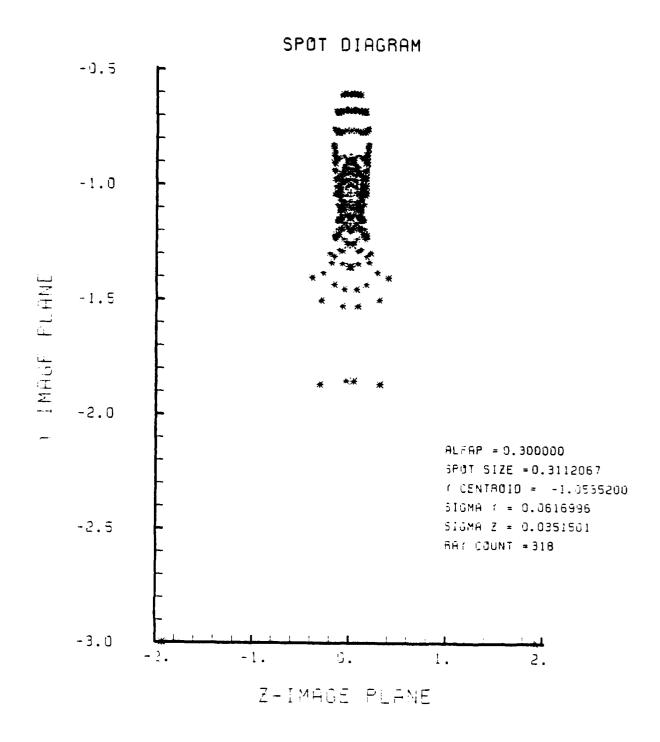


Figure F-51. Spot Diagram for Grid of Figure F-50

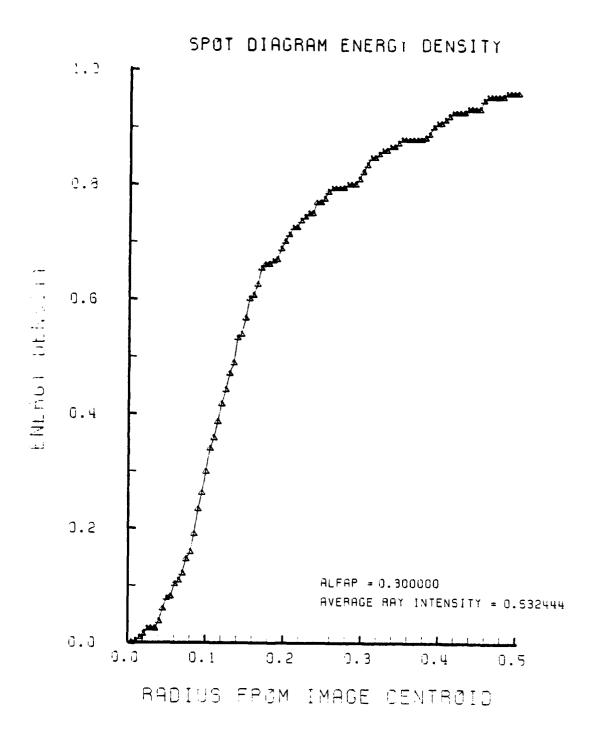


Figure F-52. Encircled Energy of Figure F-51

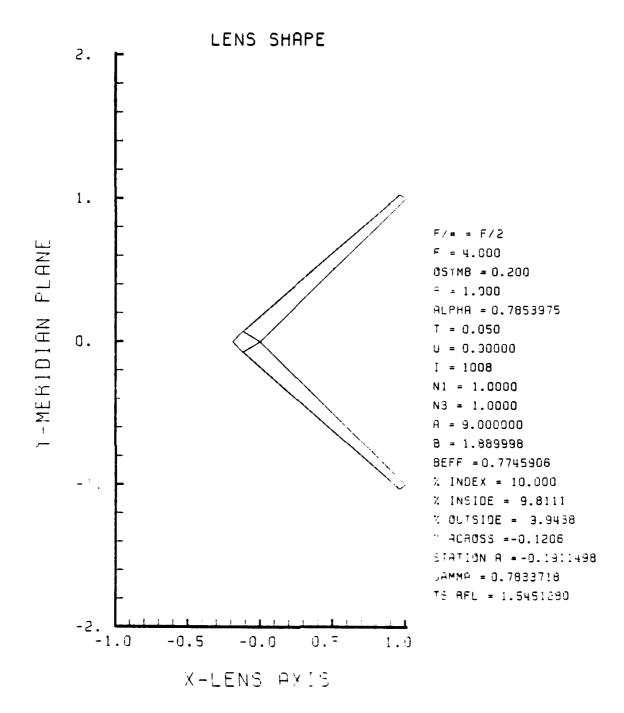


Figure F-53. GRIN Lens Shape at +10%, OB = 0.20, a = 9.00

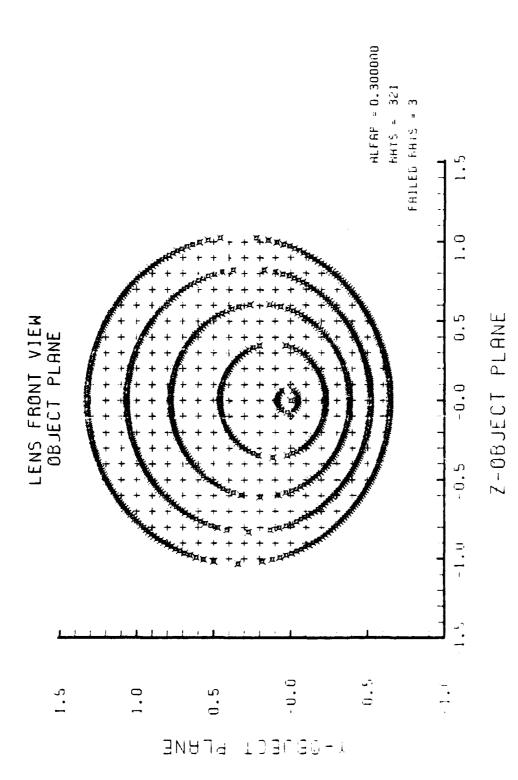


Figure F-54. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-53

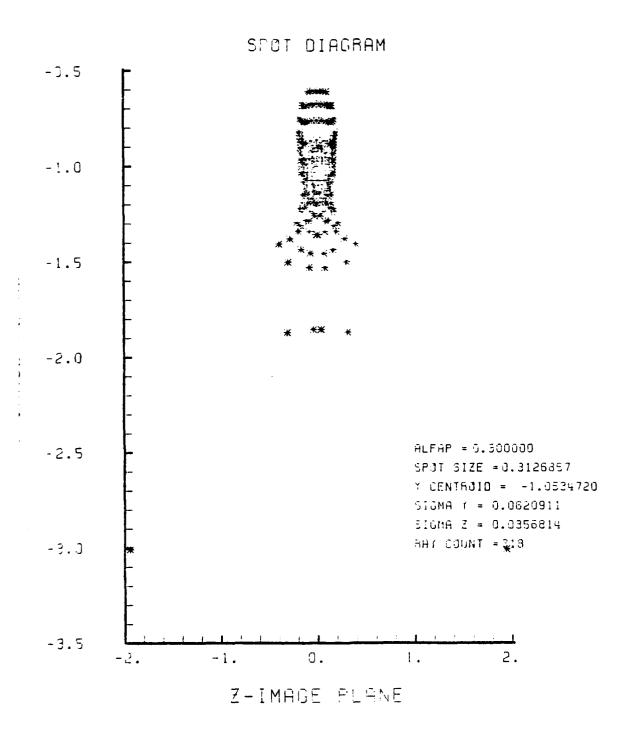


Figure F-55. Spot Diagram for Grid of Figure F-54

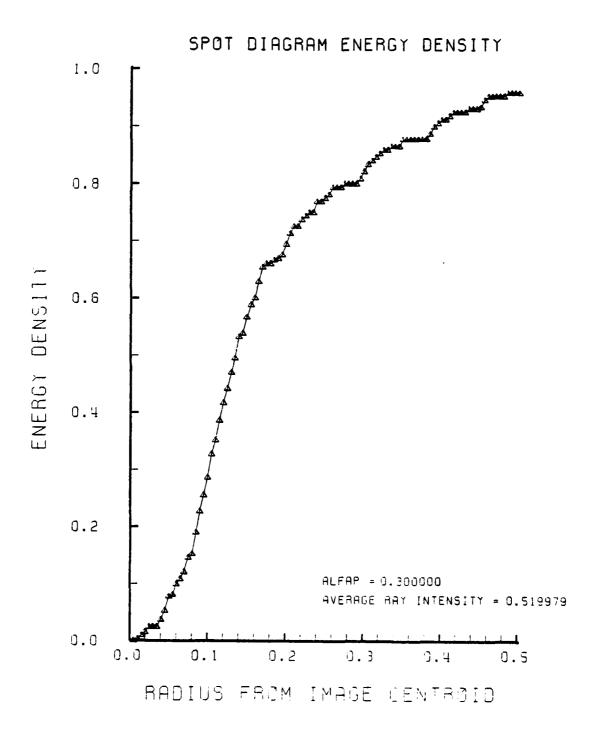


Figure F-56. Encircled Energy of Figure F-55

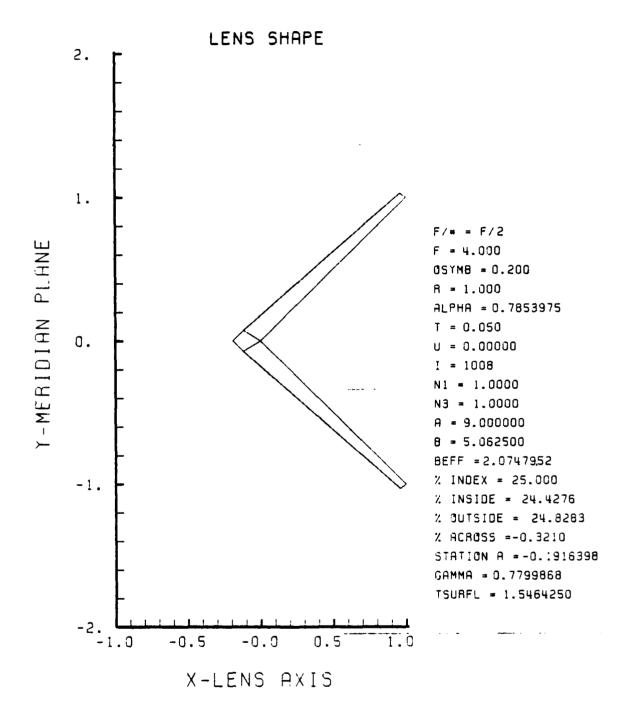
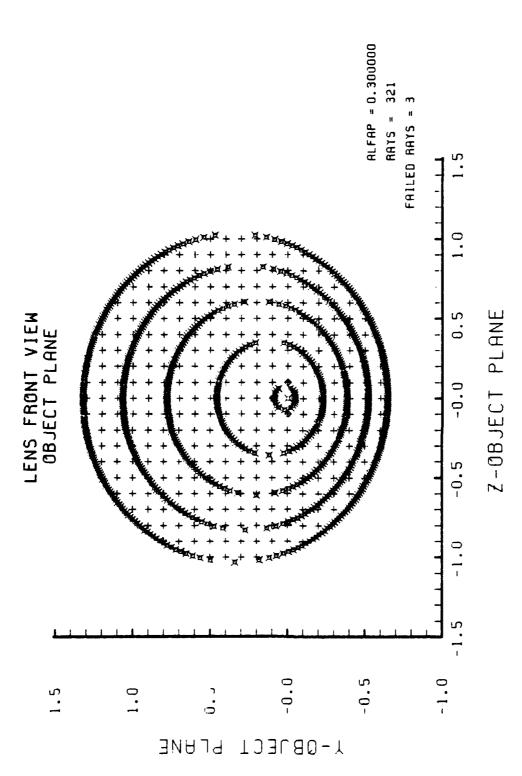


Figure F-57. GRIN Lens Shape at +25%, OB = 0.20, a = 9.00



Grid Plane at  $\alpha_p = 0.3$  for Lens of Figure F-57 Figure F-58.

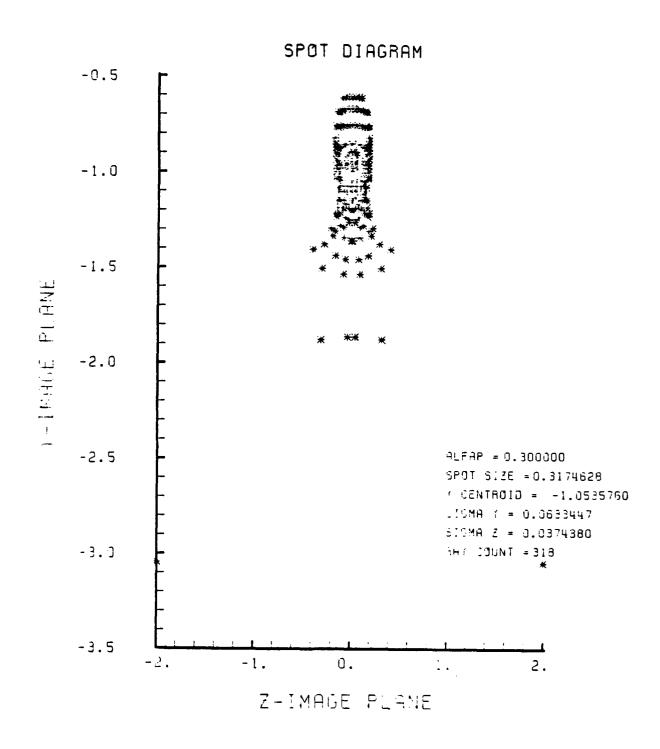


Figure F-59. Spot Diagram for Grid of Figure F-58

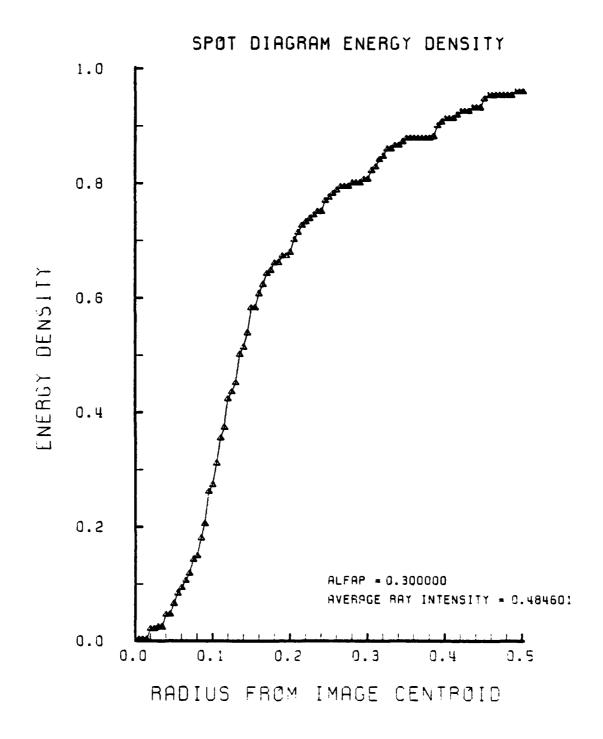


Figure F-60. Encircled Energy of Figure F-59

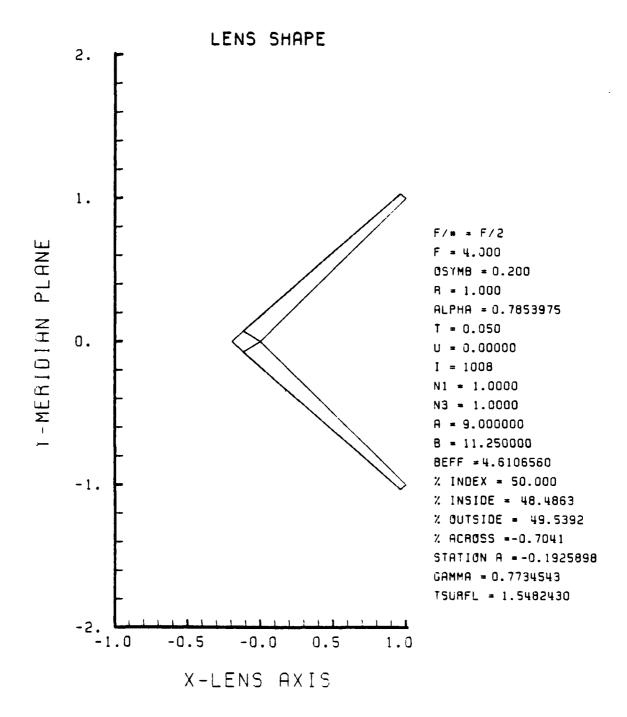
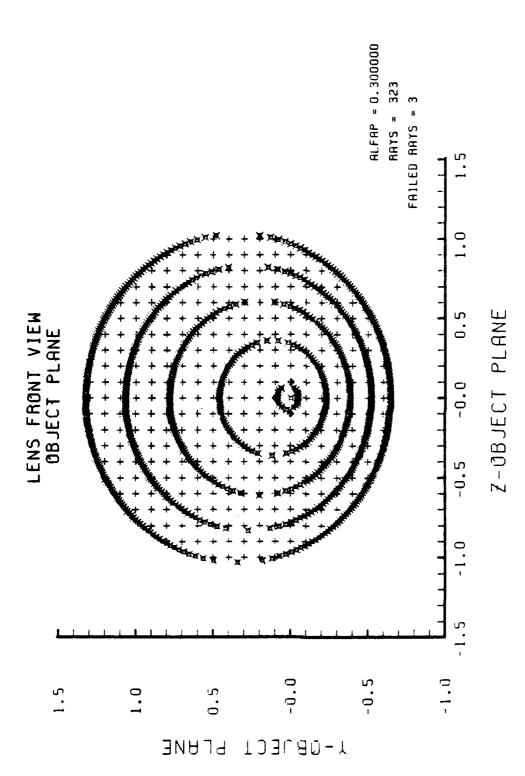


Figure F-61. GRIN Lens Shape at +50%, OB = 0.20, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-61 Figure F-62.

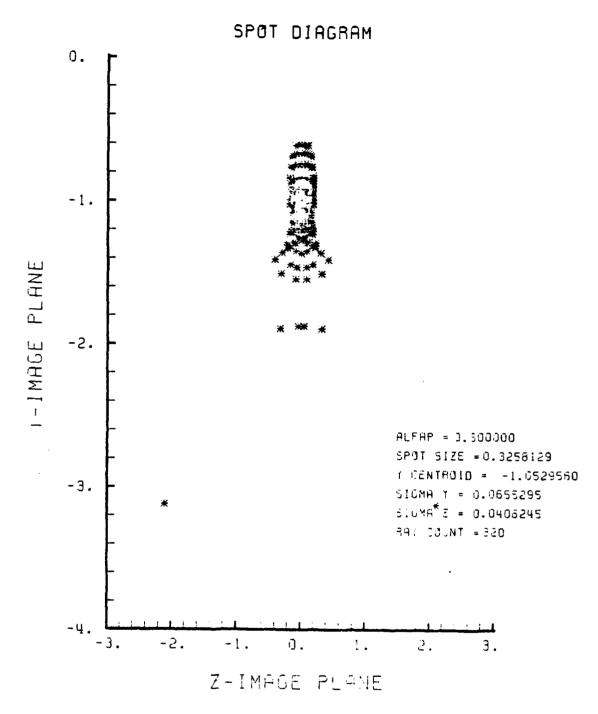


Figure F-63. Spot Diagram for Grid of Figure F-62

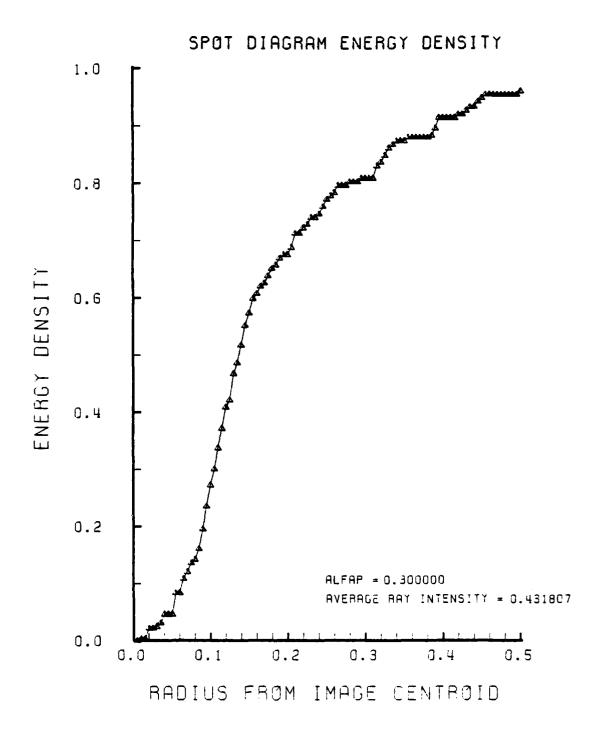


Figure F-64. Encircled Energy of Figure F-63

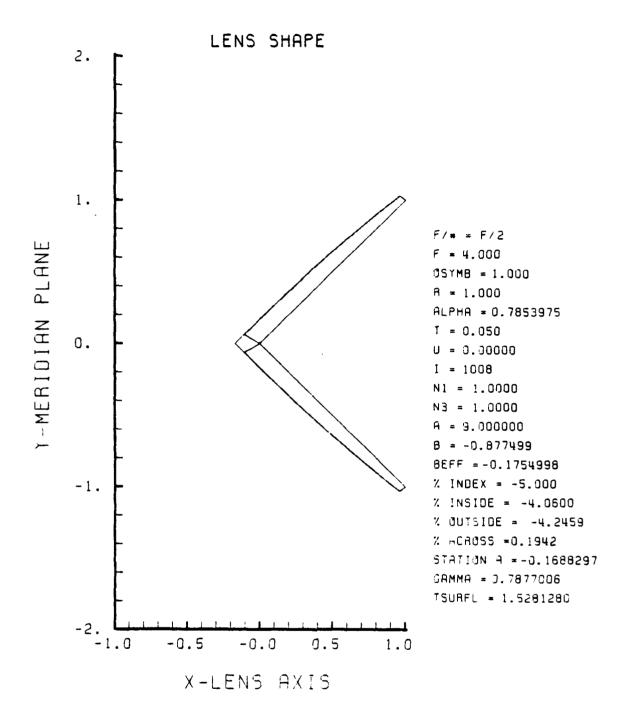


Figure F-65. GRIN Lens Shape at -5%, OB = 1.00, a = 9.00

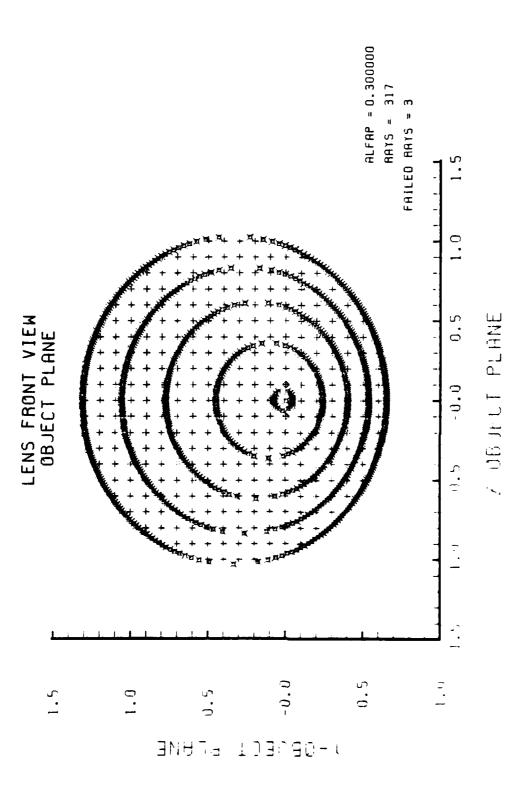


Figure 1-66. Grid Plane at  $\alpha_p=0.3$  for Lens of Figure F-65

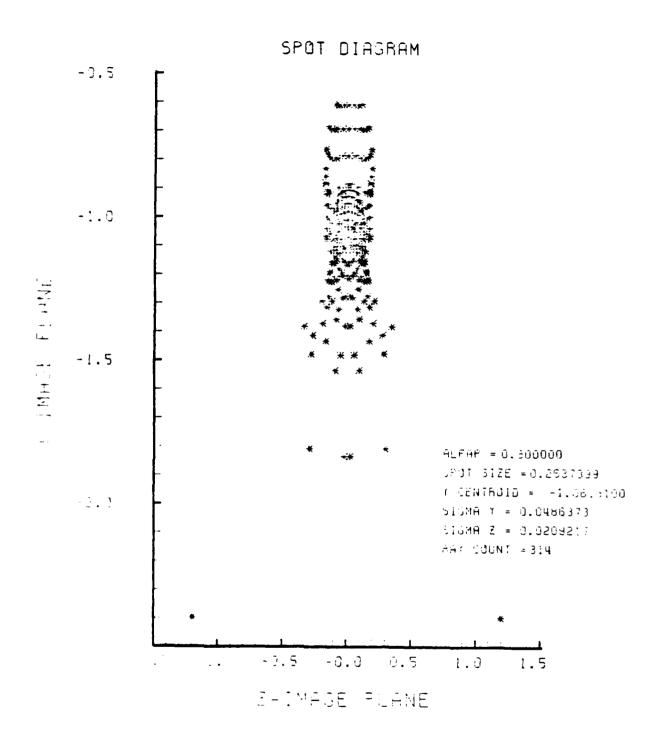


Figure F-67. Spot Diagram for Grid of Figure F-66

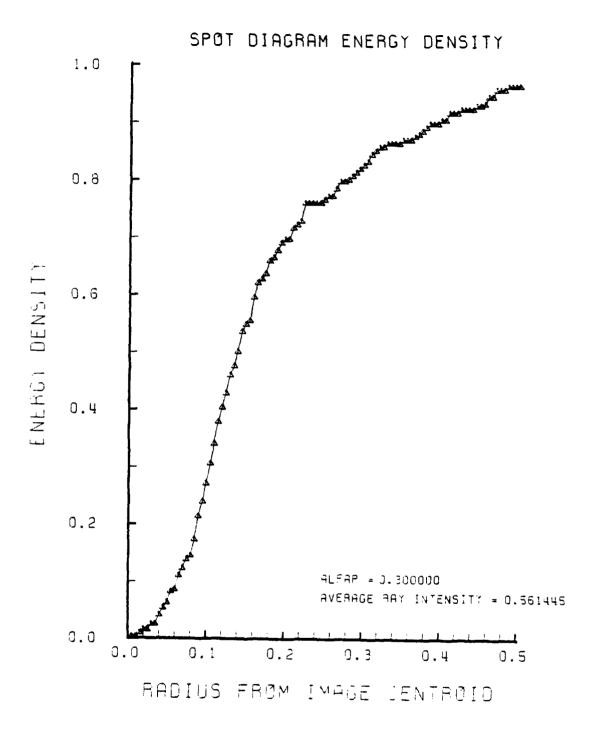


Figure F-68. Encircled Energy of Figure F-67

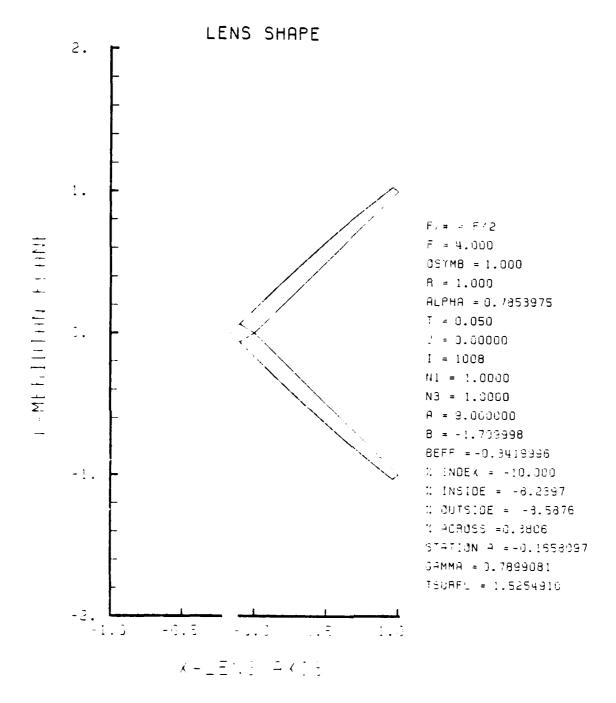
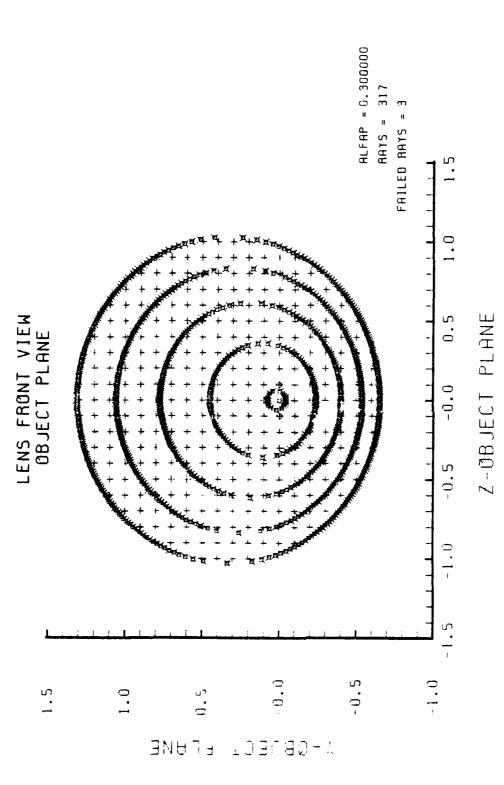


Figure F-69. GRIN Lens Shape at -10%, OB = 1.00, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-69 Figure F-70.

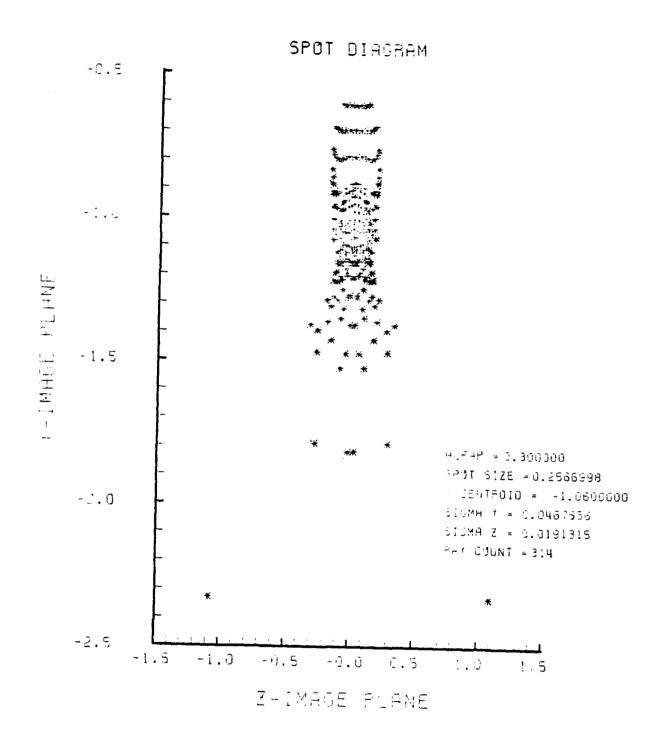


Figure F-71. Spot Diagram for Grid of Figure F-70

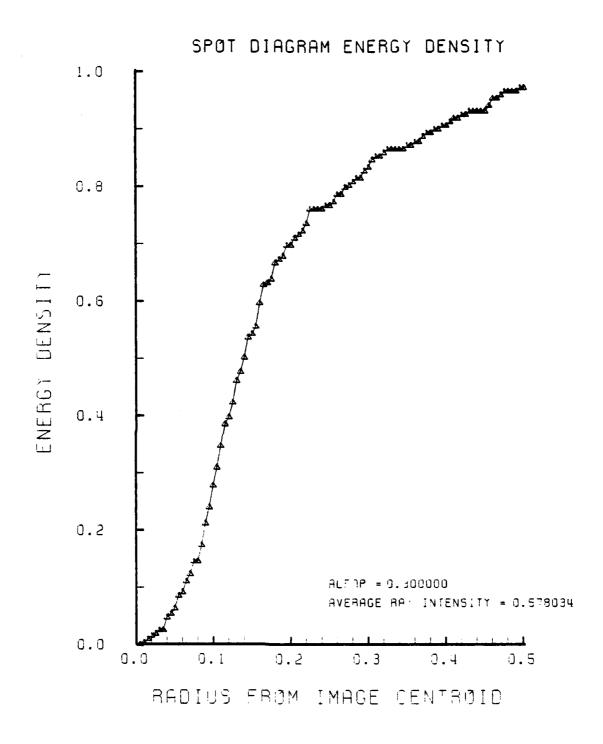


Figure F-72. Encircled Energy of Figure F-71

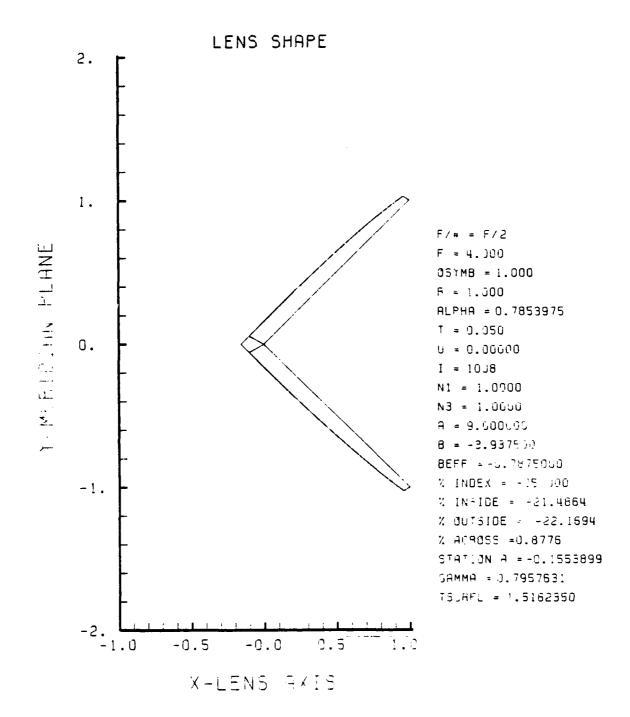
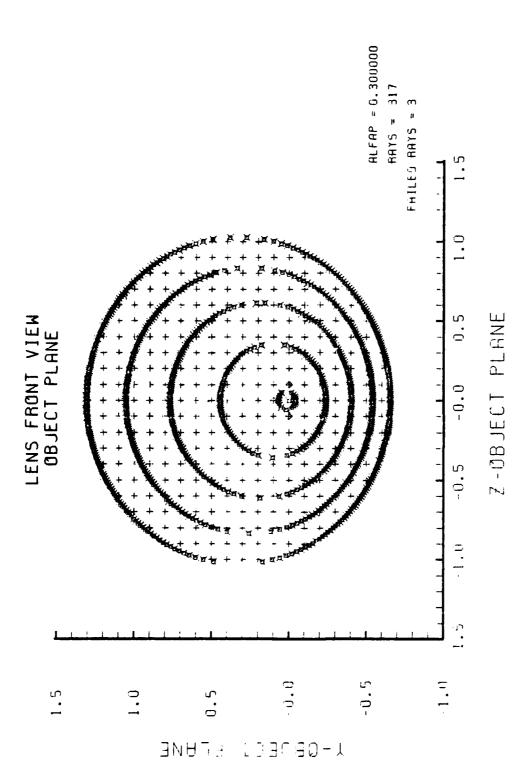


Figure F-73. GRIN Lens Shape at -25%, OB = 1.00, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-73 Figure F-74.

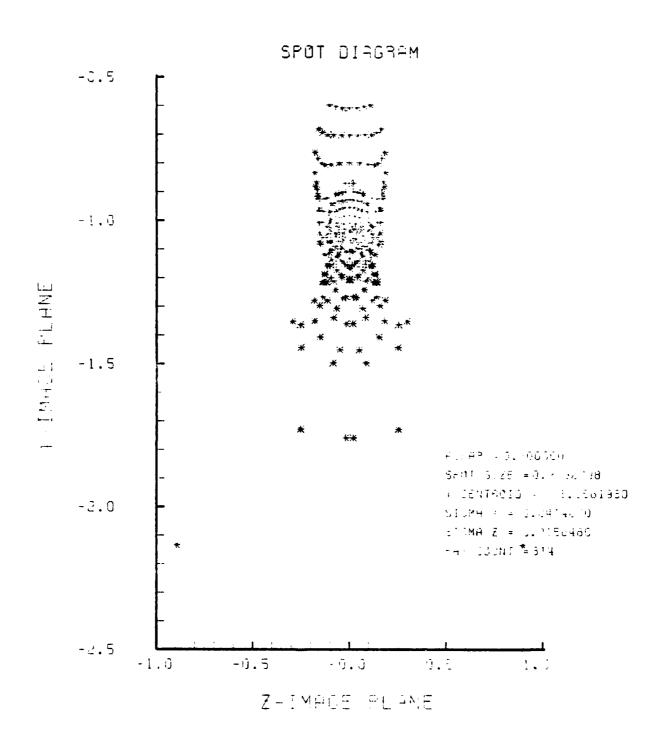


Figure F-75. Spot Diagram for Grid of Figure F-74

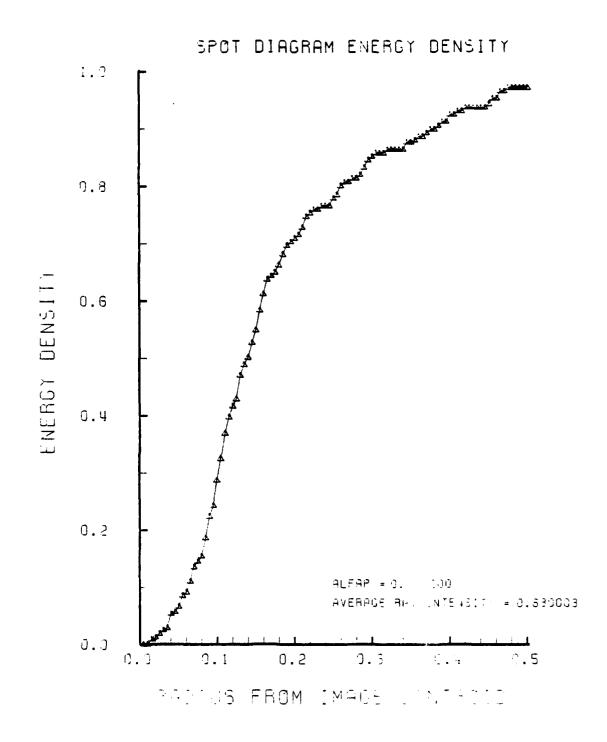


Figure F-76. Encircled Energy of Figure F-75

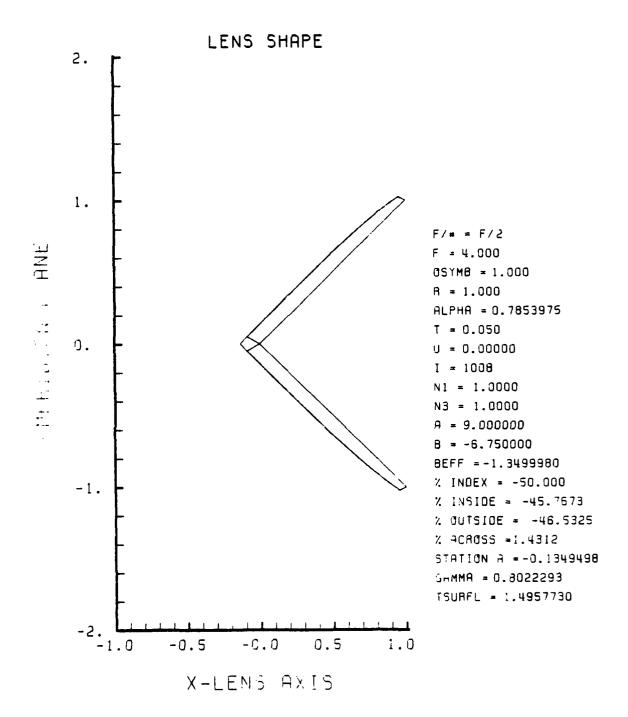
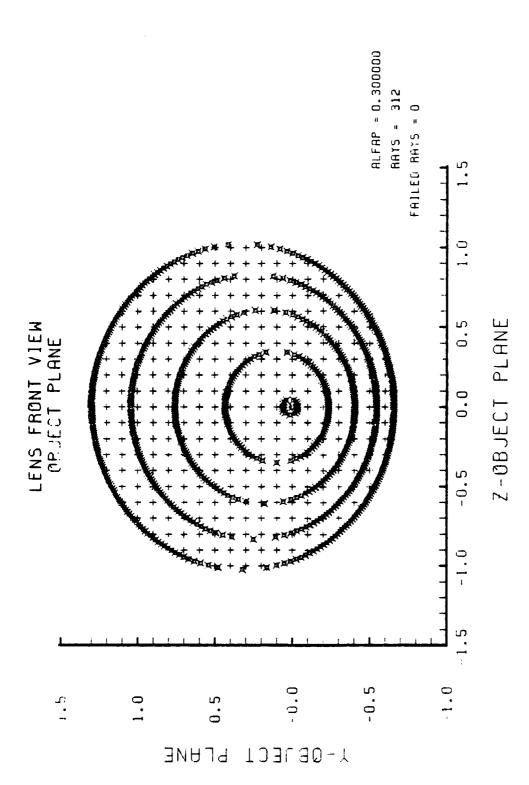


Figure F-77. GRIN Lens Shape at -50%, OB = 1.00, a = 9.00



= 0.3 for Lens of Figure F-77 Grid Plane at  $^{lpha}_{p}$ Figure F-78.

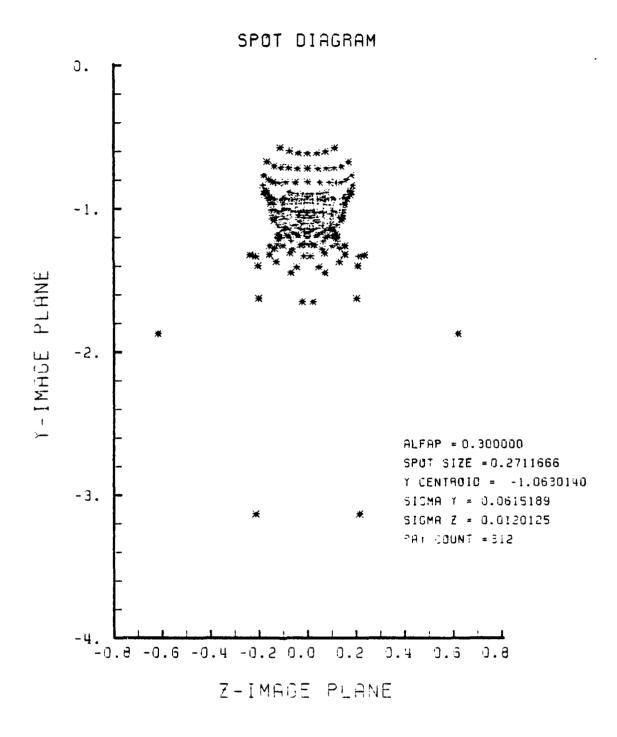


Figure F-79. Spot Diagram for Grid of Figure F-78

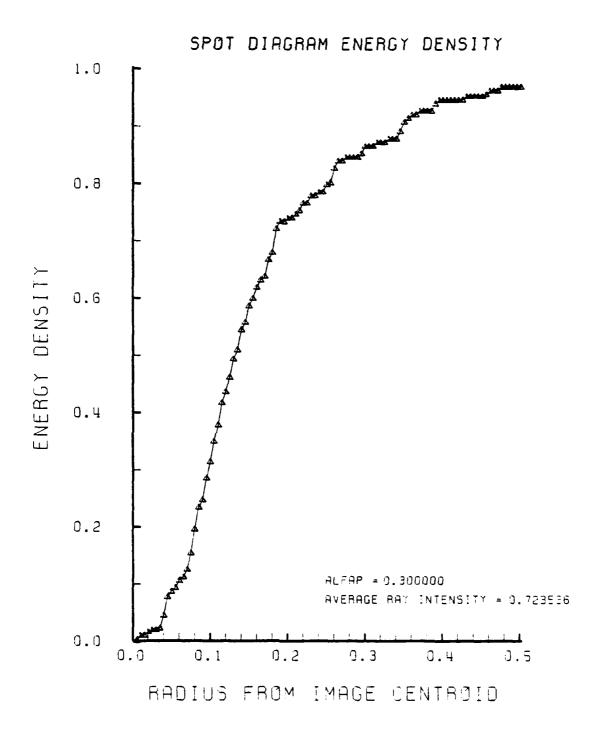
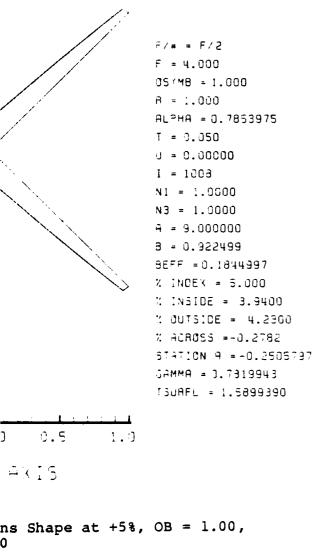
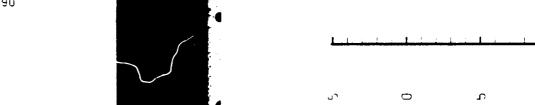


Figure F-80. Encircled Energy of Figure F-79

## SHAPE

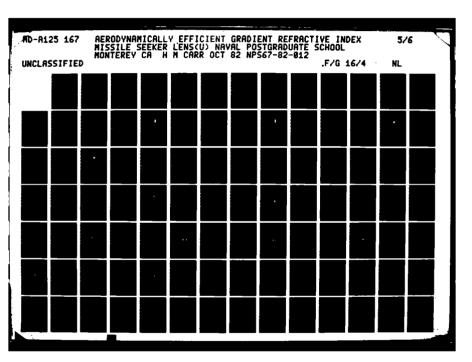


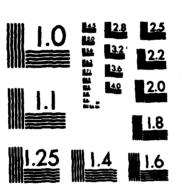


LENS FRONT VIEW OBJECT PLANE

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381





MICROCOPY RESOLUTION TEST CHART NATIONAL BUREAU OF STANDARDS-1963-A

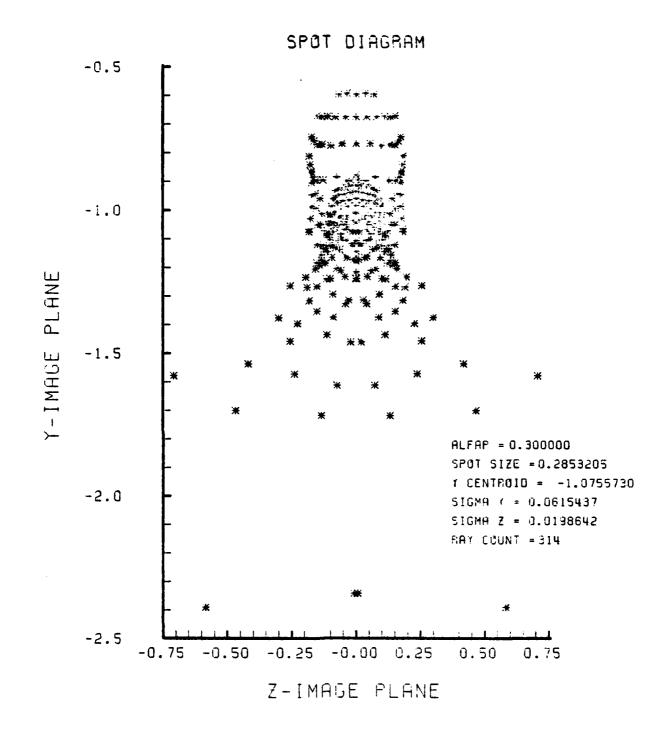


Figure F-83. Spot Diagram for Grid of Figure F-82

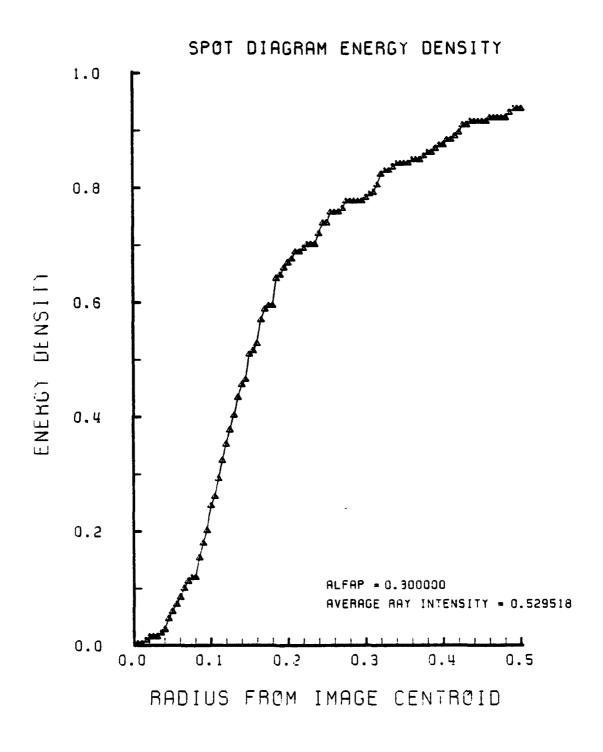
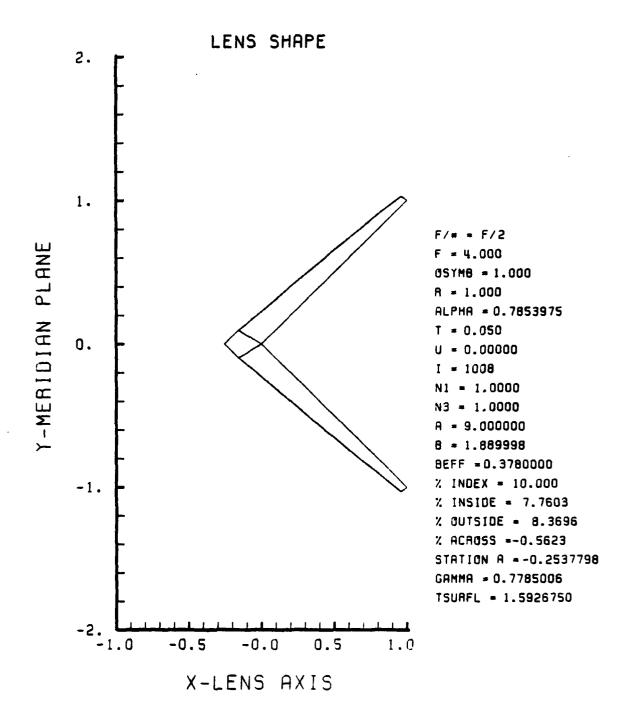


Figure F-84. Encircled Energy of Figure F-83

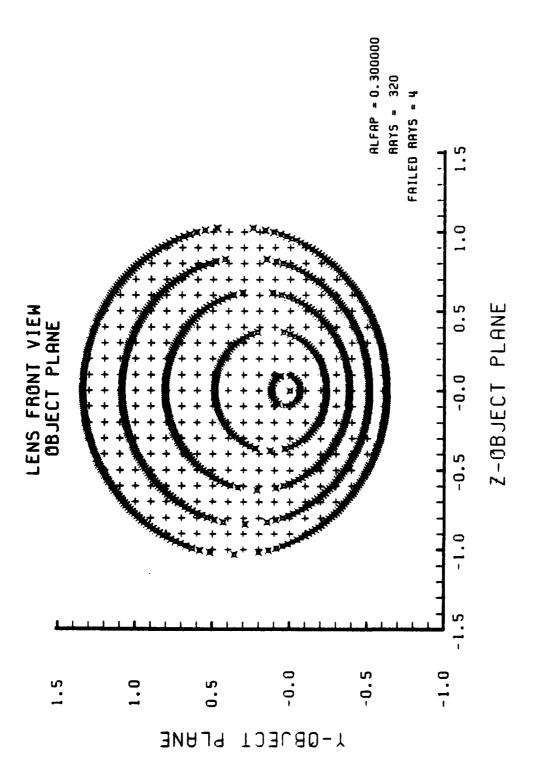


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Figure F-85. GRIN Lens Shape at +10%, OB = 1.00, a = 9.00



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Grid Plane at  $\alpha$  = 0.3 for Lens of Figure F-85 Figure F-86.

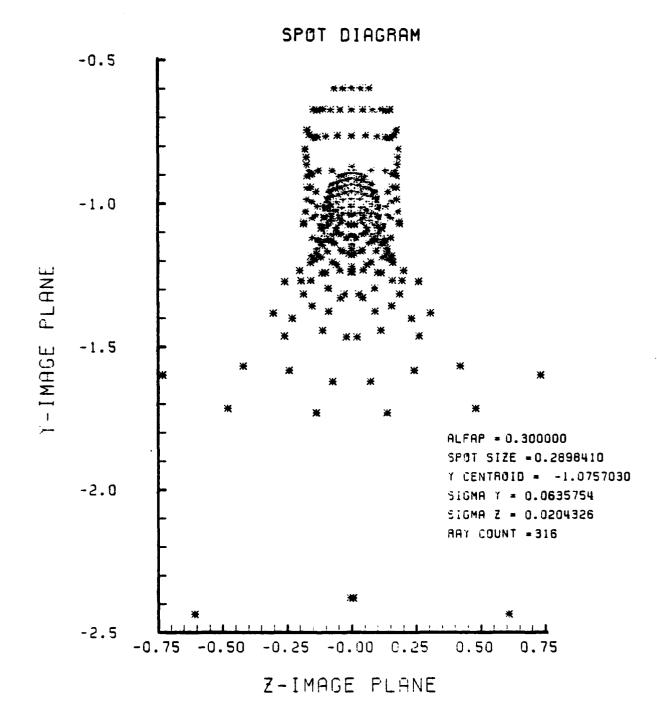


Figure F-87. Spot Diagram for Grid of Figure F-86

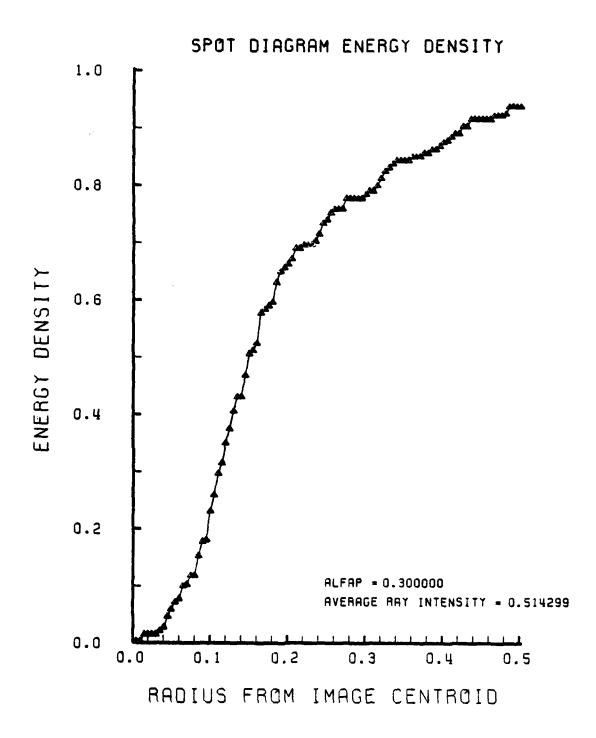


Figure F-88. Encircled Energy of Figure F-87

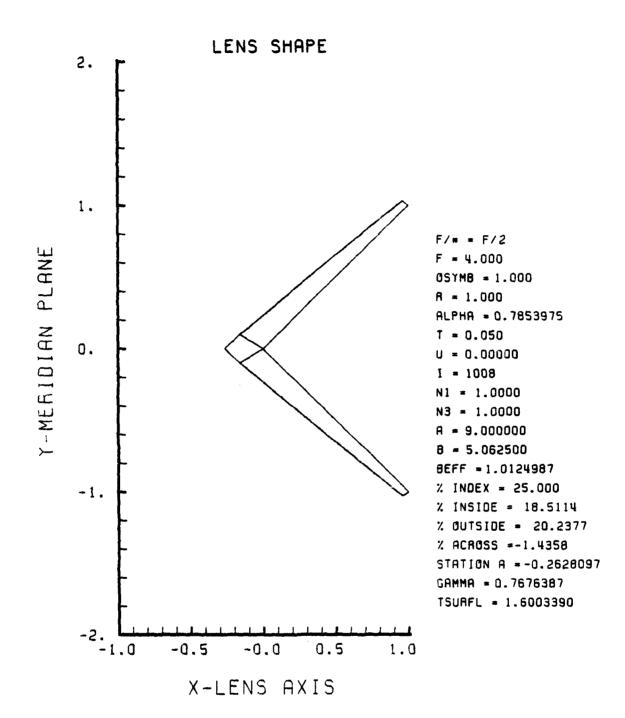
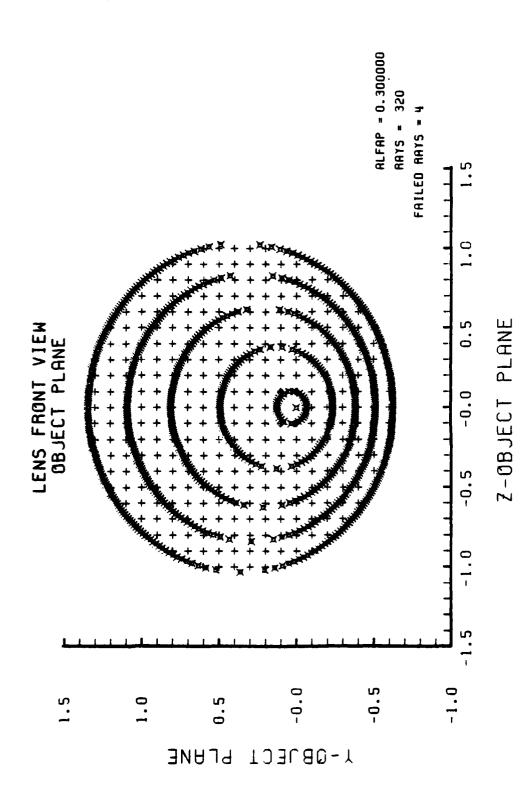


Figure F-89. GRIN Lens Shape at +25%, OB = 1.00, a = 9.00



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Figure F-90. Grid Plane at  $\alpha_{\rm p}=$  0.3 for Lens of Figure F-89

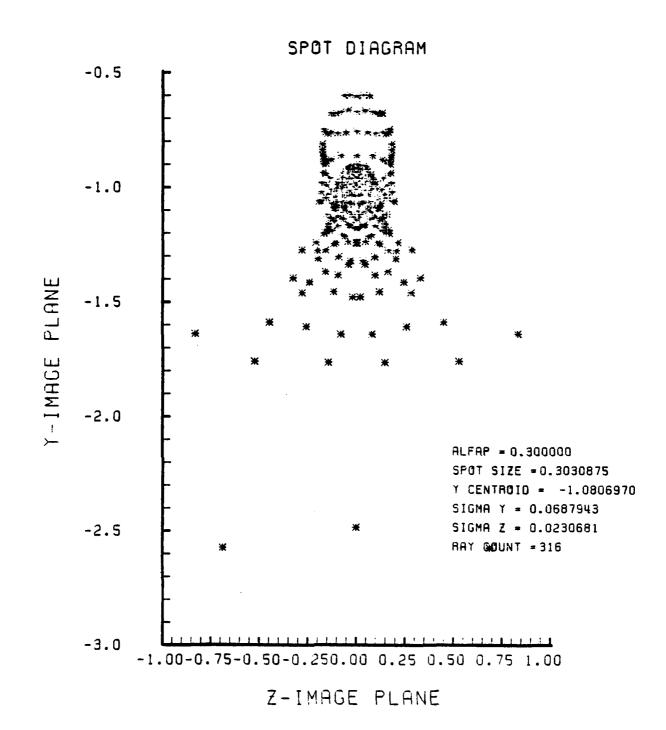


Figure F-91. Spot Diagram for Grid of Figure F-90

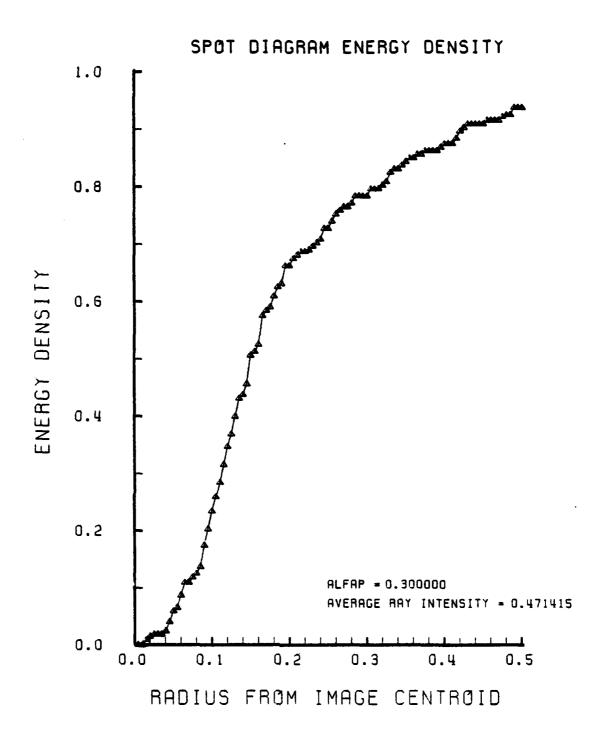


Figure F-92. Encircled Energy of Figure F-91

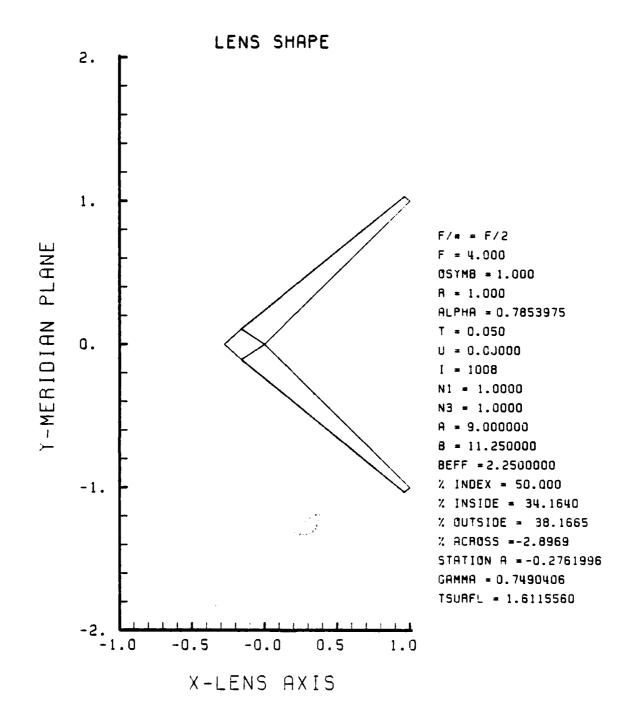
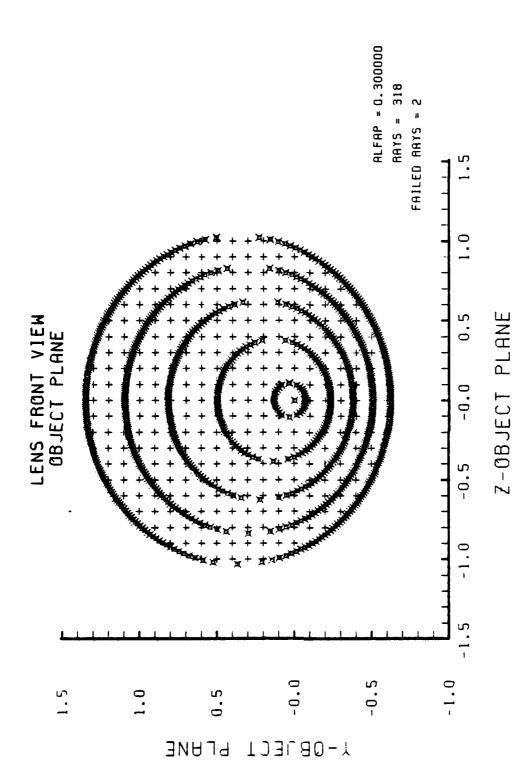


Figure F-93. GRIN Lens Shape at +50%, OB = 1.00, a = 9.00



Grid Plane at  $\alpha_p = 0.3$  for Lens of Figure F-93 Figure F-94.

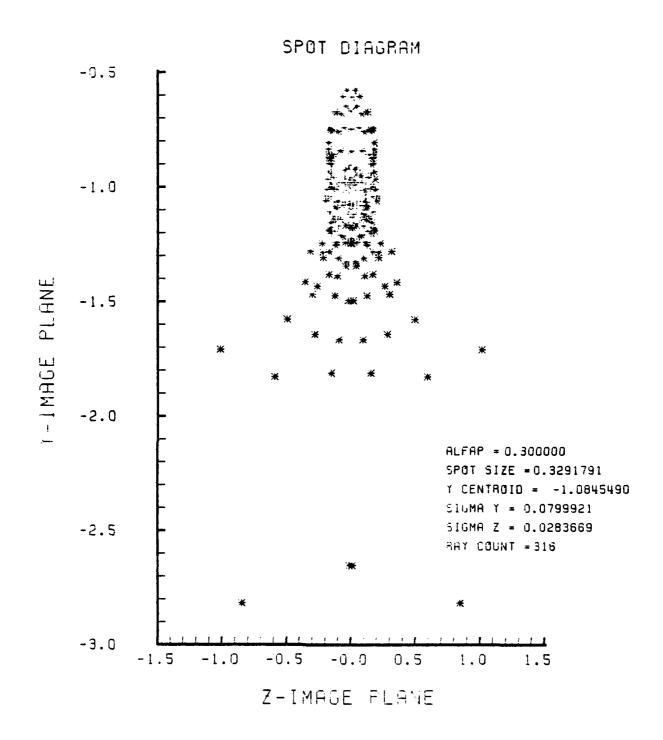


Figure F-95. Spot Diagram for Grid of Figure F-94

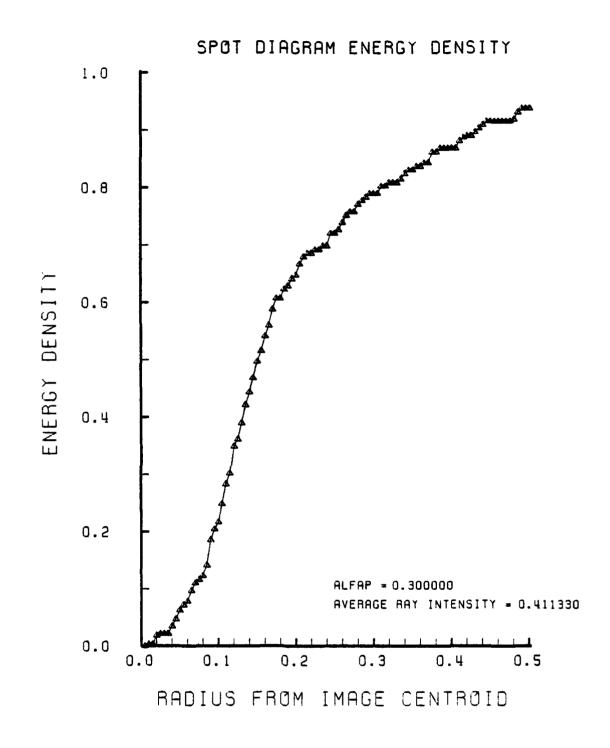


Figure F-96. Encircled Energy of Figure F-95

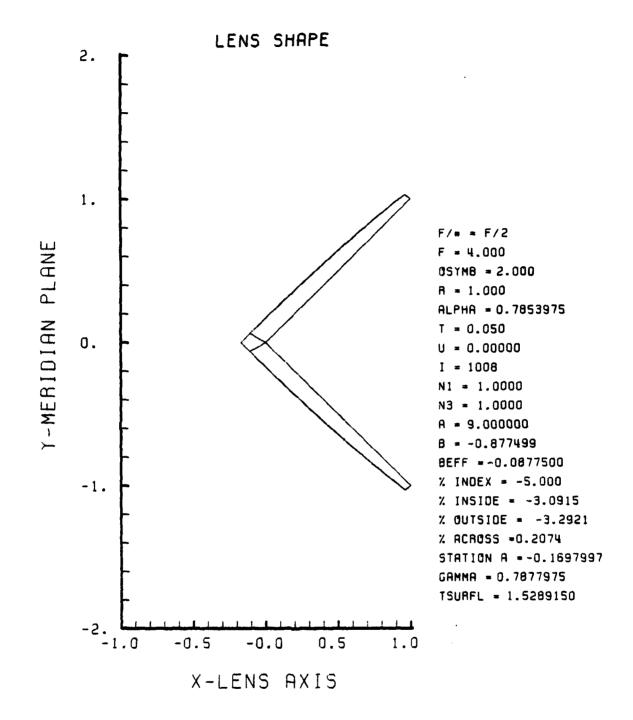


Figure F-97. GRIN Lens Shape at -5%, OB = 2.00, a = 9.00

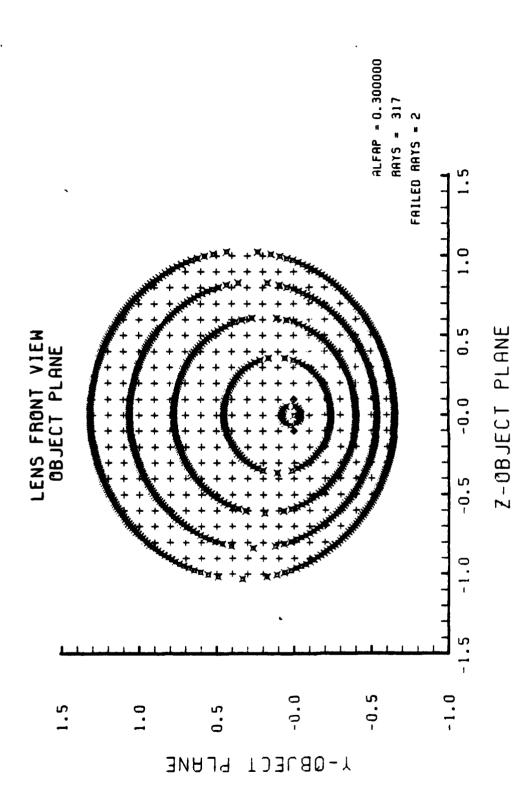


Figure F-98. Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure F-97

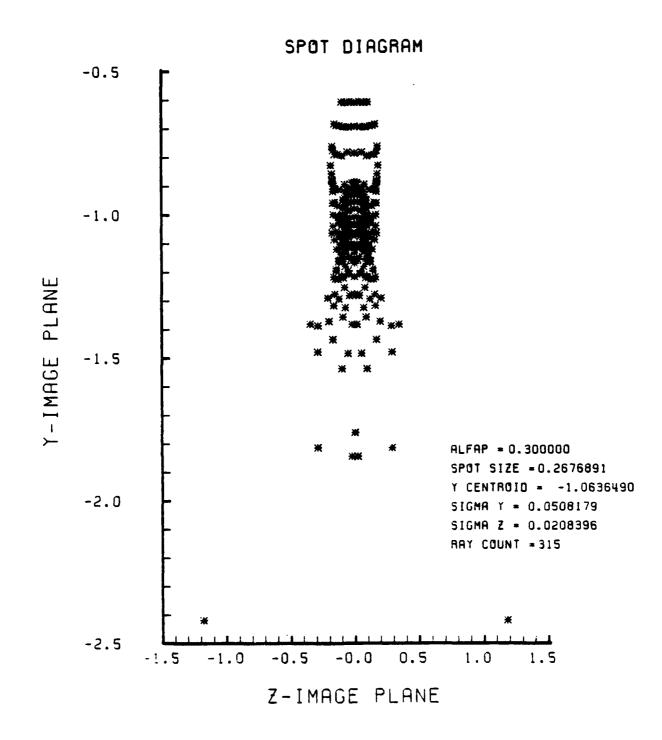


Figure F-99. Spot Diagram for Grid of Figure F-98

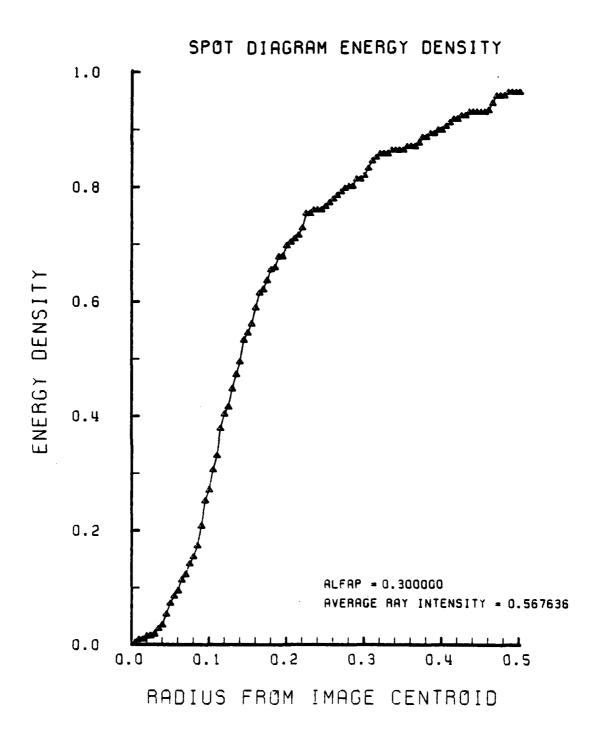


Figure F-100. Encircled Energy of Figure F-99

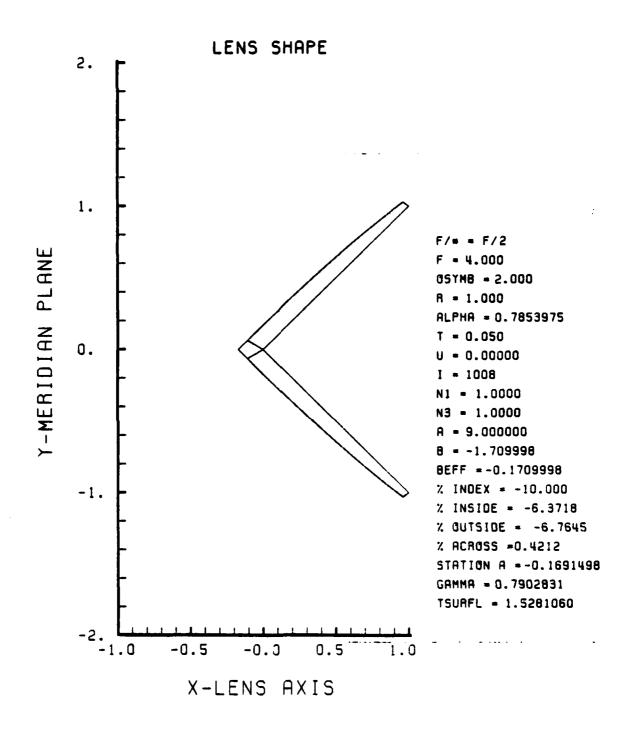
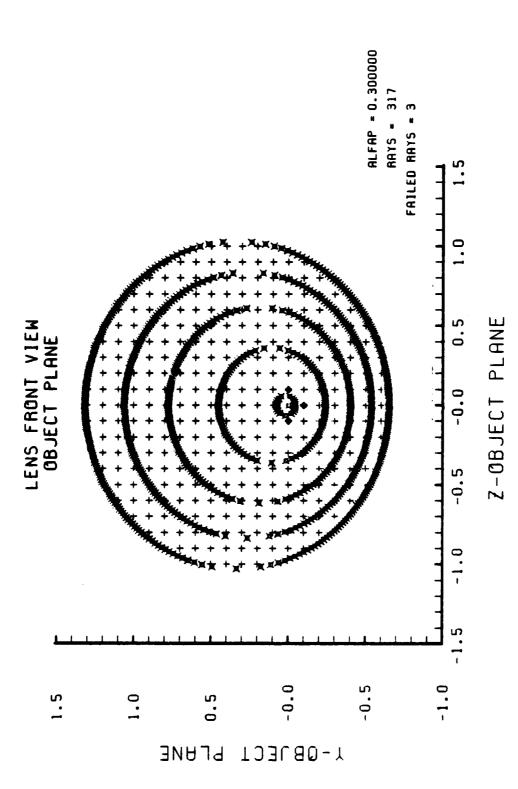


Figure F-101. GRIN Lens Shape at -10%, OB = 2.00, a = 9.00



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Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-101 Figure F-102.

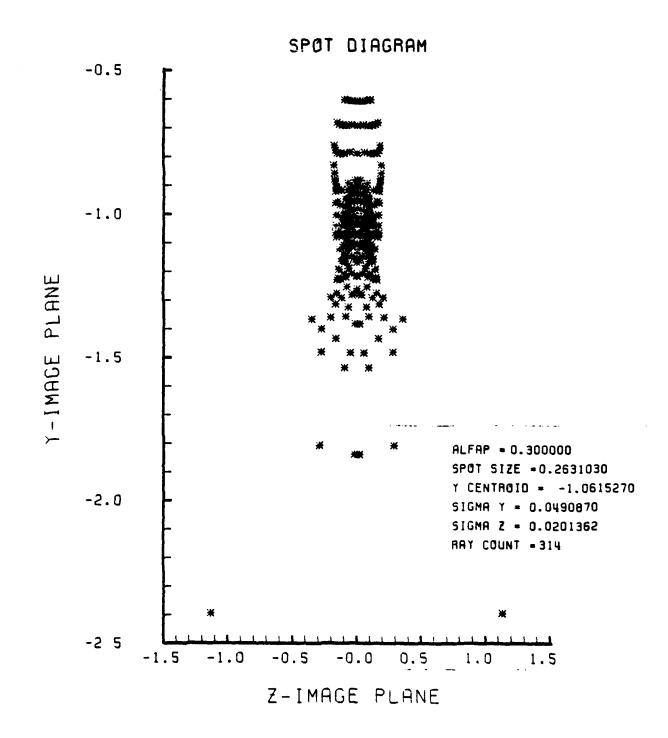


Figure F-103. Spot Diagram for Grid of Figure F-102

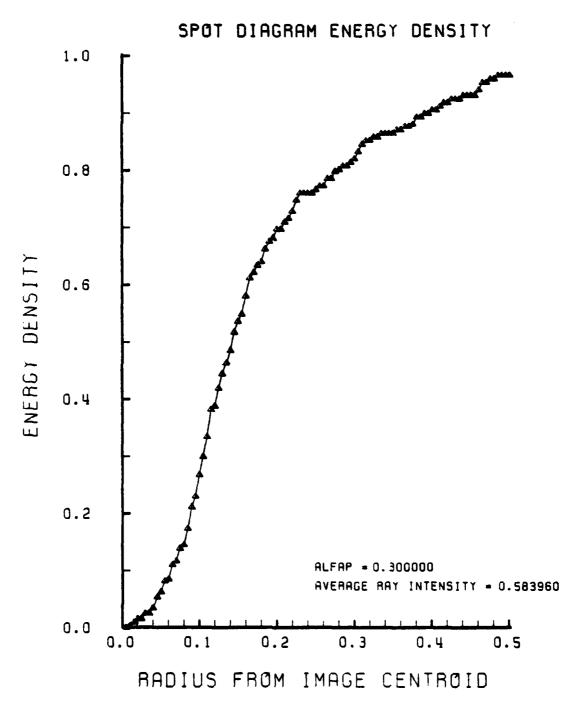


Figure F-104. Encircled Energy of Figure F-103

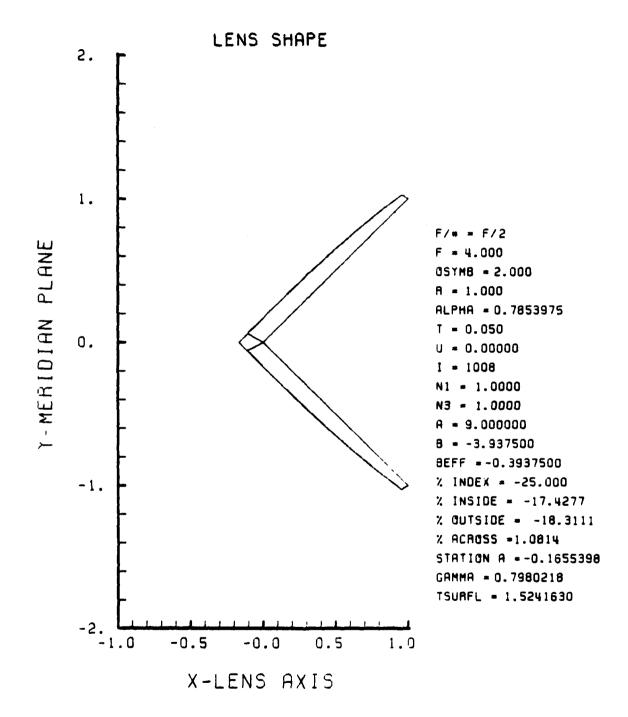
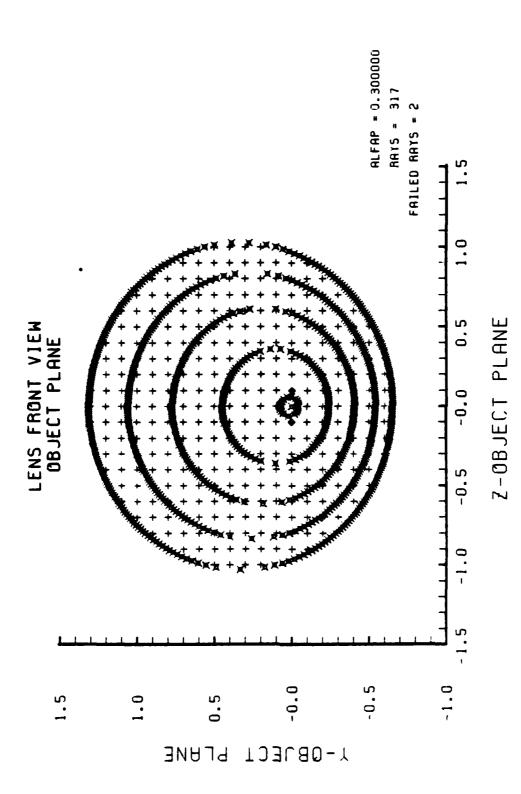


Figure F-105. GRIN Lens Shape at -25%, OB = 2.00, a = 9.00

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Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-105 Figure F-106.

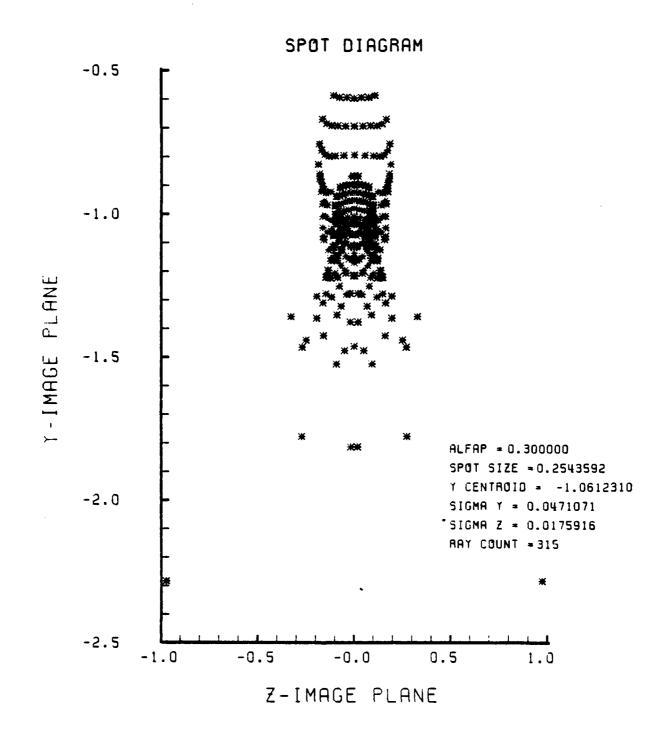


Figure F-107. Spot Diagram for Grid of Figure F-106

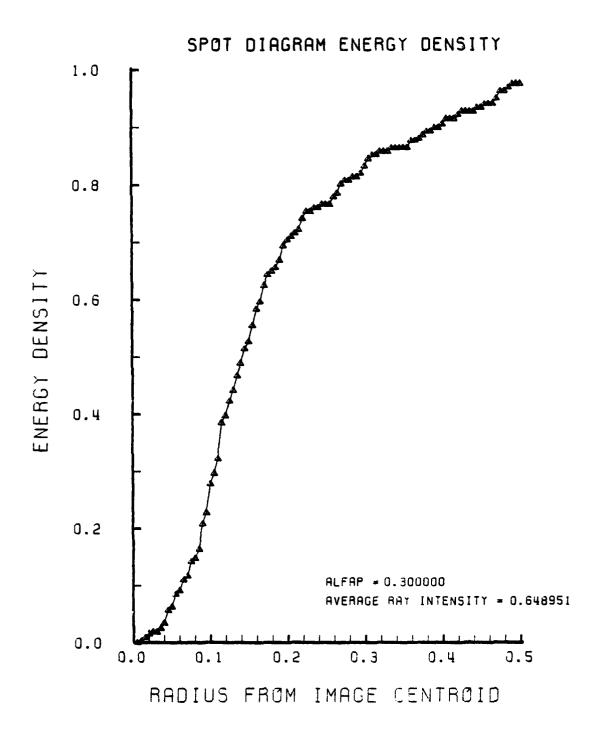


Figure F-108. Encircled Energy of Figure F-107

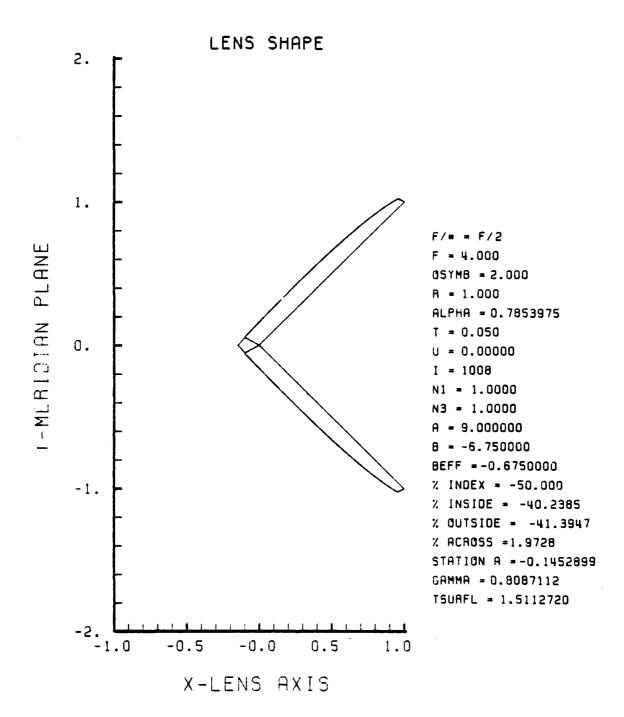
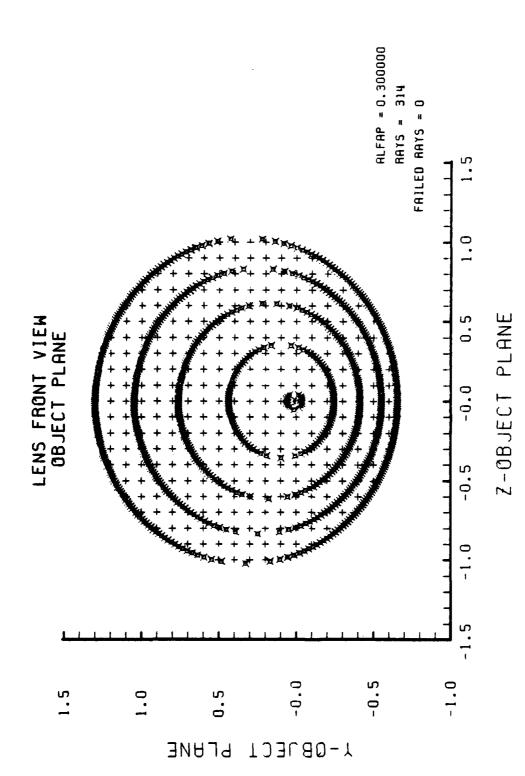


Figure F-109. GRIN Lens Shape at -50%, OB = 2.00, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Grid of Figure F-109 Figure F-110.

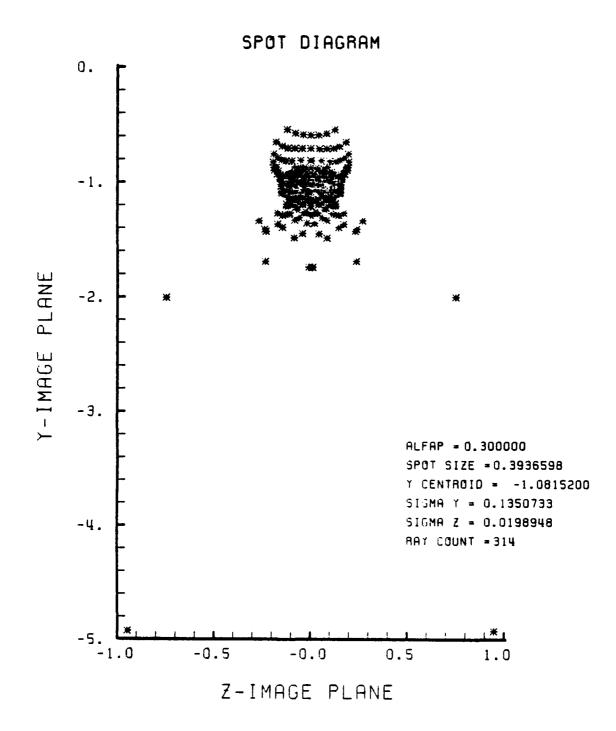


Figure F-111. Spot Diagram for Grid of Figure F-110

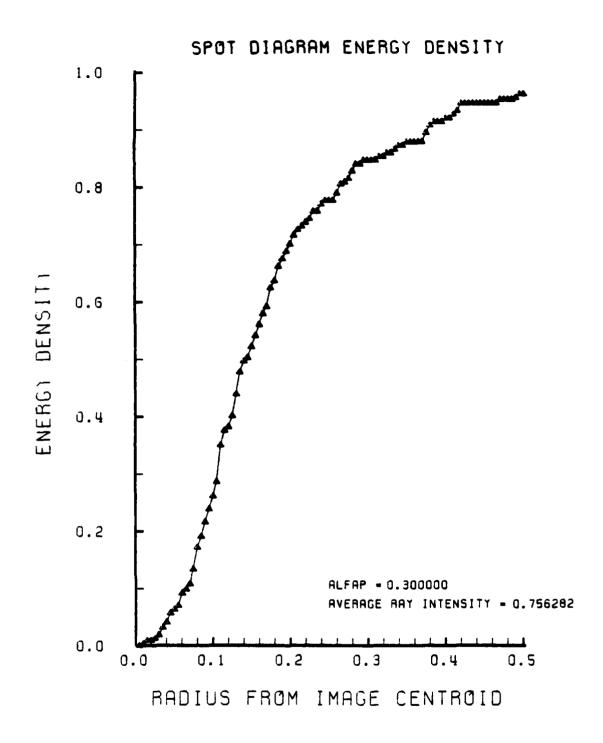


Figure F-112. Encircled Energy of Figure F-111

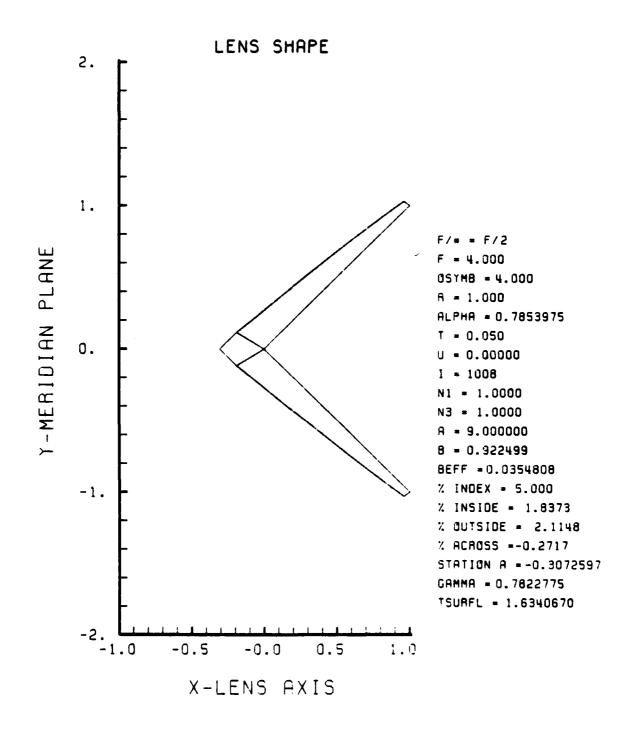


Figure F-113. GRIN Lens Shape at +5%, OB = 4.00, a = 9.00

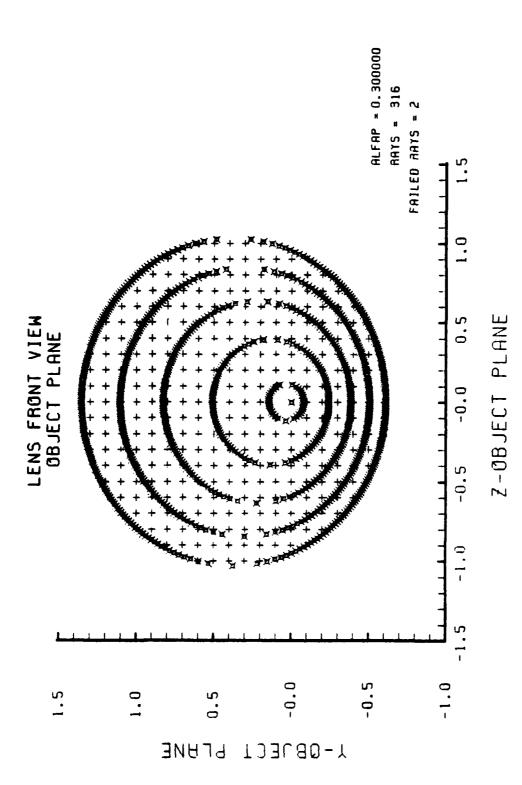


Figure F-114. Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure F-113

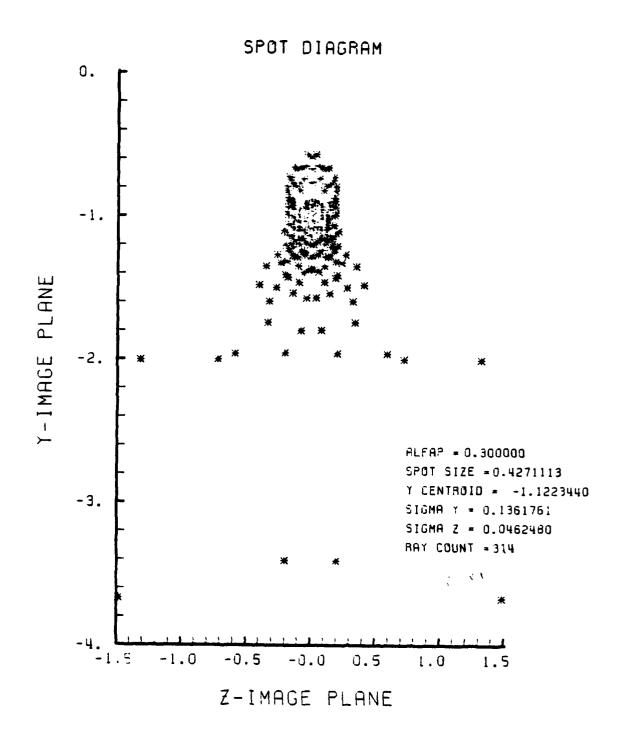


Figure F-115. Spot Diagram for Grid of Figure F-114

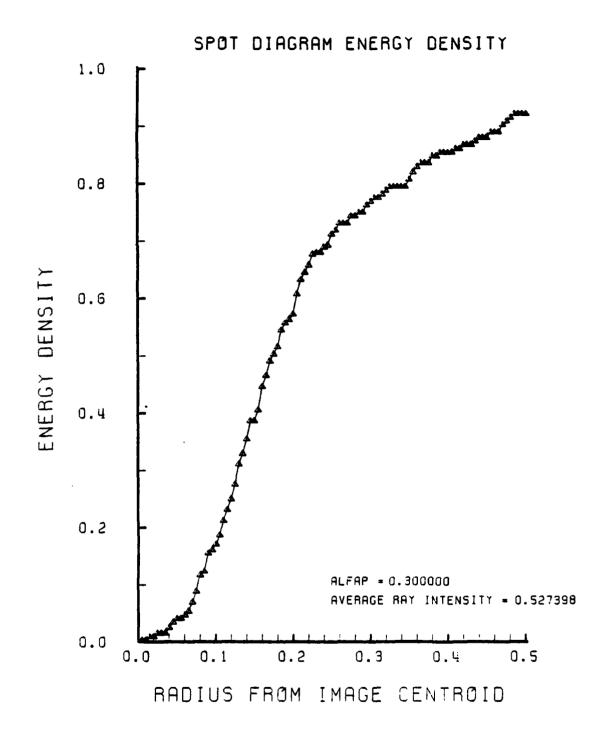


Figure F-116. Encircled Energy of Figure F-115

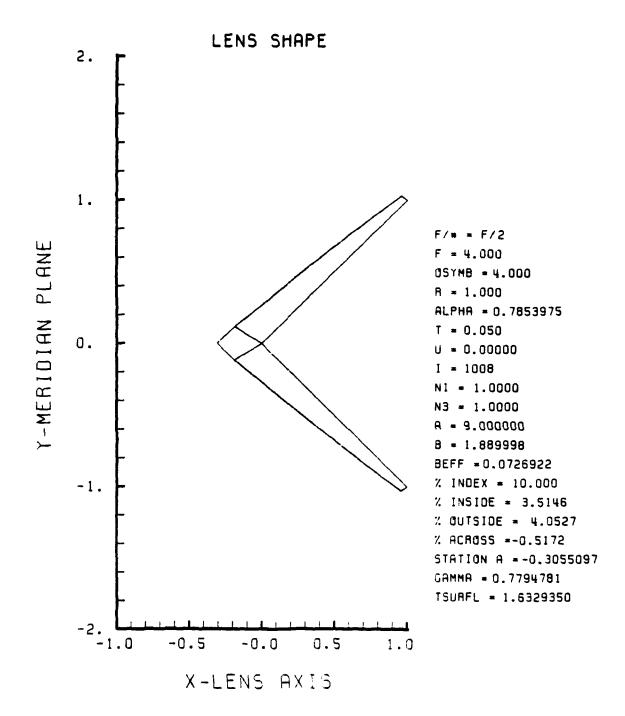
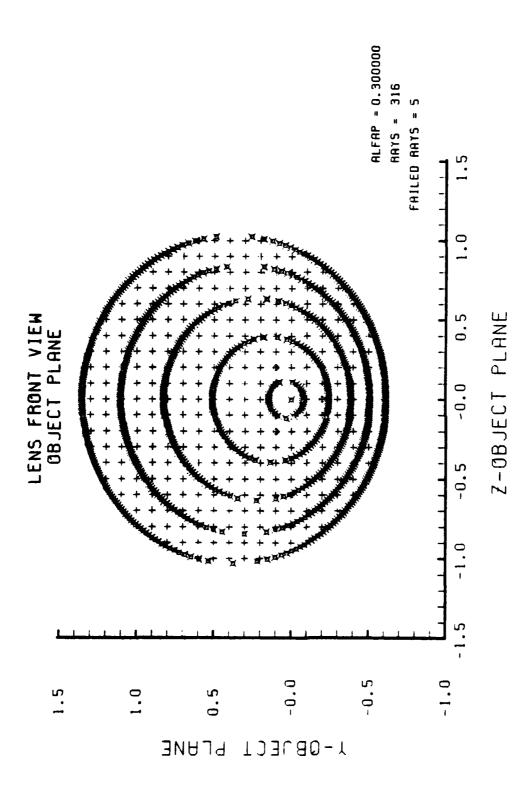


Figure F-117. GRIN Lens Shape at +10%, OB = 4.00, a = 9.00



Grid Plane at  $\alpha$  = 0.3 for Lens of Figure F-ll7 Figure F-118.

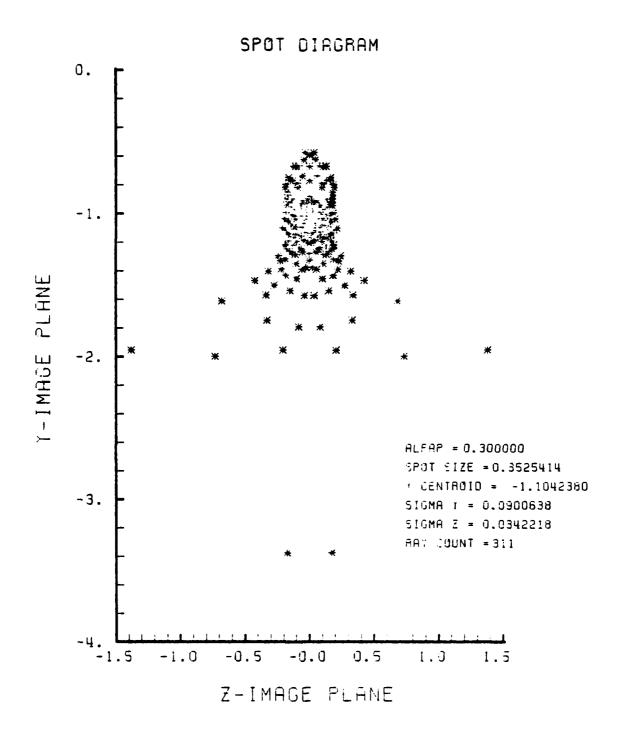


Figure F-119. Spot Diagram for Grid of Figure F-118

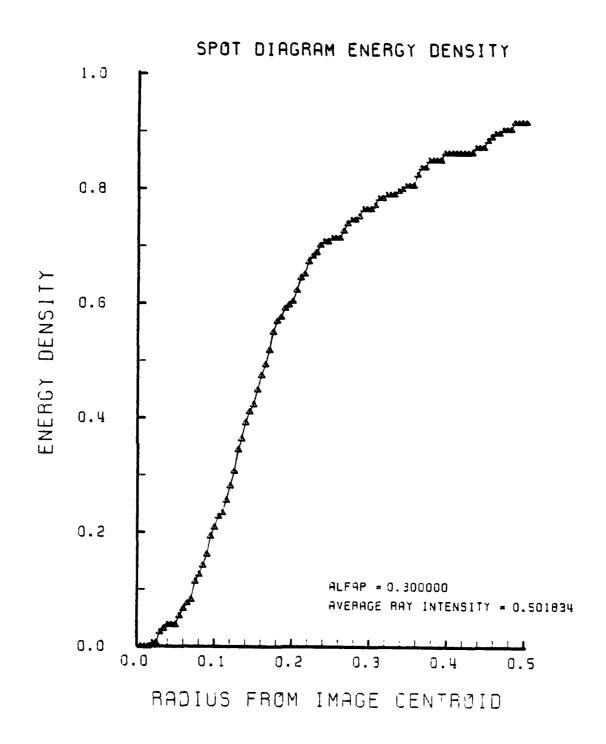


Figure F-120. Encircled Energy of Figure F-119

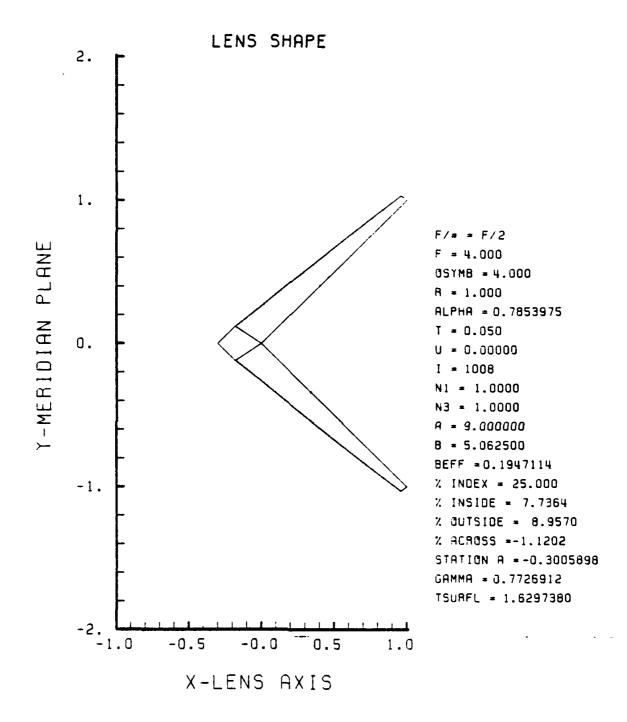
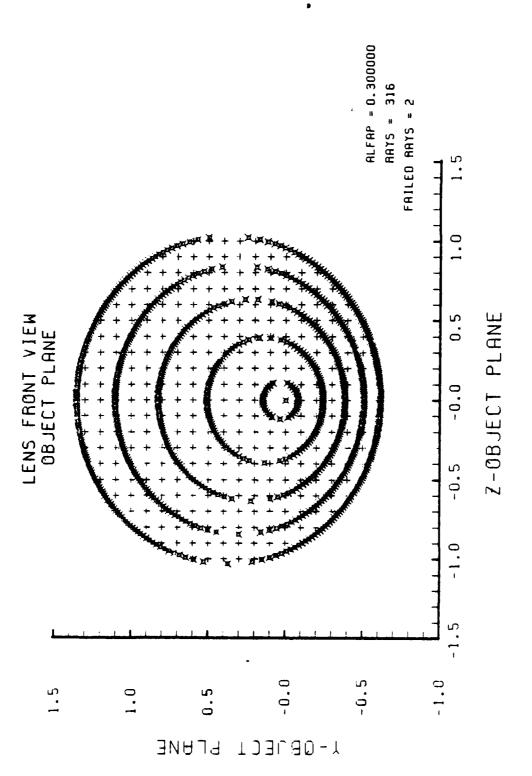


Figure F-121. GRIN Lens Shape at +25%, OB = 4.00, a = 9.00



Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-121 Figure F-122.

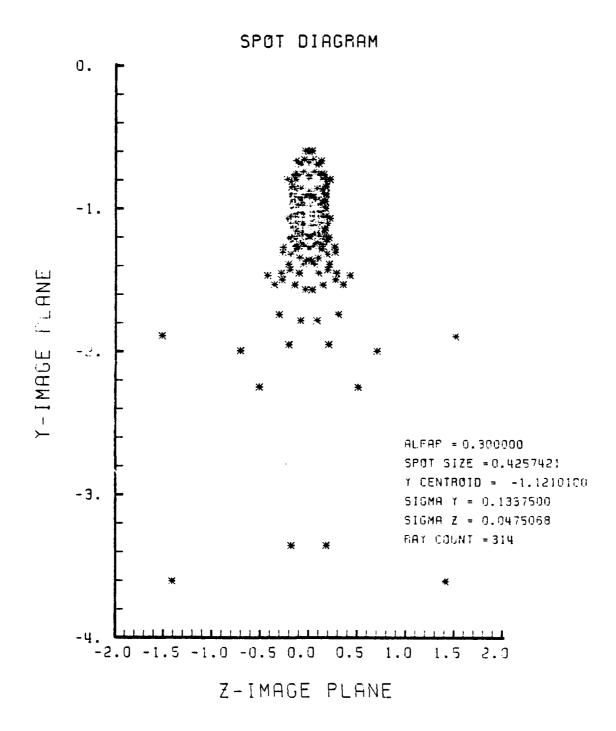


Figure F-123. Spot Diagram for Grid of Figure F-122

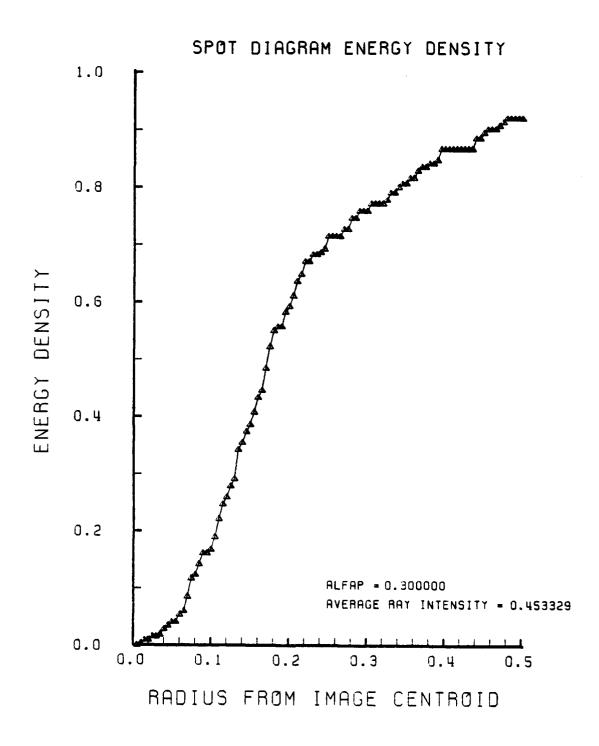


Figure F-124. Encircled Energy of Figure F-123

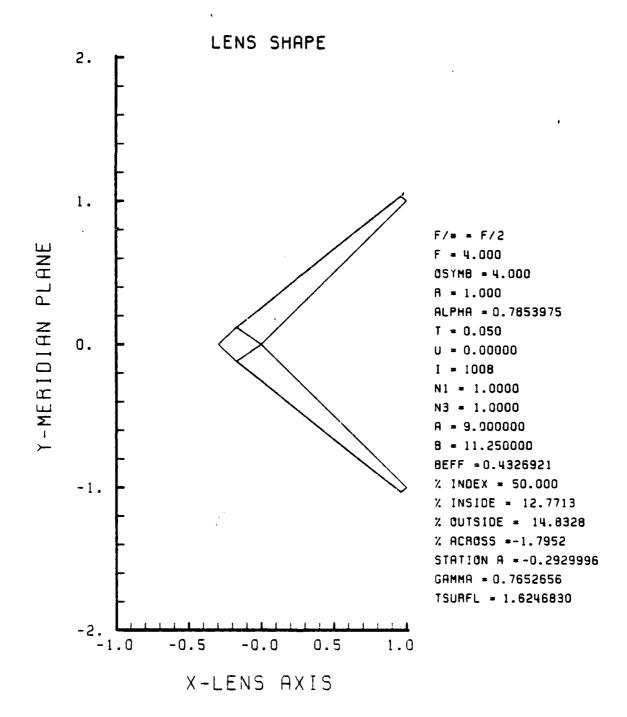


Figure F-125. GRIN Lens Shape at +50%, OB = 4.00, a = 9.00

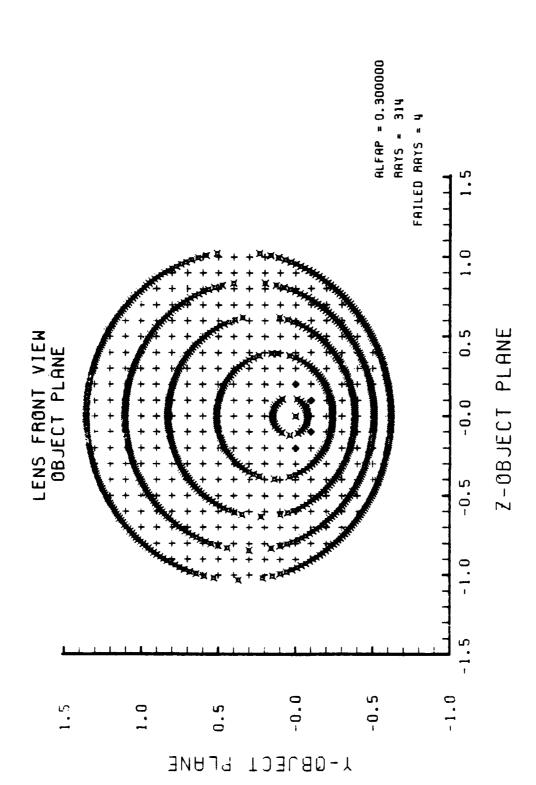


Figure F-126. Grid Plane at  $\alpha_p$  = 0.3 for Lens of Figure F-125

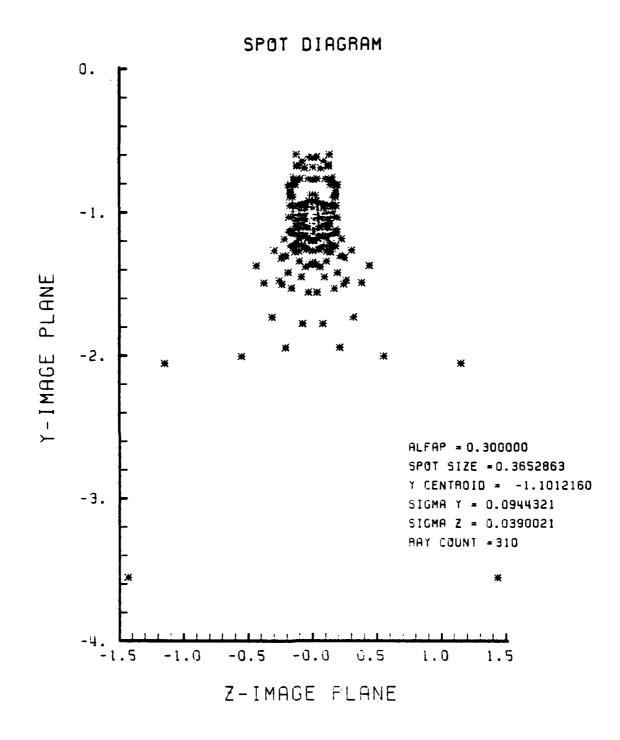


Figure F-127. Spot Diagram for Grid of Figure F-126

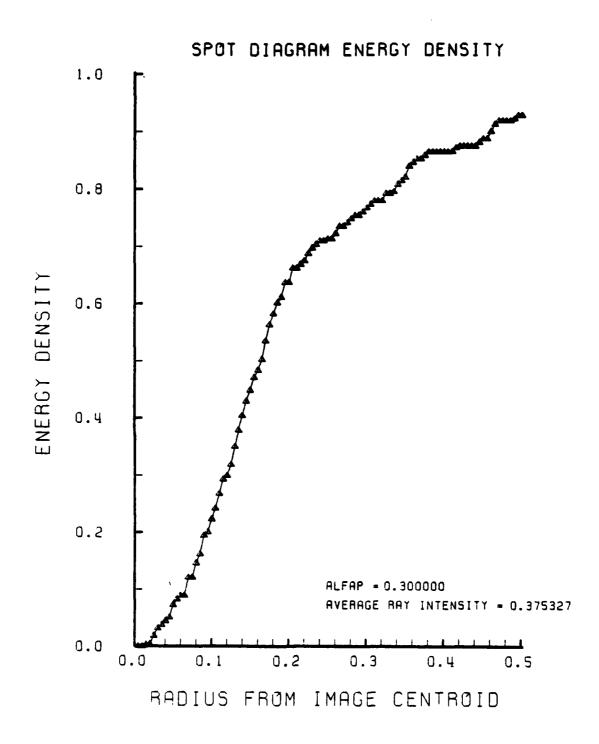


Figure F-128. Encircled Energy of Figure F-127

APPENDIX G
"BEST" GRIN LENS PERFORMANCE PLOTS IN THE F/2 CONFIGURATION

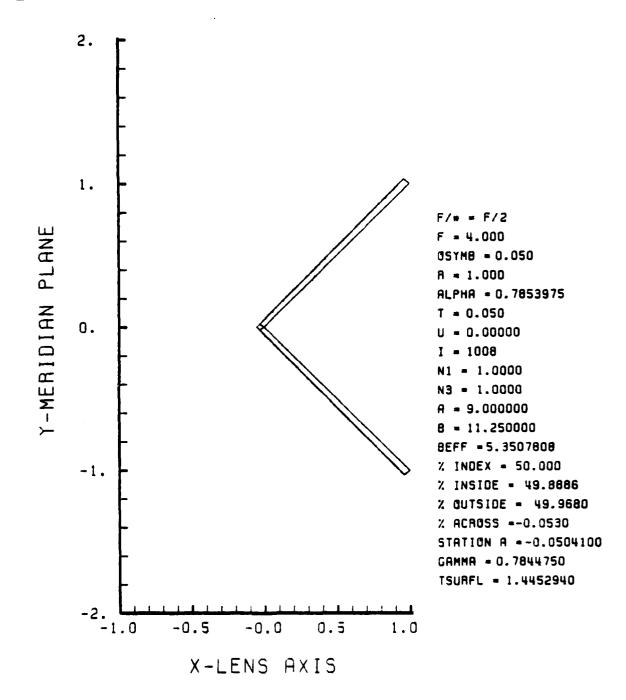
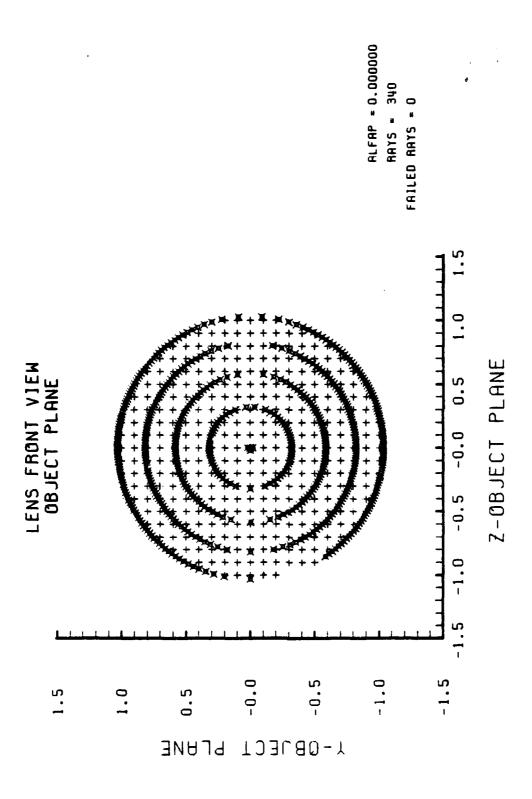


Figure G-1. "Best" GRIN Lens Shape with 50% Gradient, OB = 0.05 and a = 9.00 in the F/2 Configuration



Grid Plane at  $\alpha_p$  = 0.0 for Lens of Figure G-1 Figure G-2.

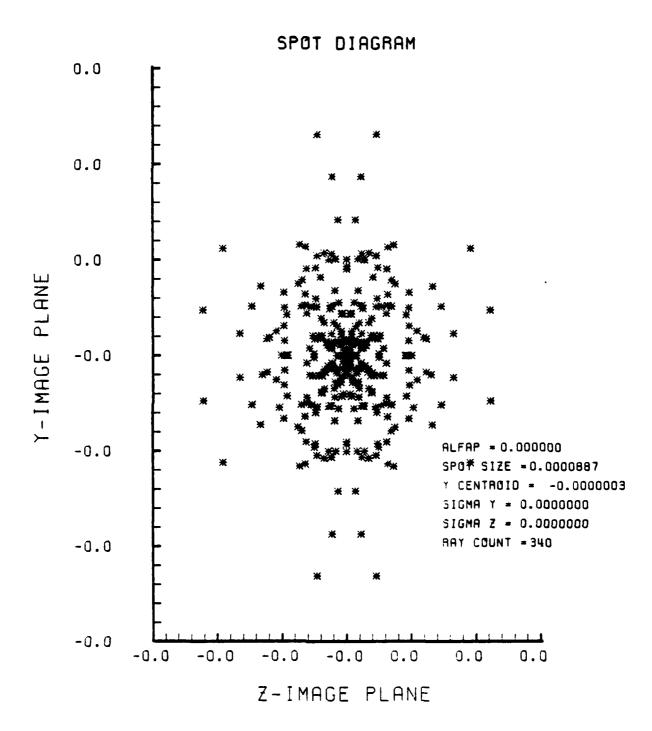


Figure G-3. Spot Diagram for Grid of Figure G-2

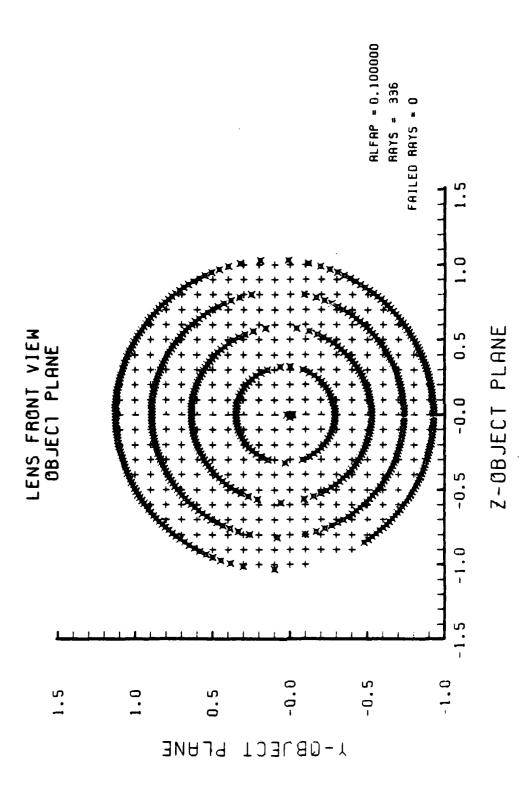


Figure G-4. Grid Plane at  $\alpha_p$  = 0.1 for Lens of Figure G-1

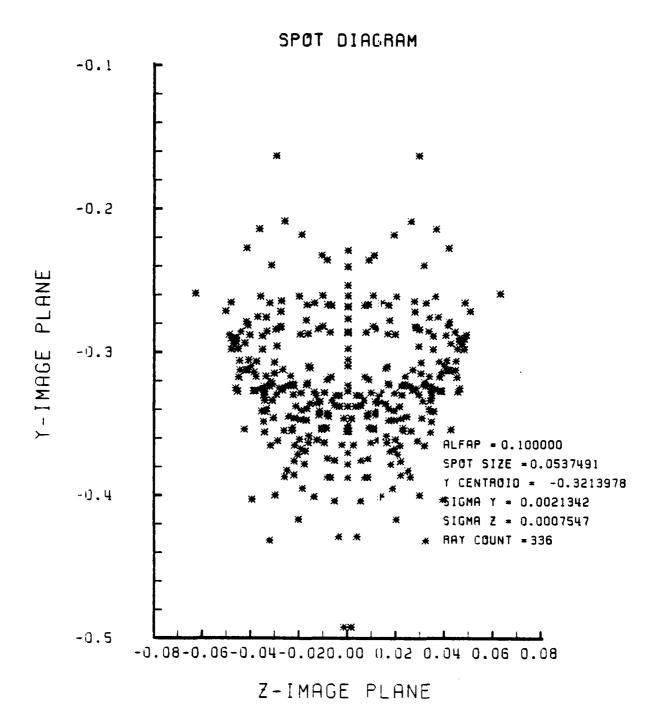


Figure G-5. Spot Diagram for Grid of Figure G-4

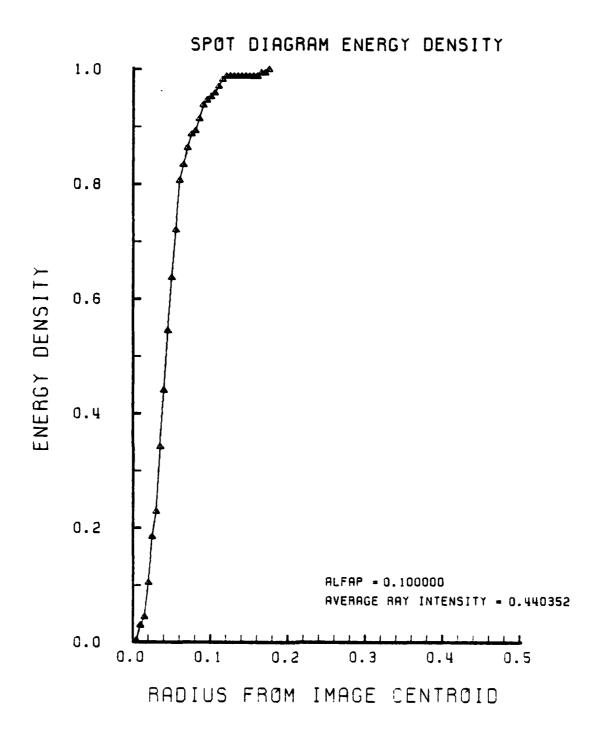
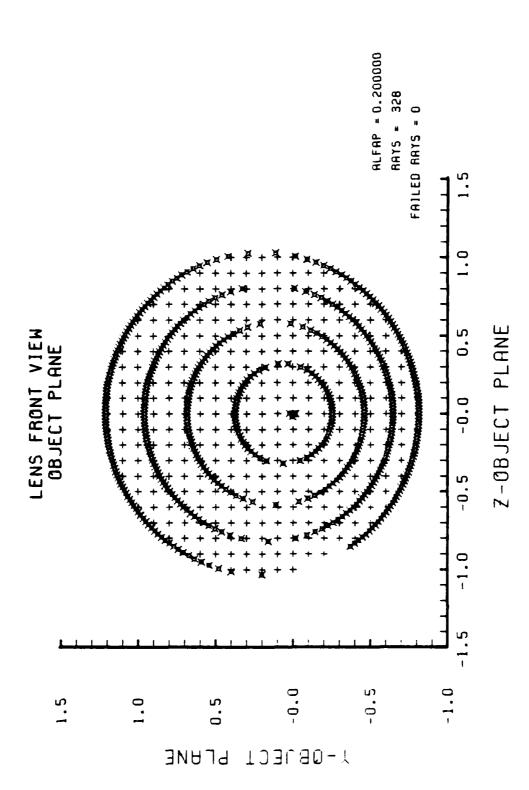


Figure G-6. Encircled Energy of Figure G-5



Grid Plane at  $\alpha_p$  = 0.2 for Lens of Figure G-1 Figure G-7.

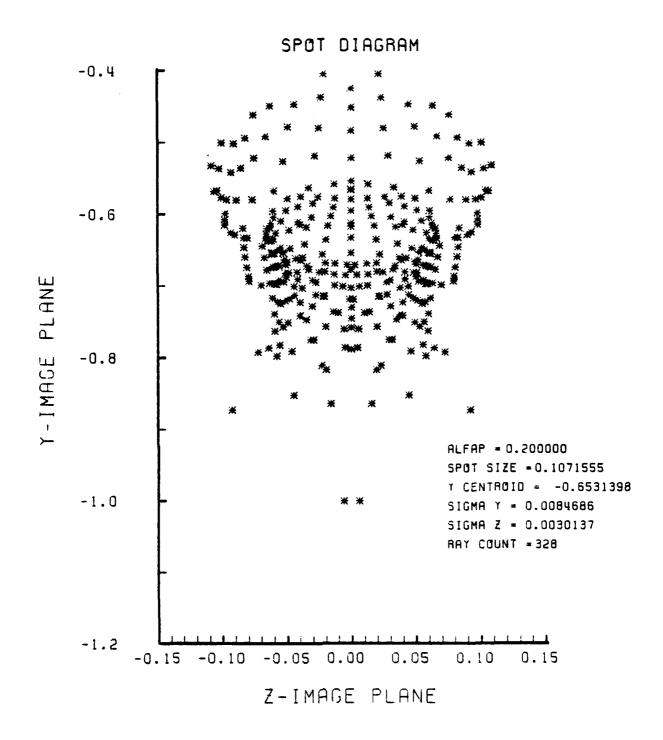


Figure G-8. Spot Diagram for Grid of Figure G-7

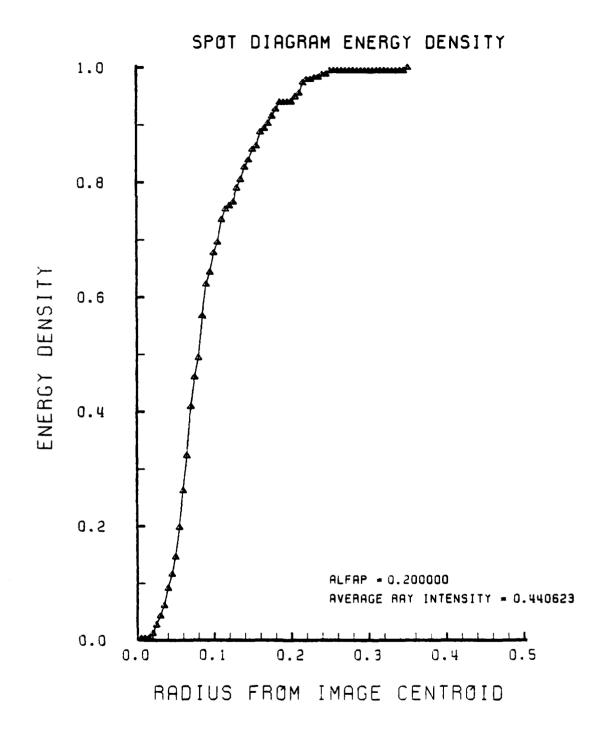


Figure G-9. Encircled Energy of Figure G-8

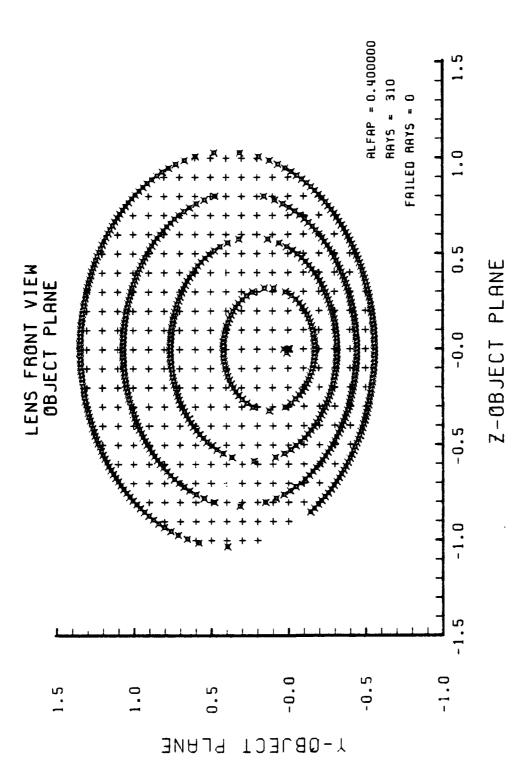


Figure G-10. Grid Plane at  $\alpha_{p}$  = 0.4 for Lens of Figure G-1

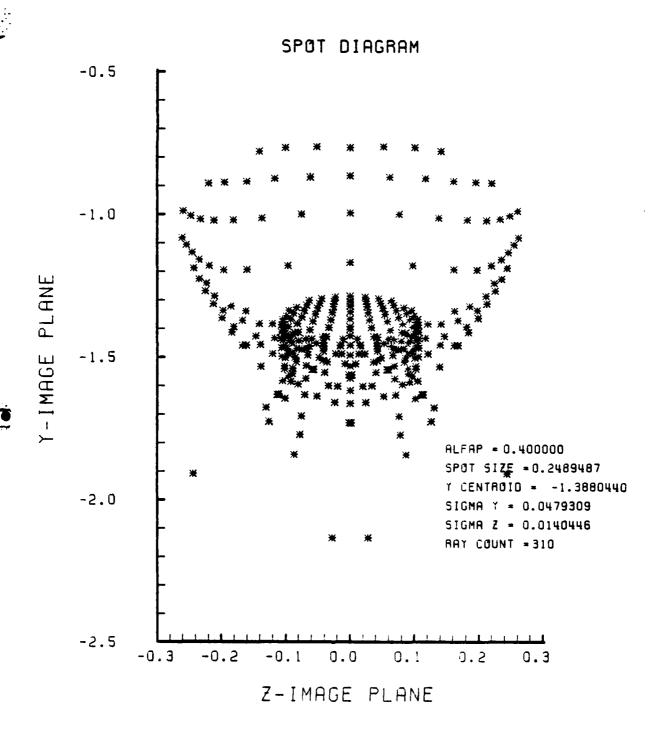


Figure G-11. Spot Diagram for Grid of Figure G-10

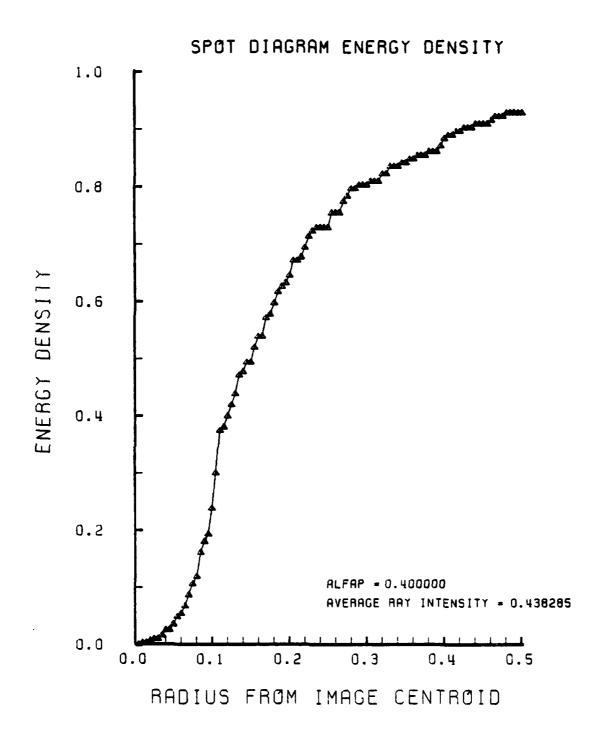
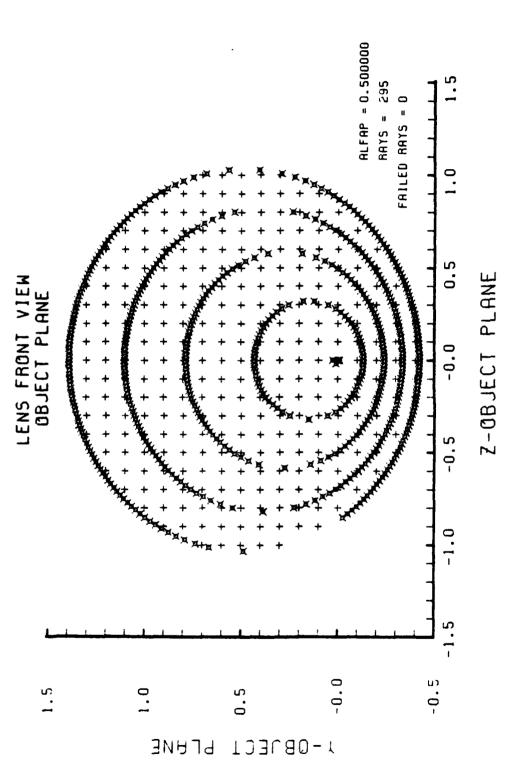


Figure G-12. Encircled Energy of Figure G-11



= 0.5 for Lens of Figure G-1 Grid Plane at  $\alpha_p$ Figure G-13.

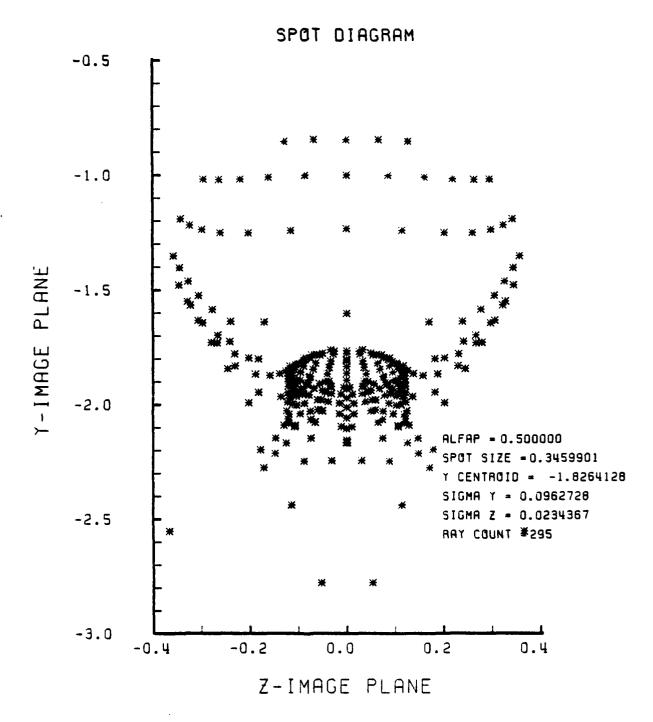
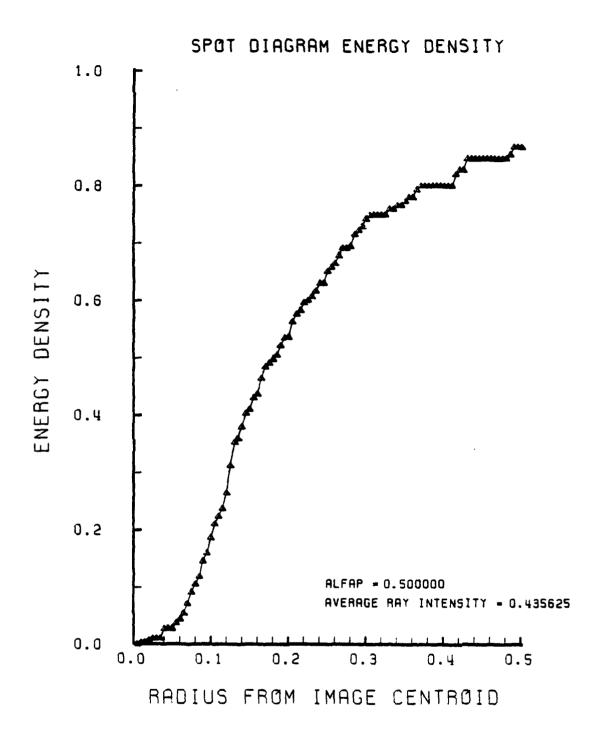


Figure G-14. Spot Diagram for Grid of Figure G-13



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Figure G-15. Encircled Energy of Figure G-14

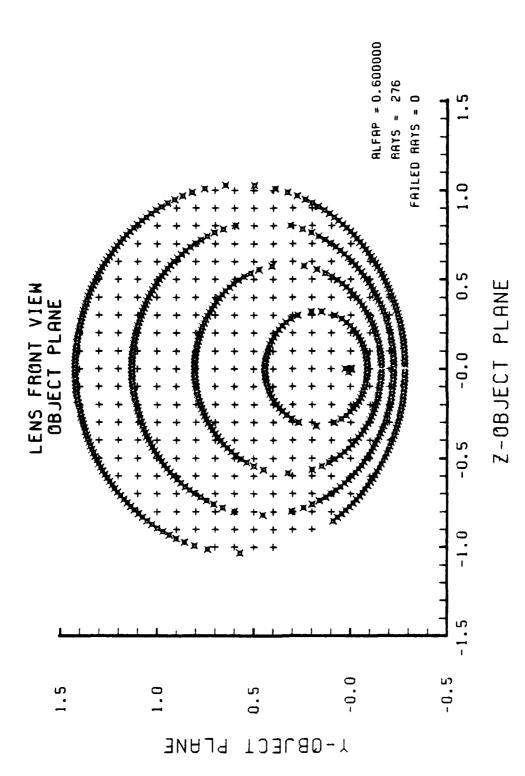


Figure G-16. Grid Plane at  $\alpha_{\rm p}$  = 0.7 for Lens of Figure G-1

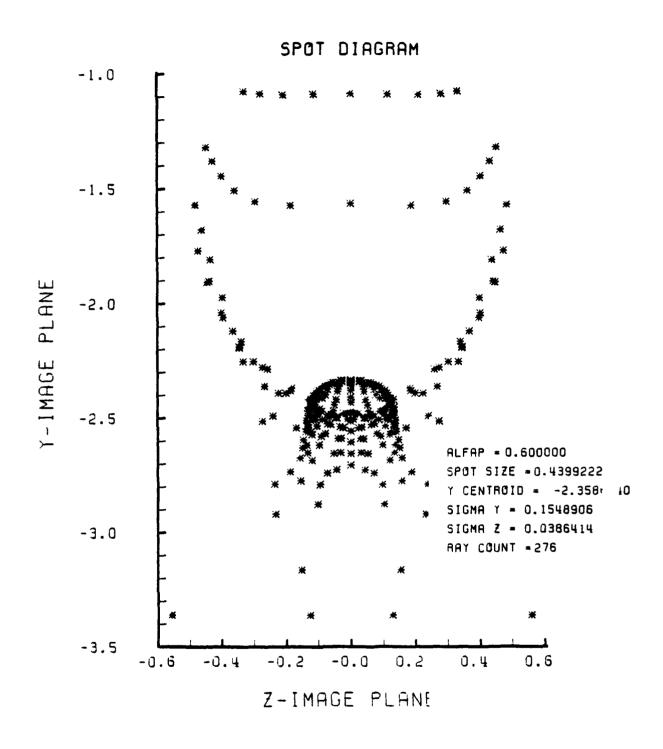


Figure G-17. Spot Diagram for Grid of Figure G-16

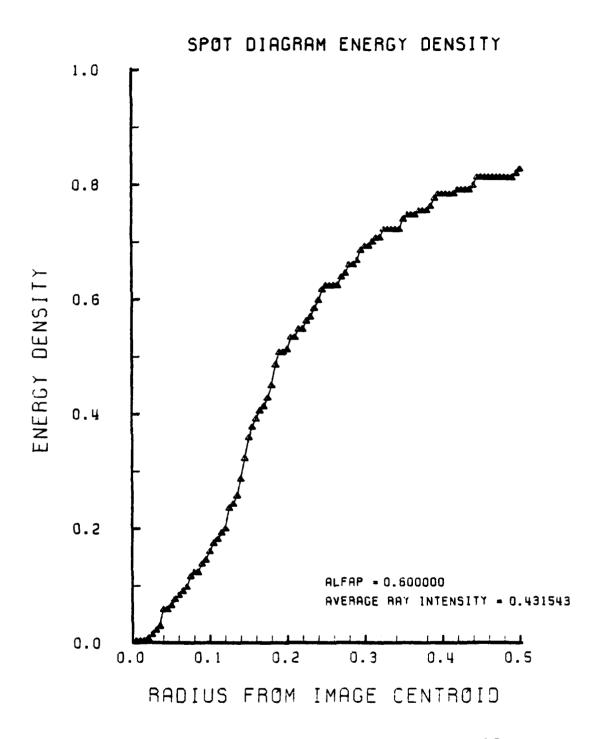
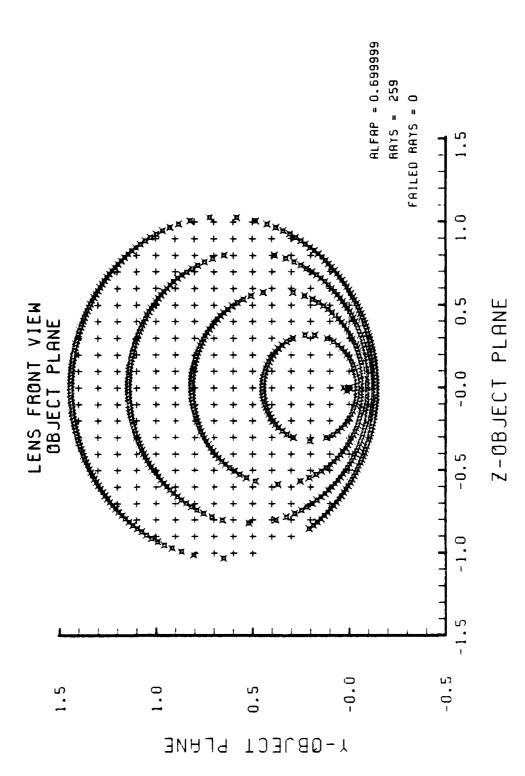


Figure G-18. Encircled Energy of Figure G-17



Grid Plane at  $\alpha_p = 0.7$  for Lens of Figure G-1 Figure G-19.

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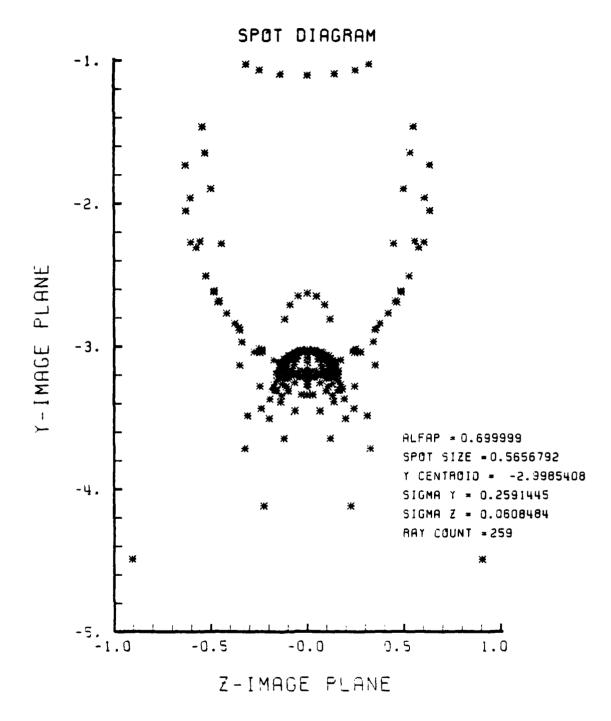


Figure G-20. Spot Diagram for Grid of Figure G-19

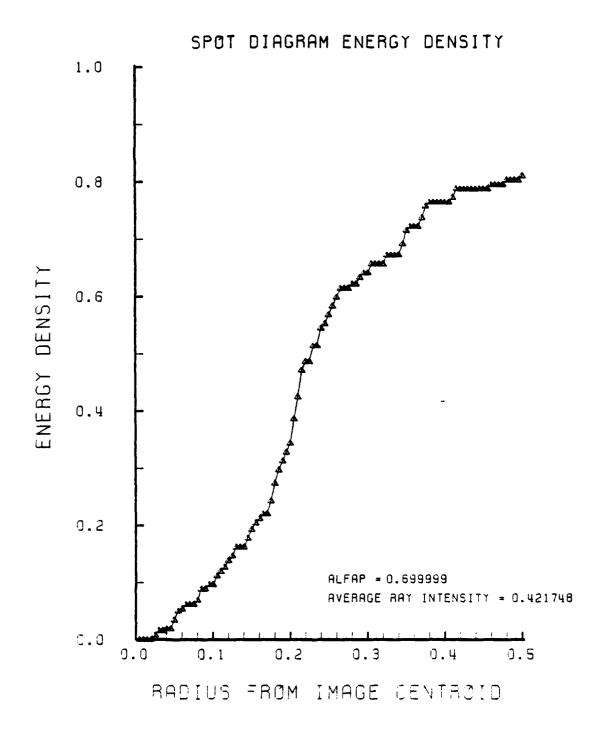
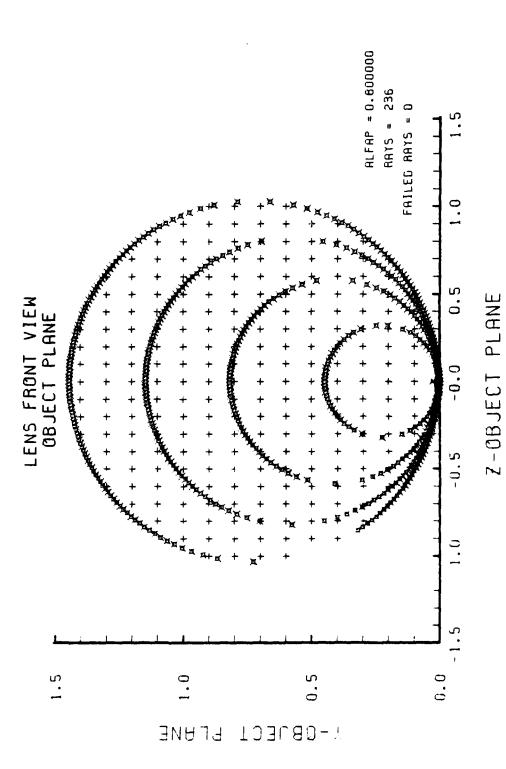


Figure G-21. Encircled Energy of Figure G-20



Grid Plane at  $\alpha_{p}$  = 0.8 for Lens of Figure G-1 Figure G-22.

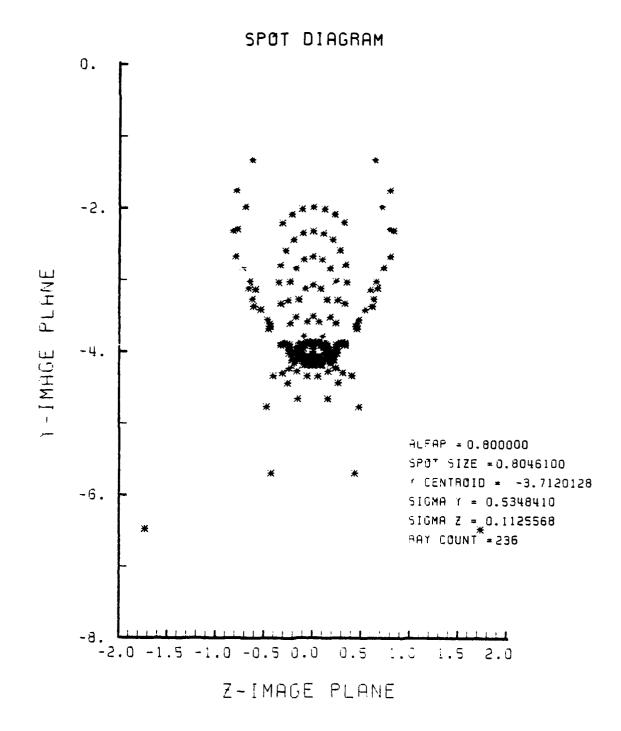


Figure G-23. Spot Diagram for Grid of Figure G-22

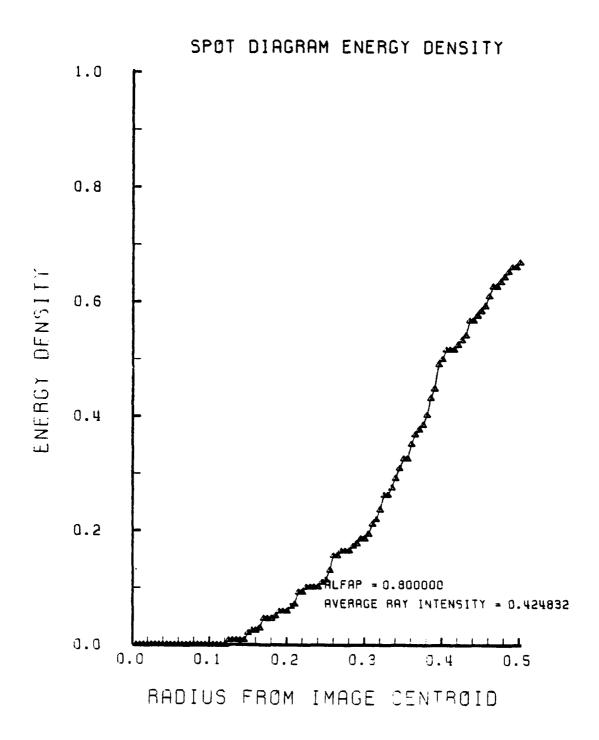


Figure G-24. Encircled Energy of Figure G-23

## APPENDIX H

## "BEST" GRIN LENS PERFORMANCE PLOTS IN THE F/1 CONFIGURATION

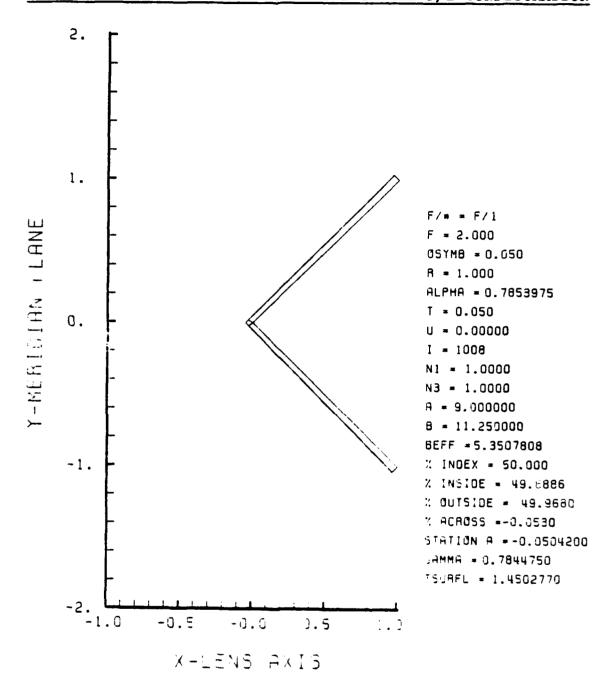
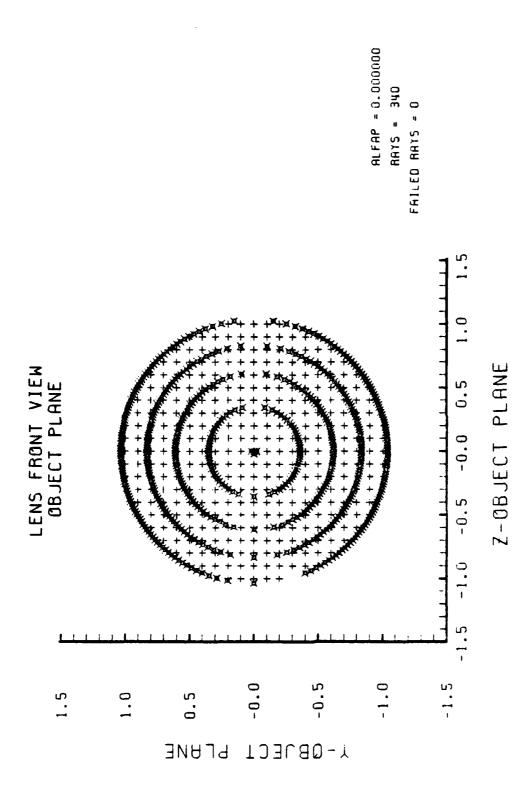


Figure H-1. "Best" GRIN Lens Shape with 50% Gradient, OB = 0.05, and a = 9.00 in the F/l Configuration



Grid Plane at  $\alpha_p$  = 0.0 for Lens of Figure H-l Figure H-2.

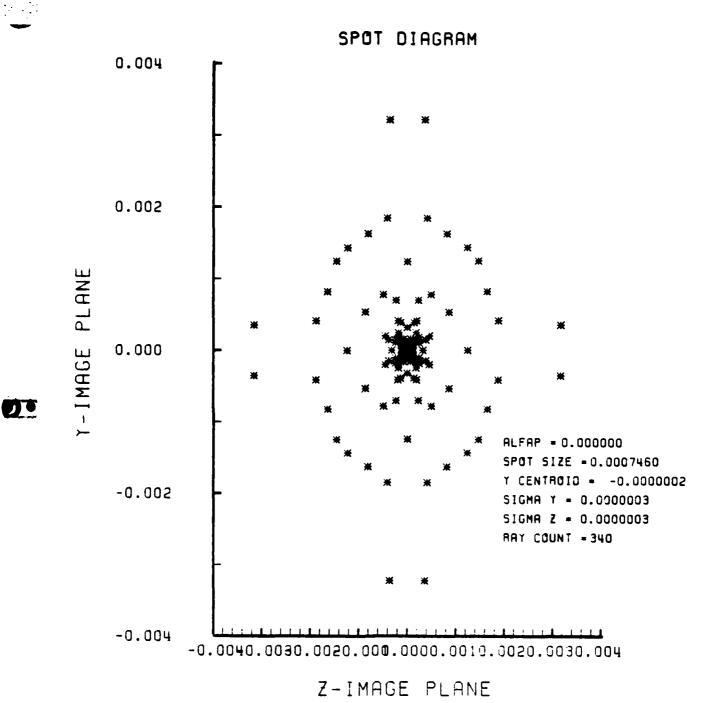
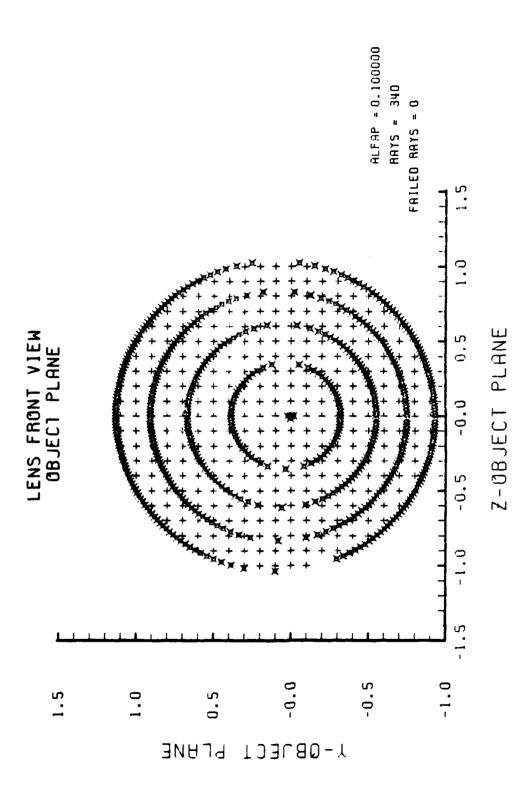


Figure H-3. Spot Diagram for Grid of Figure H-2



Grid Plane at  $\alpha_p$  = 0.1 for Lens of Figure H-1 Figure H-4.

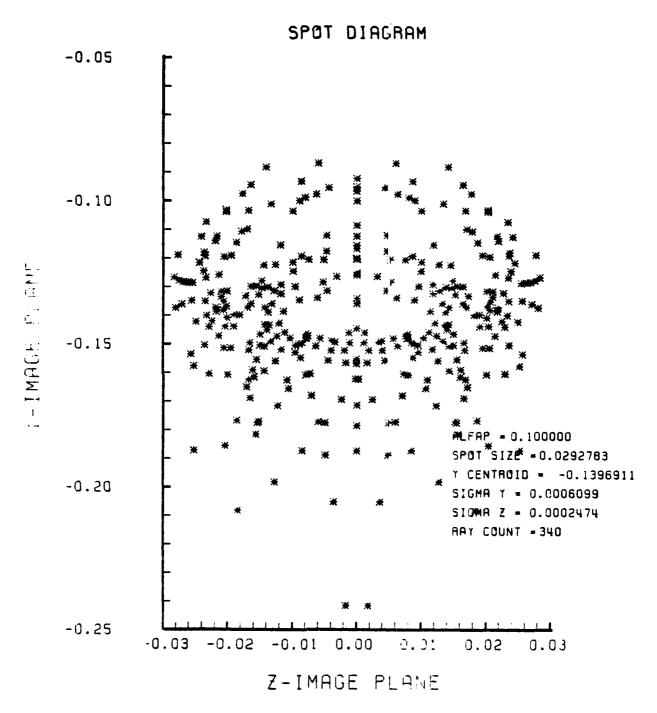


Figure H-5. Spot Diagram for Grid of Figure H-4

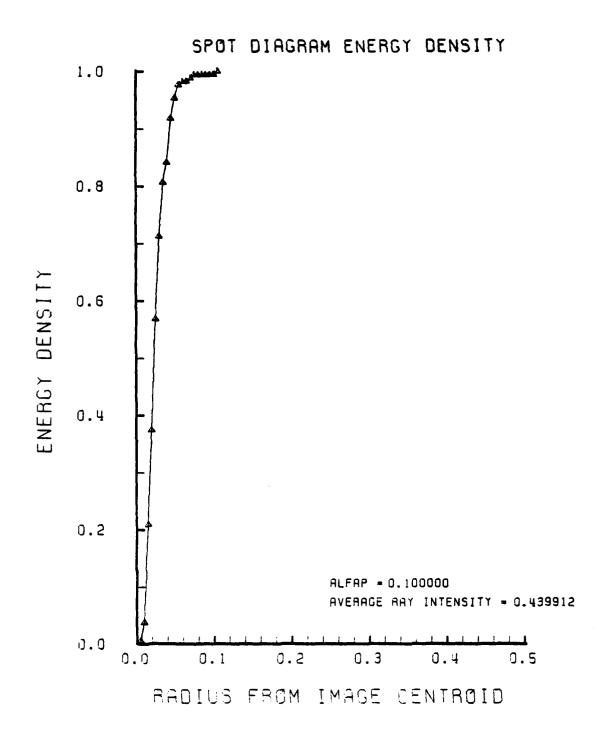
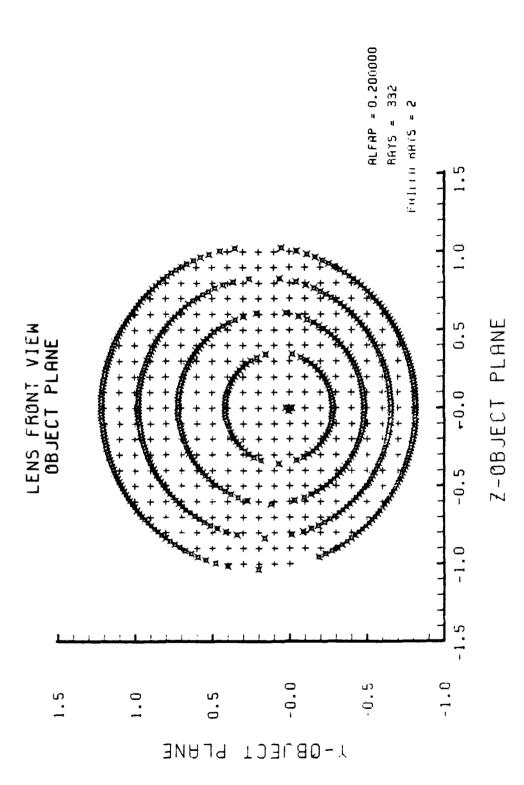


Figure H-6. Encircled Energy of Figure H-5



Grid Plane at  $\alpha_p$  = 0.2 for Lens of Figure H-l Figure H-7.

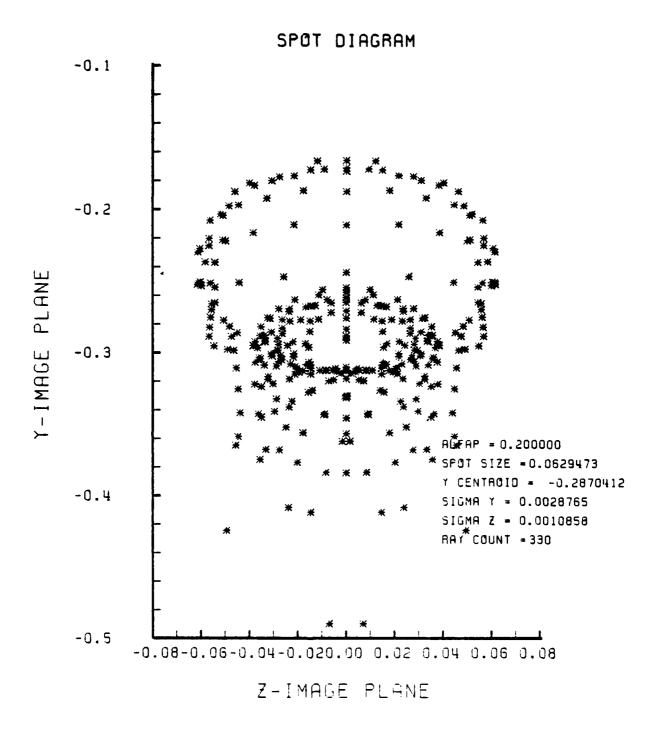


Figure H-8. Spot Diagram for Grid of Figure H-7

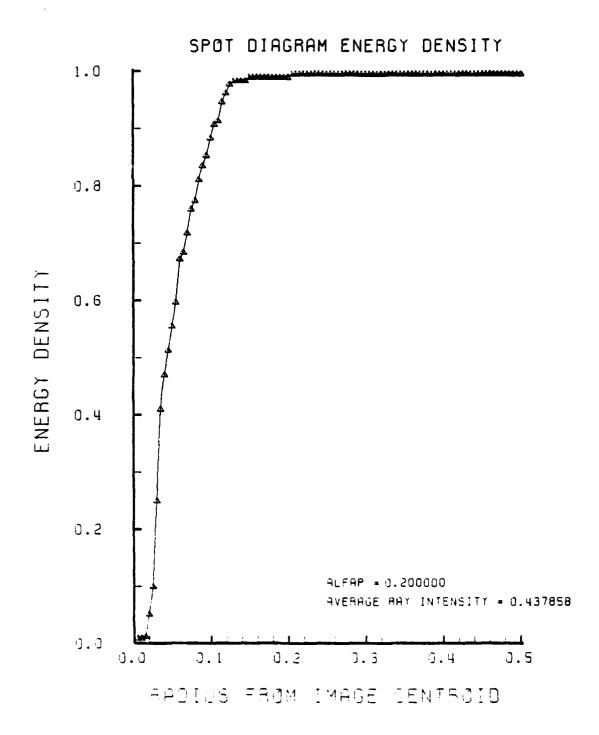
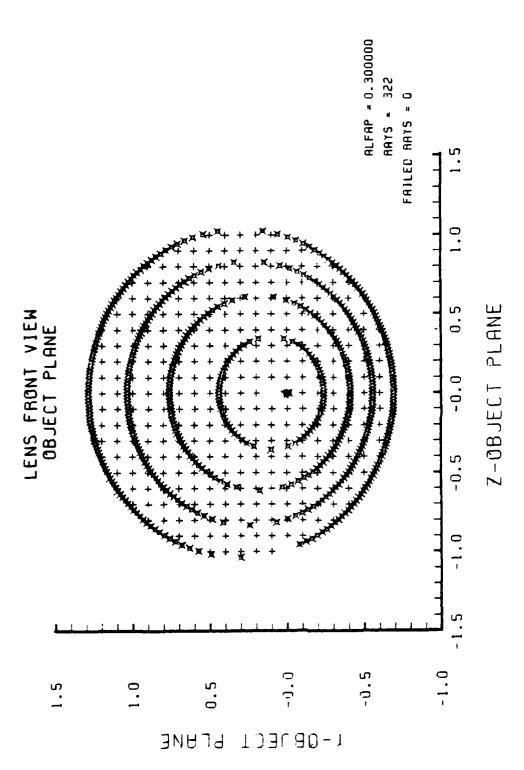


Figure H-9. Encircled Energy of Figure H-8



Grid Plane at  $\alpha_{\rm p}$  = 0.3 for Lens of Figure H-l Figure H-10.

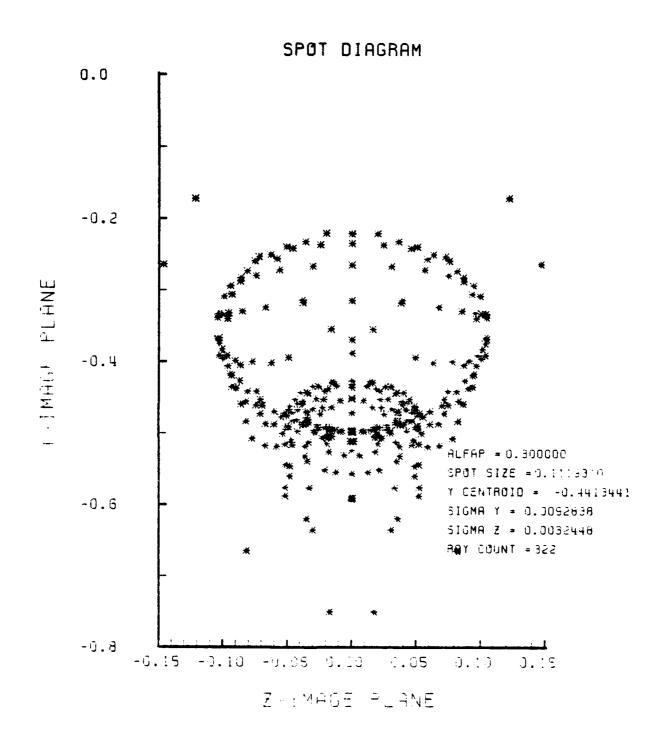


Figure H-11. Spot Diagram for Grid of Figure H-10

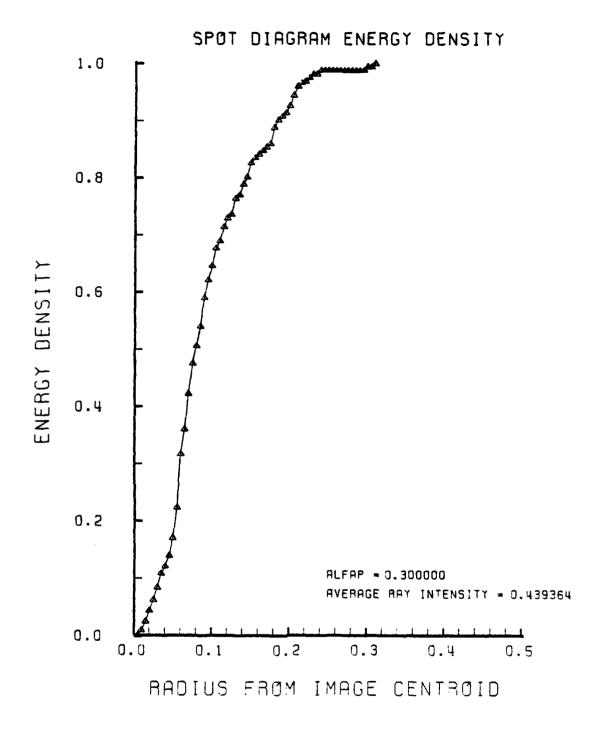


Figure H-12. Encircled Energy of Figure H-11

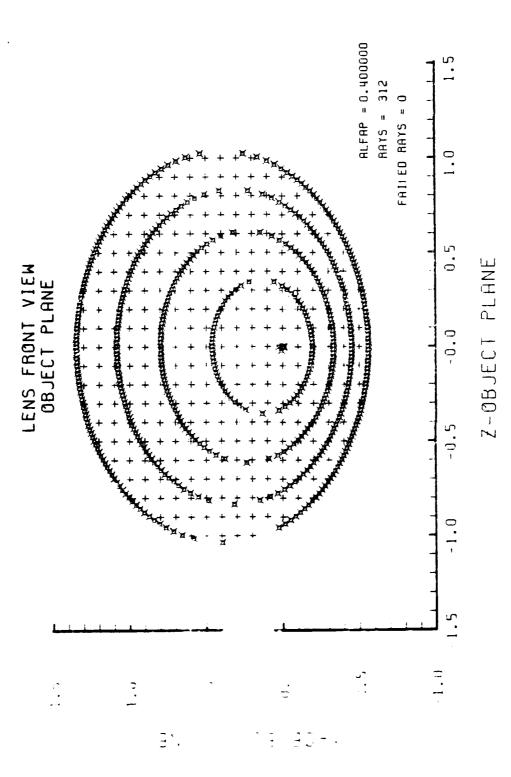


Figure H-13. Grid Plane at  $\alpha_{p}$  = 0.4 for Lens of Figure H-1

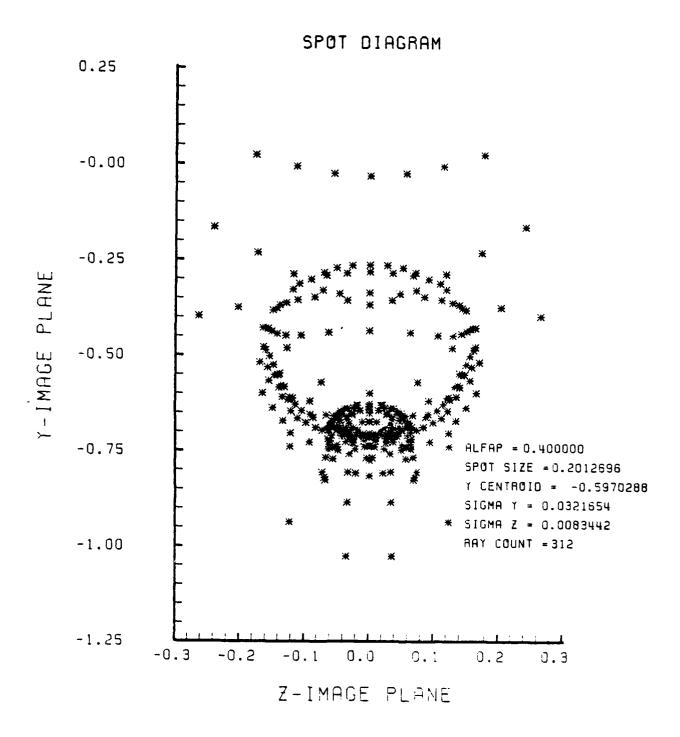


Figure H-14. Spot Diagram for Grid of Figure H-13

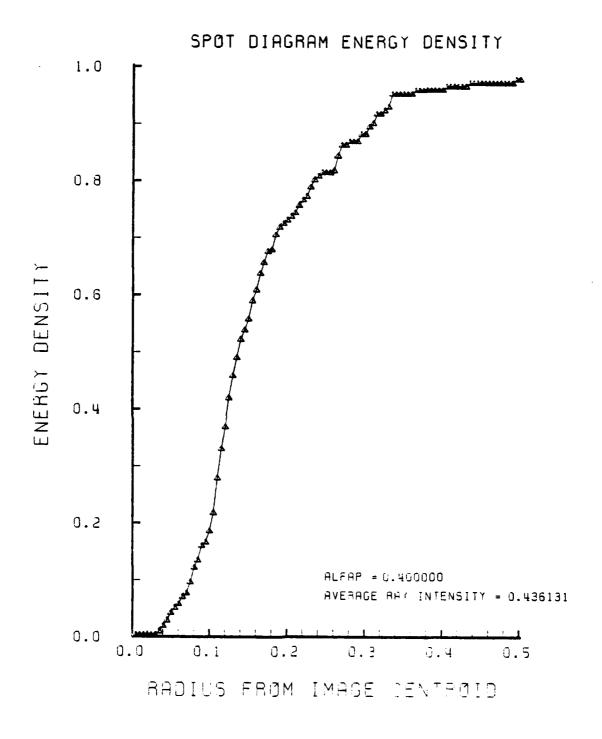
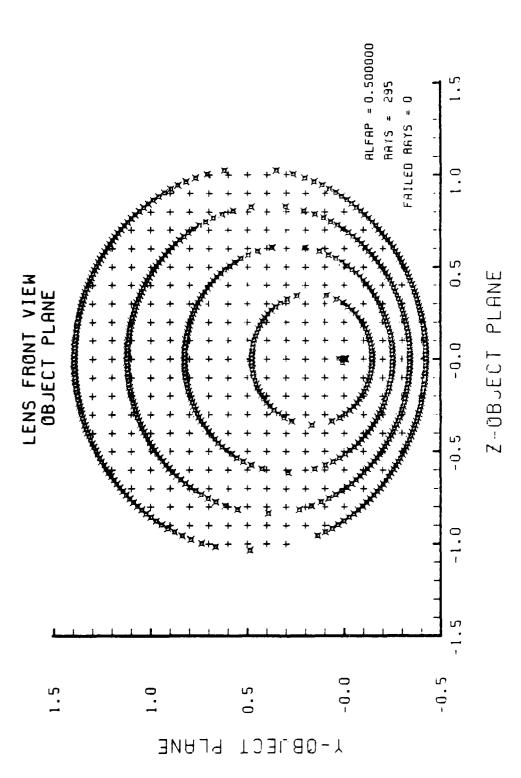


Figure H-15. Encircled Energy of Figure H-14



= 0.5 for Lens of Figure H-l Grid Plane at  $^{lpha}_{
m p}$ Figure H-16.

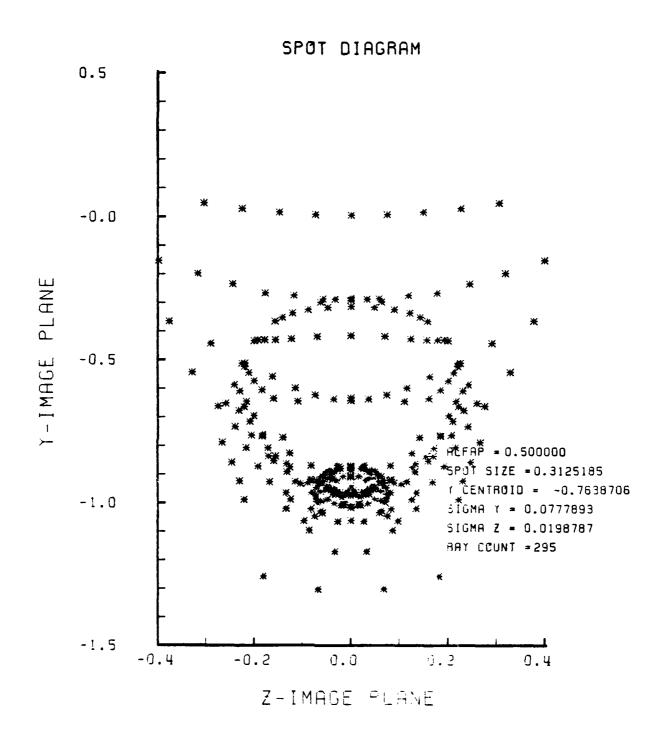


Figure H-17. Spot Diagram for Grid of Figure H-16

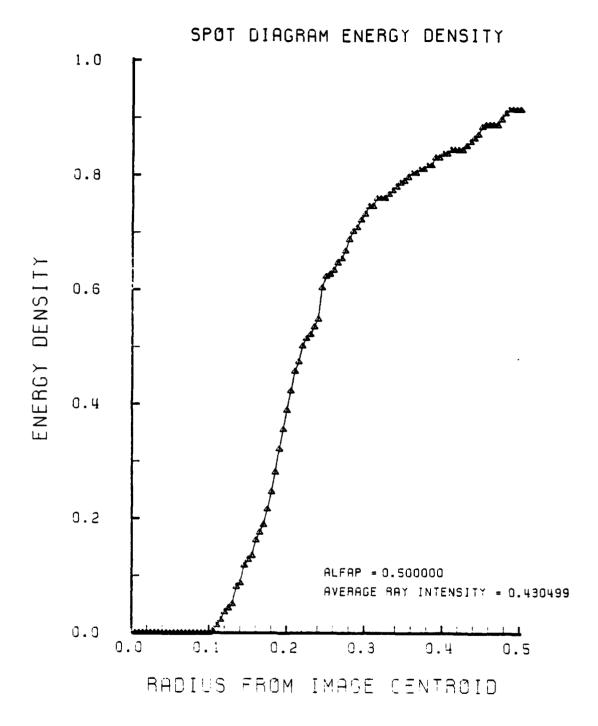
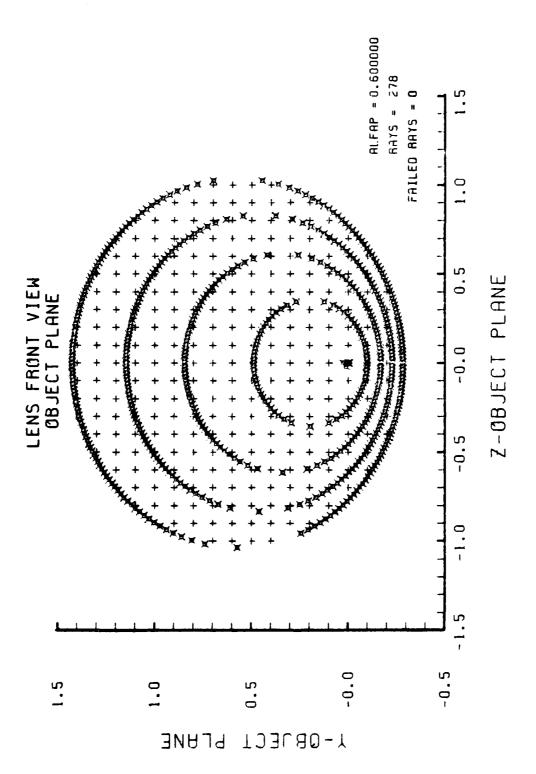


Figure H-18. Encircled Energy of Figure H-17



= 0.6 for Lens of Figure H-1 Grid Plane at  $^{lpha}_{
m p}$ Figure H-19.

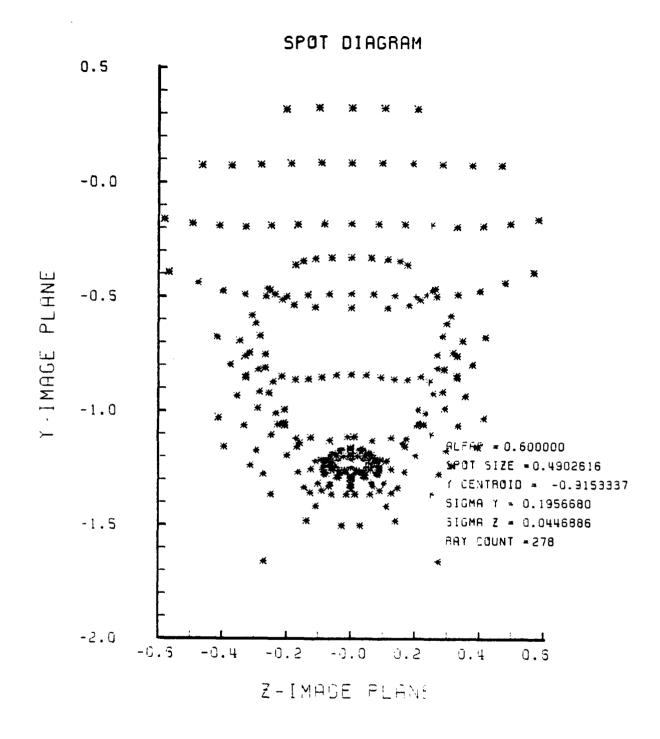


Figure H-20. Spot Diagram for Grid of Figure H-19

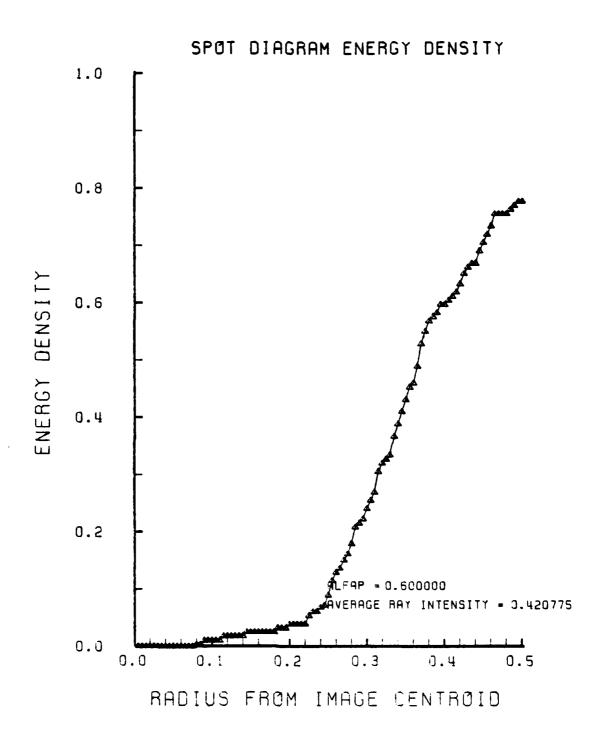
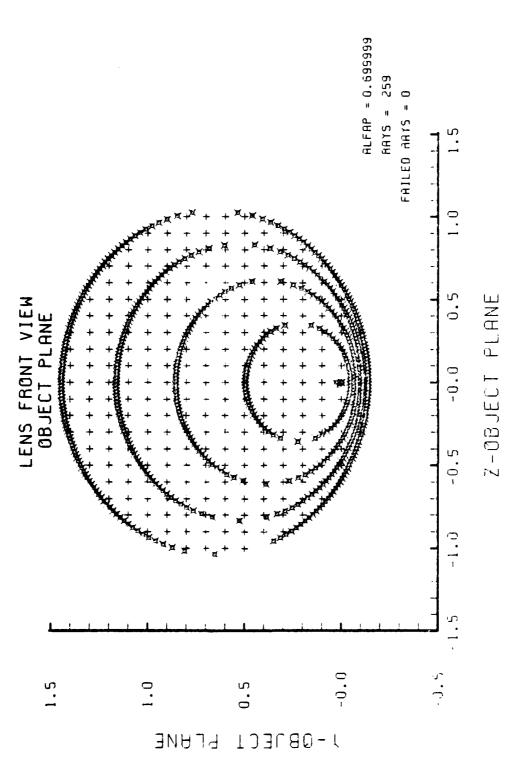


Figure H-21. Encircled Energy of Figure H-20



Grid Plane at  $\alpha_p$  = 0.7 for Lens of Figure H-l Figure H-22.

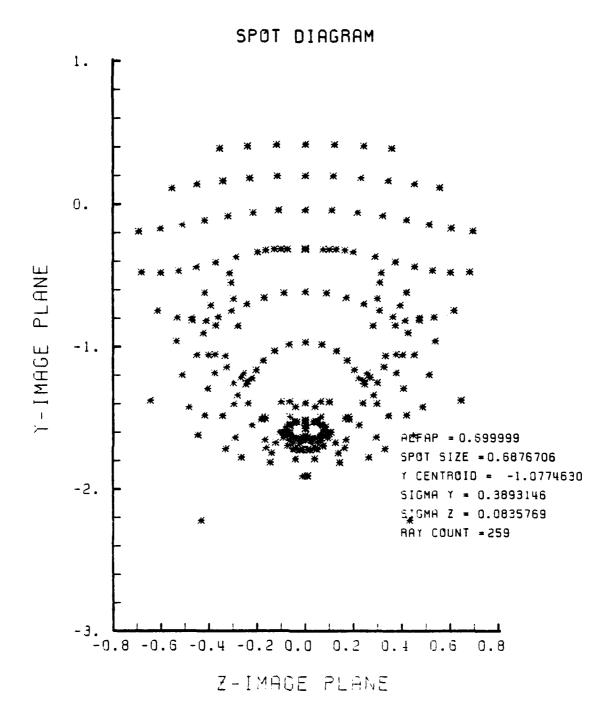


Figure H-23. Spot Diagram for Grid of Figure H-22

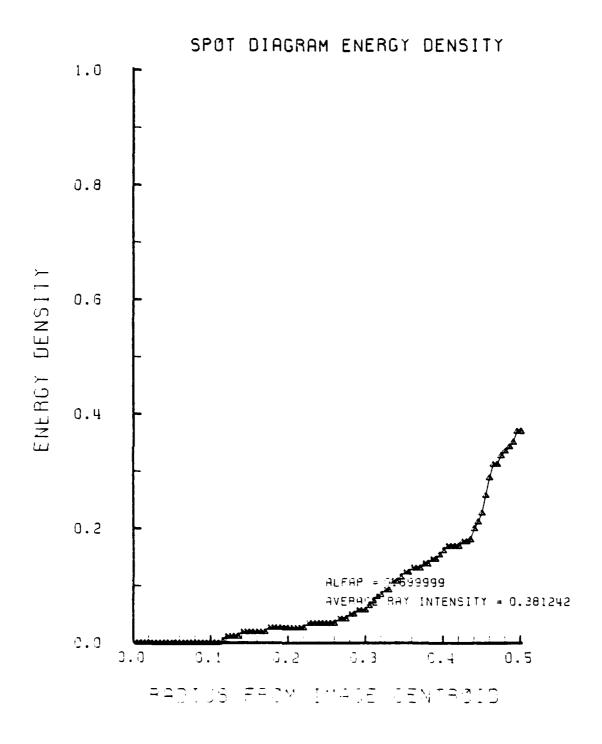
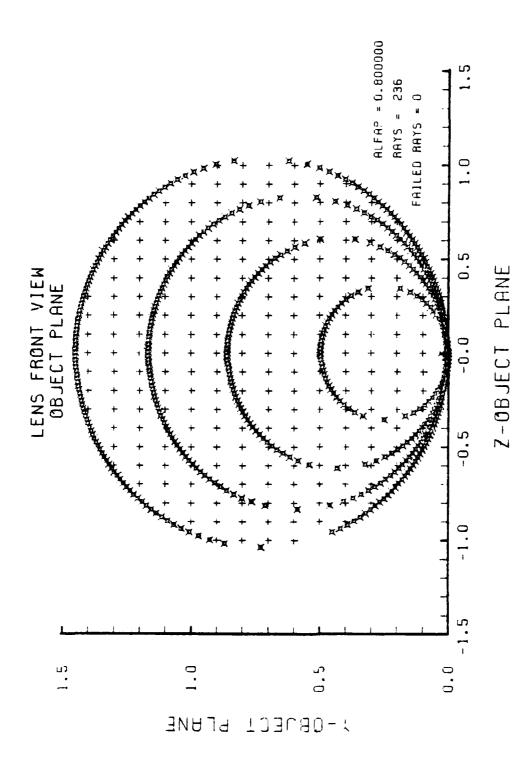


Figure H-24. Encircled Energy of Figure H-23



= 0.8 for Lens of Figure H-1 Grid Plane at  $^{lpha}_{
m p}$ Figure H-25.

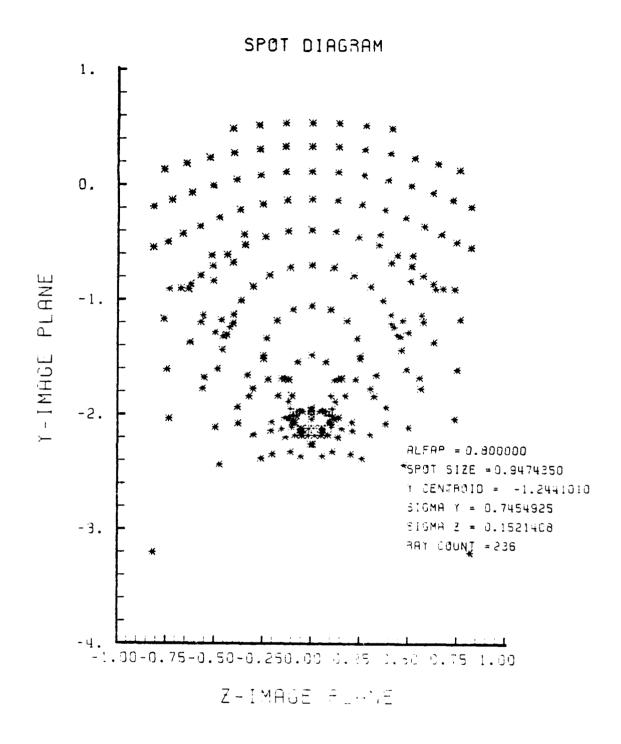
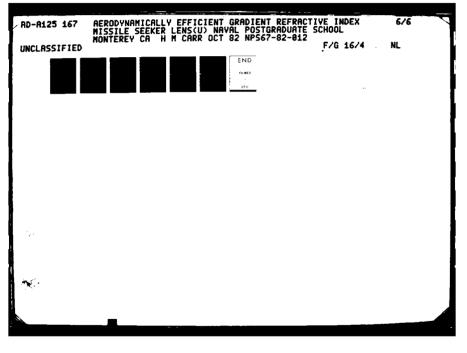
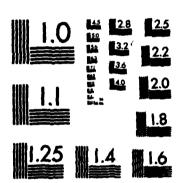


Figure H-26. Spot Diagram for Grid of Figure H-25





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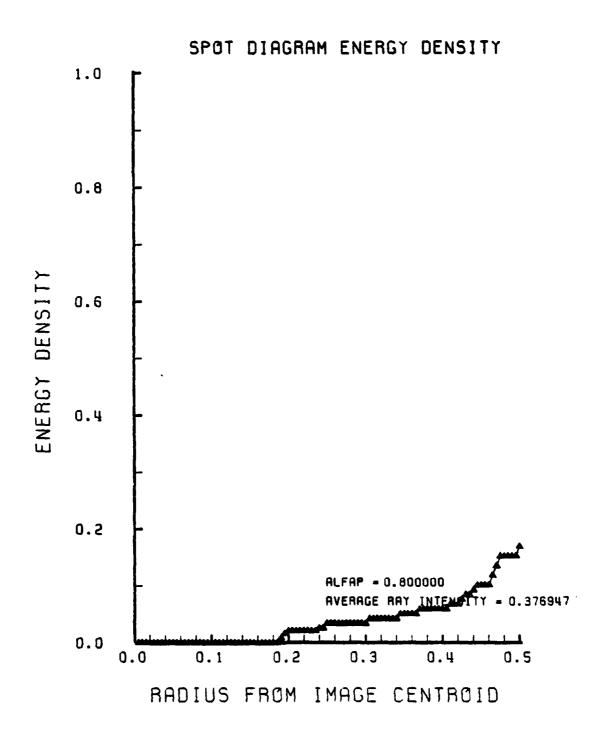


Figure H-27. Encircled Energy of Figure H-26

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